

Quasi-biennial oscillations extracted from the monthly NAO index and temperature records are phase-synchronized

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Abstract. Using the extension of Monte Carlo Singular System Analysis (MC SSA), based on evaluating and testing the regularity of the dynamics of the SSA modes against the colored noise null hypothesis, we demonstrate detection of oscillatory modes with period of about 27 months in records of monthly average near-surface air temperature from several European locations, as well as in the monthly North Atlantic Oscillation index. According to their period, the detected modes can be attributed to the quasi-biennial oscillations (QBO). The QBO modes extracted from the temperature and from the NAO index underwent synchronization analysis and their phase synchronization has been confirmed with high statistical significance.

1 Introduction

The North Atlantic Oscillation (NAO) is a dominant pattern of atmospheric circulation variability in the extratropical Northern Hemisphere, accounting for about 60% of the total sea-level pressure variance. The NAO has a strong effect on European weather conditions, influencing meteorological variables including temperature (Hurrell et al., 2001). The NAO-temperature relationship, however, is not straightforward and its mechanism is not yet fully understood. Recent studies show that the relationship between the NAO and the winter Northern Hemisphere surface temperature (NHT) might be a function of the NAO state (Pozo-Vázquez et al., 2001). There are energetic modes of coherent variability between temperature and the NAO in the time domain of 2–6 years for 1857–1879 and 1978–1984, and in the domain of 6–10 years from 1961 to 1991 (Pozo-Vázquez et al., 2001). Changes of the strength and character of correlations be-

tween the NAO and the NHT seems to be modulated by the phase of the solar activity cycle (Gimeno et al., 2003).

The long-term dynamical behaviour of the NAO phenomenon itself is far from being understood and is also investigated in connection with studies of predictability of the NAO and the ability of climate numerical models of simulating it (Eden et al., 2002; Bojariu and Gimeno, 2003). Some authors are pessimistic, claiming that the NAO index can be considered just as pink noise with very little predictability (Fernandez et al., 2003). On the other hand, Gámiz-Fortis et al. (2002) have applied the Monte Carlo Singular System Analysis (MC SSA) to the winter NAO index, i.e., yearly sampled values obtained by averaging December, January and February index values, and, using an embedding window of length $n=40$ years, were able to identify an oscillatory mode with a period of about 7.7 years. Paluš and Novotná (2004) have introduced the so called enhanced MC SSA, based on evaluating and testing the regularity of the dynamics of the SSA modes against the colored noise null hypothesis, in addition to the test based on variance (eigenvalues). The application of the regularity index, computed from a coarse-grained estimation of mutual information, enhances the test sensitivity and reliability in the detection of dynamical modes which are relatively more regular than those obtained by decomposition of colored noise, in particular, in the detection of irregular oscillations embedded in red noise. This enhanced MC SSA was successfully applied in the detection of oscillatory modes of period 7.8 years in records of monthly average near-surface air temperature from several European locations, as well as in the monthly North Atlantic Oscillation index (Paluš and Novotná, 2004).

In this paper we continue and refine the enhanced MC SSA analysis of the temperature records and the monthly North Atlantic Oscillation index by focusing on higher (than near-decadal) frequencies and demonstrate the presence of several oscillatory modes with periods in the range from 20 to 64 months. The most prominent oscillatory mode (besides the

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above mentioned mode with period 7.8 years) is the mode with the average period 27 months. We identify this mode with the quasi-biennial oscillation (QBO) and extract this oscillatory mode from both the temperature data and the NAO index. The extracted oscillatory modes (projections on the SSA/EOF basis) are further studied by the means of synchronization analysis (Pikovsky et al., 2001) and the existence of phase synchronization between the QBO modes from the temperature and from the NAO index is demonstrated and statistically proven.

A brief introduction into Monte Carlo singular system analysis and its enhancement is given in Sect. 2. The analyzed data are described in Sect. 3. Section 4 summarizes the application of enhanced MC SSA to monthly average near-surface temperature records and to the monthly NAO index. Synchronization analysis is introduced and applied to the QBO modes in Sect. 5. The results are discussed in Sect. 6 and conclusions are given in Sect. 7.

2 Monte Carlo singular system analysis

Singular system (or singular spectrum) analysis (SSA) in its original form (also known as principal component analysis, or Karhunen-Loève decomposition) is a method for the identification and distinction from noise of important information in multivariate data. It is based on an orthogonal decomposition of a covariance matrix of multivariate data under study. SSA provides an orthogonal basis onto which the data can be transformed, thus making individual data components (“modes”) linearly independent. Each of the orthogonal modes (projections of the original data onto new orthogonal basis vectors) is characterized by its variance, which is given by the related eigenvalue of the covariance matrix.

Here we will deal with a univariate version of SSA in which the analyzed data is a univariate time series and the decomposed matrix is a time-lag covariance matrix, i.e., instead of several components of multivariate data, a time series and its time-lagged versions are considered. This form of SSA, which has frequently been used in the field of meteorology and climatology (Vautard and Ghil, 1989; Ghil and Vautard, 1991; Keppenne and Ghil, 1992; Yiou et al., 1994; Allen and Smith, 1994), can provide a decomposition of the studied time series into orthogonal components (modes) with different dynamical properties, and thus “interesting” phenomena such as slow modes (trends) and regular or irregular oscillations (if present in the data) can be identified and retrieved from the background of noise and/or other “uninteresting” non-specified processes.

In traditional SSA, the distinction of “interesting” components (signal) from noise is based on finding a threshold (jump-down) to a “noise floor” in a sequence of eigenvalues given in a descending order. This approach might be problematic if the signal-to-noise ratio is not sufficiently large, or the noise present in the data is not white but “colored”.

For such cases, statistical approaches utilizing Monte Carlo simulation techniques have been proposed (Ghil and Vautard, 1991; Vautard et al., 1992) for reliable signal/noise separation. The particular case of Monte Carlo SSA (MC SSA) that considers “red” noise, usually present in geophysical data, has been introduced by Allen and Smith (1996).

Now, we present a few necessary details of the SSA method in the form of a technical recipe:

Take the analyzed time series $\{y(i)\}$, $i=1, \dots, N_0$, and construct a map (“embedding”) into a space of n -dimensional vectors $\mathbf{x}(i)$ with components $x^k(i)$, given as

$$x^k(i) = y(i + k - 1), \quad (1)$$

where $k=1, \dots, n$; and n is the embedding dimension.

Construct a symmetric $n \times n$ matrix $\mathbf{C}=\mathbf{X}^T\mathbf{X}$, with elements:

$$c_{kl} = (1/N) \sum_{i=1}^N x^k(i)x^l(i), \quad (2)$$

where $1/N$ is the proper normalization and the components $x^k(i)$, $i=1, \dots, N$, are supposed to have a zero mean. The symmetric matrix \mathbf{C} can be decomposed as

$$\mathbf{C} = \mathbf{V}\mathbf{\Sigma}\mathbf{V}^T, \quad (3)$$

where the $n \times n$ matrix $\mathbf{V}=\{v_{ij}\}$ gives an orthonormal basis in the space of vectors $\mathbf{x}(i)$, $\mathbf{\Sigma}=\text{diag}(\sigma_1, \sigma_2, \dots, \sigma_n)$, σ_i are non-negative eigenvalues giving the variance of orthogonal modes

$$\xi^k(i) = \sum_{l=1}^n v_{lk}x^l(i), \quad (4)$$

into which the original series can be decomposed. For more details, see, e.g., Vautard et al. (1992).

Of course, the original time series $x^k(i)$ can be reconstructed from the modes, as

$$x^k(i) = \sum_{l=1}^n v_{kl}\xi^l(i). \quad (5)$$

In Eq. (5), the modes $\xi^k(i)$ can also be interpreted as time-dependent coefficients and the orthogonal vectors $\mathbf{v}_k=\{v_{kl}\}$ as basis functions, usually called the empirical orthogonal functions (EOFs).

The clear signal/noise distinction based on the eigenvalues $\sigma_1, \sigma_2, \dots, \sigma_n$ can only be obtained in particularly idealized situation when the signal/noise ratio is large enough and the background consists of white noise. In many geophysical processes, however, so-called “red” noise with power spectrum of the $1/f$ type is present (Allen and Smith, 1996). Its SSA eigenspectrum also has the $1/f$ character (Gao et al., 2003), i.e., an eigenspectrum of red noise is equivalent to a coarsely discretized power spectrum, where the number of frequency bins is given by the embedding dimension n . The eigenvalues related to the slow modes are much larger than

the eigenvalues of the modes related to higher frequencies. Thus, in the classical SSA approach applied to red noise, the eigenvalues of the slow modes might incorrectly be interpreted as a (nontrivial) signal, or, on the other hand, a nontrivial signal embedded in red noise might be neglected if its variance is smaller than the slow-mode eigenvalues of the background red noise. Therefore, Allen and Smith (1996) proposed comparing the SSA spectrum of the analyzed signal with SSA spectra of a red-noise model fitted to the studied data. Such a red-noise process can be modeled by using an AR(1) model (autoregressive model of the first order):

$$u(i) - \hat{u} = \alpha(u(i-1) - \hat{u}) + \gamma z(i), \quad (6)$$

where \hat{u} is the process mean, α and γ are process parameters, and $z(i)$ is Gaussian white noise with a zero mean and a unit variance.

In order to correctly detect a signal in red noise, the following approach has been proposed (Allen and Smith, 1996):

First, the eigenvalues are plotted not according to their values, but according to a frequency associated with a particular mode (EOF), i.e., the eigenspectrum in this form becomes a sort of a (coarsely) discretized power spectrum in general, not only in the case of red noise (when the eigenspectrum naturally has this form, as mentioned above).

Second, an eigenspectrum obtained from studied data is compared, in a frequency-by-frequency way, with eigenspectra obtained from a set of realizations of an appropriate noise model (such as the AR(1) model (6)), i.e., an eigenvalue related to a particular frequency bin obtained from the data is compared with a range of eigenvalues related to the same frequency bin, obtained from the set of realizations of the chosen AR(1) model.

The detection of a nontrivial signal in an experimental time series becomes a statistical test in which the null hypothesis that the experimental data were generated by a chosen noise model is tested. The realizations of the considered noise model (“null hypothesis”), i.e., the artificial data generated by the chosen noise model, are usually called “surrogate data” (Theiler et al., 1992; Allen and Smith, 1996; Paluš, 1995; Paluš and Novotná, 2004). When an eigenvalue associated with some frequency bin differs with a statistical significance from the range of related noise model eigenvalues, then one can infer that the studied data cannot be fully explained by the null hypothesis and could contain an additional (nontrivial) signal.

The above MC SSA is a sophisticated technique, but it still assumes that the signal of interest has been linearly added to a specified noise background and therefore that the variance in the frequency band, characteristic of the signal, is significantly greater than the typical variance in this frequency band obtained from the noise model. If the studied signal has a more complicated origin, e.g., when an oscillatory mode is embedded into a background process without significantly increasing variance in a particular frequency band, the standard MC SSA can fail. In order to be able to detect any inter-

esting dynamical mode independently of its (relative) variance, Paluš and Novotná (2004) have proposed testing also dynamical properties of the SSA modes against the modes obtained from surrogate data. In their particular implementation, the dynamics of the modes is characterized by their predictability (or regularity) measured by means of information theory.

The mutual information $I(X; Y)$ of two random variables X and Y is given by $I(X; Y) = H(X) + H(Y) - H(X, Y)$, where the entropies $H(X)$, $H(Y)$, $H(X, Y)$ are defined in the usual Shannonian sense (Cover and Thomas, 1991). For a time series $\{x(t)\}$, considered as a realization of a stationary and ergodic stochastic process $\{X(t)\}$, $t=1, 2, 3, \dots$, we compute the mutual information $I(x; x_\tau)$ as a function of time lag τ . In the following, we will mark $x(t)$ as x and $x(t+\tau)$ as x_τ . Let us find such τ_{\max} that for $\tau' \geq \tau_{\max}$: $I(x; x_{\tau'}) \approx 0$ for the analyzed datasets. Then we define the regularity index to be the norm of the mutual information:

$$\|I(x; x_\tau)\| = \frac{\Delta\tau}{\tau_{\max} - \tau_{\min} + \Delta\tau} \sum_{\tau=\tau_{\min}}^{\tau_{\max}} I(x; x_\tau) \quad (7)$$

with $\tau_{\min} = \Delta\tau = 1$ (sample) as a usual choice.

Since the mutual information $I(x; x_\tau)$ measures the average amount of information contained in the process $\{X\}$ about its future τ time units ahead, the regularity index $\|I(x; x_\tau)\|$ gives an average measure of predictability of the studied signal and is inversely related to the signal's entropy rate, i.e., to the rate at which the system, or process, producing the studied signal “forgets” information about its previous states (Paluš, 1996a).

Finally, we realize the enhanced MC SSA as follows:

1. The studied time series undergoes SSA as briefly described above, or, in detail in Paluš and Novotná (2004), i.e., using an embedding window of length n , the $n \times n$ lag-correlation matrix \mathbf{C} is decomposed using the SVD-CMP routine (Press et al., 1986). In the eigenspectrum, the position of each eigenvalue on the abscissa is given by the dominant frequency associated with the related EOF, i.e., detected in the related mode. That is, the studied time series is projected onto the particular EOF, the power spectrum of the projection (mode) is estimated, and the frequency bin with the highest power is identified. This spectral coordinate is mapped onto one of the n frequency bins, which equidistantly divide the abscissa of the eigenspectrum.
2. An AR(1) model is fitted to the series under study, and the residuals are computed.
3. The surrogate data are generated using the above AR(1) model, where “scrambled” (randomly permuted in temporal order) residuals are used as innovations, i.e., the noise term $\gamma z(i)$ in Eq. (6).

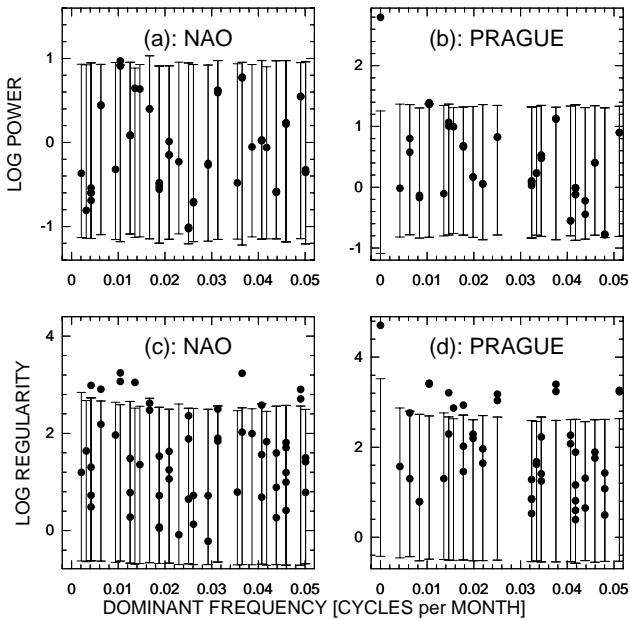


Fig. 1. Enhanced MC SSA analysis of the NAO index (a, c) and Prague near-surface air temperature series (b, d). Low-frequency parts of eigenspectra – logarithms of eigenvalues (“LOG POWER”) (a, b) and regularity index spectra (c, d). Bursts – eigenvalues or regularity indices for the analysed data; bars – 95% of the surrogate eigenvalues or regularity index distribution, i.e., the bar is drawn from the 2.5th to the 97.5th percentiles of the surrogate eigenvalues/regularity indices distribution. Both datasets span the period 1824–2002, the embedding dimension $n=480$ months was used.

4. Each realization of the surrogates undergoes SSA as described in step 1. Then, the eigenvalues for the whole surrogate set, in each frequency bin, are sorted and the values for the 2.5th and 97.5th percentiles are found. In eigenspectra, the 95% range of the surrogates’ eigenvalue distribution is illustrated by a horizontal bar between the above percentile values.
5. For each frequency bin, the eigenvalue obtained from the studied data is compared with the range of the surrogate eigenvalues. If an eigenvalue lies outside the range given by the above percentiles, the null hypothesis of the AR(1) process is rejected, i.e., there is a probability $p < 0.05$ that the data can be explained by the null noise model.
6. For each SSA mode (a projection of the data onto a particular EOF), the regularity index is computed, as well as for each SSA mode for all the realizations of surrogate data. The regularity indices are processed and statistically tested in the same way as the eigenvalues. The regularity index is based on mutual information obtained by a simple box-counting approach with marginal equiquantization (Paluš, 1995, 1996a, 1997a).

3 The data

The NAO index is traditionally defined as the normalized pressure difference between the Azores and Iceland. The NAO data used here and their description are available at <http://www.cru.uea.ac.uk/cru/data/nao.htm>.

In the initial stage of this study, we used monthly average near-surface air temperature time series from ten European stations (see Paluš and Novotná, 2004, for details) obtained from the Carbon Dioxide Information Analysis Center Internet server (<ftp://cdiac.esd.ornl.gov/pub/ndp041>) and a time series from the Prague–Klementinum station from the period 1781–2002. The long-term monthly averages were subtracted from the data, so that the annual cycle was effectively filtered out. In this paper, which should be considered as a methodological one, i.e., introducing the method of detection and extraction of oscillatory modes from raw data and their subsequent synchronization analysis, we focus on the latter data from the Prague-Klementinum station, taken from the period 1824–2002 in order to have the same time span as the used NAO index data. Further results from other stations as well as from NCEP/NCAR reanalysis series will be published elsewhere.

4 Enhanced MC SSA: detection and extraction of oscillatory modes

The results from the enhanced MC SSA for the considered NAO index and the Prague near-surface air temperature time series are presented in Fig. 1. In order to have the results comparable with previous studies (Paluš and Novotná, 2004; Gámiz-Fortis et al., 2002), we used the embedding dimension $n=480$ months. In the standard MC SSA, the only eigenvalue undoubtedly distinct from the surrogate range is the trend (zero frequency) mode in the temperature (Fig. 1b). Further, there are two modes at the frequency 0.0104 just above the surrogate bar in both the temperature and NAO test (Figs. 1a, b). The eigenvalue for the temperature is slightly closer to significance than in Paluš and Novotná (2004). This better distinction is due to the larger number of surrogate realizations (5000) used in this study. These results, however, are still “on the edge” of significance and are not very convincing.

A quite different picture is obtained from the analyses based on the regularity index (Figs. 1c, d). In Paluš and Novotná (2004), slow modes were the main interest, so that in the estimation of the regularity index (7), the parameter τ_{\max} was extended up to four decades. We expected that such a choice could diminish the test sensitivity for faster oscillations and therefore in the present study we set $\tau_{\max}=120$ months. This strongly increased the significance of the period 8 yr modes (frequency 0.0104), and in addition several other new significant modes appeared in both the NAO and temperature (Figs. 1c, d). The distinction of the regularity

indices of these modes from the related surrogate ranges is clear and even the simultaneous statistical inference (see Paluš, 1995, and references therein) cannot jeopardize the significance of the results. The significant modes in the NAO are located at the frequencies (in cycles per month) 0.004, 0.006, 0.0104, 0.014, 0.037 and 0.049, corresponding to the periods of 240, 160, 96, 73, 27 and 20 months, respectively. Besides the zero frequency (trend) mode, the significant modes in the temperature are located at the frequencies 0.0104, 0.014, 0.016, 0.018, 0.025, 0.037 and 0.051, corresponding to the periods of 96, 68, 64, 56, 40, 27 and 20 months, respectively. The modes with period 8 years were extracted and presented in Paluš and Novotná (2004), their mean frequency was estimated with higher precision as 7.8 years. Besides the latter modes (and the trend mode in the temperature), the highest regularity index was obtained for the modes with the period of 27 months (frequency 0.037). This frequency lies within the range of the quasi-biennial oscillations (QBO), therefore we will use the term “QBO modes”. The QBO modes were extracted from both the NAO index and the temperature, i.e., the raw data were projected on the related eigenvector, and are illustrated in Figs. 2a, b. The extracted modes are shorter than the original time series by the embedding dimension $n=480$ months, and there is an uncertainty of the exact timing of the modes equal to the embedding window of 40 years. We adjusted the temporal coordinate of the modes by maximizing their correlation with the original data, although this approach does not always give an unambiguous result.

5 Synchronization analysis: the method and results

Based on the concept of phase synchronization of chaotic oscillators (Rosenblum et al., 1996; Pikovsky et al., 2001), a novel technique has been developed to analyze complex, even non-stationary, bivariate data (Pikovsky et al., 2001; Paluš, 1997b). First, we calculate the instantaneous phases $\phi_1(t)$ and $\phi_2(t)$ of analyzed signals, in this case, of the QBO modes extracted from the NAO and the temperature data. The instantaneous phase and amplitude of a signal $s(t)$ can be determined by using the analytic signal concept of Gabor (1946), recently introduced into the field of nonlinear dynamics within the context of chaotic synchronization (Rosenblum et al., 1996). The analytic signal $\psi(t)$ is a complex function of time defined as

$$\psi(t) = s(t) + j\hat{s}(t) = A(t)e^{j\phi(t)}. \tag{8}$$

Usually, the imaginary part $\hat{s}(t)$ of the analytic signal $\psi(t)$ can be obtained by using the Hilbert transform of $s(t)$

$$\hat{s}(t) = \frac{1}{\pi} \text{P.V.} \int_{-\infty}^{\infty} \frac{s(\tau)}{t - \tau} d\tau. \tag{9}$$

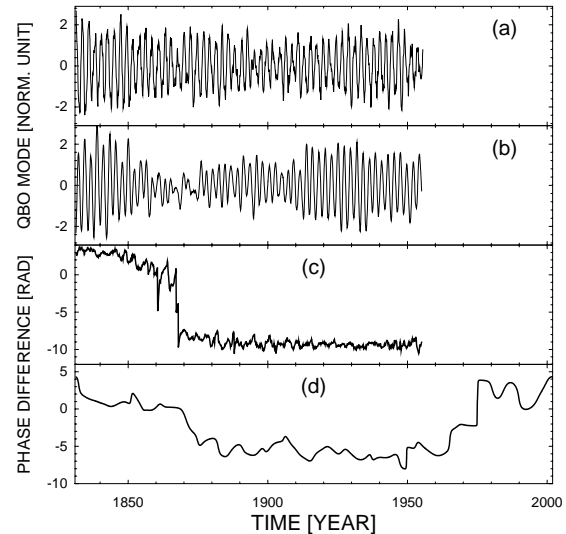


Fig. 2. The QBO mode extracted from the monthly NAO index (a) and the monthly Prague near-surface air temperature series (b). The instantaneous phase difference of the NAO and the temperature QBO modes obtained from the instantaneous phases computed from the SSA modes (c) and the same phase difference computed from the instantaneous phases estimated using the complex continuous wavelet transform (d).

(P.V. means that the integral is taken in the sense of the Cauchy principal value.) $A(t)$ is the instantaneous amplitude and the instantaneous phase $\phi(t)$ of the signal $s(t)$ is

$$\phi(t) = \arctan \frac{\hat{s}(t)}{s(t)}. \tag{10}$$

Having the analyzed signals extracted as the SSA modes, each oscillatory mode usually exists together with its orthogonal ($\pi/2$ -delayed or advanced) version. These two modes can be considered to be the real and imaginary parts of the analytic signal and the phase $\phi(t)$ can be obtained according to Eq. (10). Another approach, also considered here, is based on the wavelet transform (Torrence and Compo, 1998). Applying a continuous complex wavelet transform directly to the NAO index and the temperature time series, the complex coefficients related to the scale (frequency) of the studied cycles (the period of 27 months) can directly be used in Eq. (10) to estimate the phase $\phi(t)$. Having the instantaneous phases $\phi_1(t)$ and $\phi_2(t)$ of the NAO and the temperature QBO modes, respectively, we define the instantaneous phase difference

$$\Delta\phi(t) = \phi_1(t) - \phi_2(t). \tag{11}$$

In the classical case of periodic self-sustained oscillators, phase synchronization is defined as phase locking, i.e., the phase difference is constant. In the case of phase-synchronized chaotic or other complex noisy systems, fluctuations of phase difference typically occur. Therefore, the criterion for phase synchronization is that the absolute values of $\Delta\phi$ must be bounded (Rosenblum et al., 1996). When

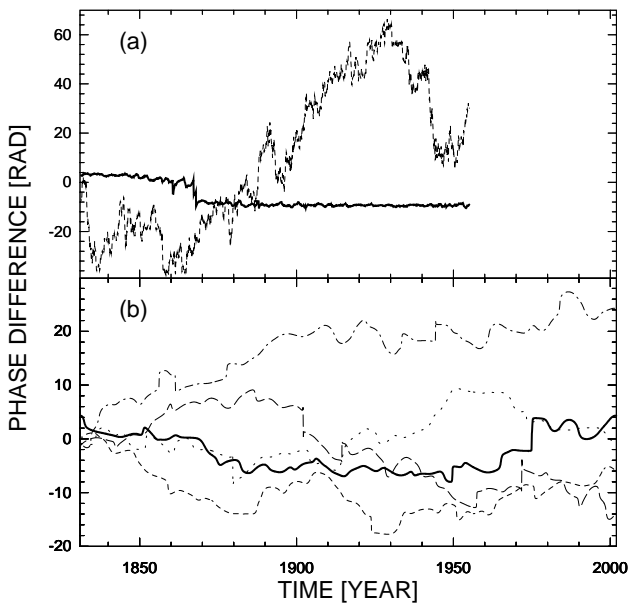


Fig. 3. (a) The instantaneous phase difference of the NAO and the temperature QBO modes obtained from the instantaneous phases computed from the SSA modes (thick solid line) compared with the instantaneous phase difference of the oscillatory modes of frequency 0.0407 (thin dashed line), again extracted by SSA from the monthly NAO index and the monthly Prague near-surface air temperature series. (b) The instantaneous phase difference of the NAO and the temperature QBO modes computed from the instantaneous phases estimated using the complex continuous wavelet transform (thick solid line) compared with the instantaneous phase differences computed in the same way, but with realization of the isospectral surrogate data used instead of the temperature time series (thin dashed/dotted lines).

the instantaneous phases are not represented as cyclic functions in the interval $[0, 2\pi]$ but as monotonously increasing functions on the whole real line, then the instantaneous phase difference $\Delta\phi(t)$ is also defined on the real line and is an unbounded (increasing or decreasing) function of time for asynchronous states of systems, while epochs of phase synchronization appear as plateaus in the $\Delta\phi(t)$ vs. time plots. The phase difference $\Delta\phi(t)$ of the NAO and the temperature QBO modes obtained from the instantaneous phases computed from the SSA modes as the real and the imaginary parts of the analytic signals is presented in Fig. 2c. After an initial decrease and a jump, there is a clear plateau starting at about the year 1870, extending to 1955, when the SSA estimate of the phases ends due to the embedding dimension $n=480$ months. A larger part of the data can be exploited using the wavelet transform, which gives the phase difference $\Delta\phi(t)$ of the NAO and the temperature QBO modes presented in Fig. 2d. Paluš et al. (2005) compared the three above mentioned methods for the estimation of instantaneous phase and demonstrated their equivalency. There are, however, some differences in the numerical properties of the es-

timates. SSA uses the natural basis of the empirical orthogonal functions and thus is more precise and sensitive to local extrema. Its drawback is the above mentioned uncertainty of the temporal localization of the modes and the phases. The cross-correlation method for the temporal localization does not always give unambiguous results and can hardly be fully automatized for the processing of large amounts of data. The wavelet method caused some smoothing of the estimates given by the shape of the wavelet basis functions leading to decreased precision of the estimates, however, the estimated phases are exactly localized in time and thus automated processing of large amounts of data is possible. Nevertheless, in the period 1870–1955, both methods consistently estimate the plateau in the $\Delta\phi(t)$ vs. time plot, i.e. they indicate phase synchronization between the QBO modes in the NAO and in the temperature. The flatness of the plateau in the $\Delta\phi(t)$ vs. time plot is indeed visually convincing, demonstrating the bounded character of $\Delta\phi(t)$. One can doubt, however, whether this is indeed a demonstration of the physical phenomenon of phase synchronization. The extraction of the SSA modes is in fact a very narrow band filtration of the analyzed signal. Two sine waves with the same frequency would be formally phase-locked without any physical interaction. In order to demonstrate that the plateau in Figs. 1c, d is not a trivial consequence of the filtration, we compute the instantaneous phase difference for two SSA modes of close frequency. The modes on the frequency 0.0407 in both the NAO and the temperature series (Fig. 1) are not significant in our enhanced MC SSA test and thus do not represent any physically existing oscillation but just the narrow-band filtration of the background noise. Their phase difference (Fig. 3a, thin dashed line) behaves as an unbounded random walk.

We can go further and compute the instantaneous phase difference of oscillatory modes on the same frequency (period 27 months) using the wavelet method applied to the NAO index, but instead of the temperature data we use so-called isospectral surrogate data (Theiler et al., 1992; Paluš, 1995). In the above MC SSA, we used the AR(1) model as the surrogate data which preserved the $1/f$ character of the spectrum of the data but not possible cycles. In this case, we need to preserve the possible existence of cycles, i.e., we need to preserve the whole spectrum of the signal, but to randomize phases of such cycles in order to destroy any dependence on the NAO possibly contained in the temperature data. The isospectral surrogate data are constructed from the real temperature data by means of the Fast Fourier Transform (FFT): The FFT is computed, the magnitude of the FFT coefficients (i.e. the power spectrum) is preserved, but the Fourier phases are randomized. After inverse FFT, we obtain a signal with the same power spectrum as the temperature series, but independent of the NAO or any other real atmospheric process.

The instantaneous phase differences $\Delta\phi(t)$ obtained using four different (independently randomized) realizations of the isospectral surrogate data are presented in Fig. 3b as thin

dashed and/or dotted lines. They are apparently unbounded, they look smoother due to their estimation by means of the wavelet transform.

We will not rely on the above visual evidence but will prove the existence of phase synchronization in a statistical way. First, we express the flatness of the plateau in the $\Delta\phi(t)$ vs. time plot in a quantitative way. Now, we consider $\Delta\phi(t)$ as $\Delta\phi(t) \bmod(2\pi)$ and use a simple trigonometric statistic, known also as the mean resultant length (MRL) (Allefeld and Kurths, 2004):

$$\text{MRL} = \sqrt{\langle \cos(\Delta\phi(t)) \rangle^2 + \langle \sin(\Delta\phi(t)) \rangle^2} \quad (12)$$

where $\langle \rangle$ means the temporal average. The MRL tends to zero for asynchronous processes and to one for phase locked systems. Considering real noisy data, neither 0 nor 1 is reached, but we can expect that the MRL for synchronized oscillations will be significantly larger than for unsynchronized (surrogate) processes.

Paluš (1997b) demonstrated that the instantaneous phases ϕ_1, ϕ_2 of phase-synchronized systems are confined in strip-like structures when plotted in the plane $(0, 2\pi) \times (0, 2\pi)$, while the instantaneous phases ϕ_1, ϕ_2 of asynchronous processes almost homogeneously fill the plane $(0, 2\pi) \times (0, 2\pi)$. The dependence between ϕ_1 and ϕ_2 , i.e., the structure in the ϕ_1, ϕ_2 plot, can be quantified by using various statistical or information-theoretic approaches. Here we use the mutual information (Cover and Thomas, 1991), defined as

$$I(\phi_1, \phi_2) = \int_0^{2\pi} \int_0^{2\pi} p_{1,2}(\phi_1, \phi_2) \log \frac{p_{1,2}(\phi_1, \phi_2)}{p_1(\phi_1)p_2(\phi_2)} d\phi_1 d\phi_2, \quad (13)$$

where $p_1(\phi_1)$ and $p_2(\phi_2)$ are probability distributions of the phases ϕ_1 and ϕ_2 , respectively, and $p_{1,2}(\phi_1, \phi_2)$ is their joint distribution. Theoretically, independence of the phases, i.e., the absence of phase synchronization, means homogeneous distribution of the points (ϕ_1, ϕ_2) and $I(\phi_1, \phi_2)=0$; while for phase synchronization, i.e., a mutual dependence of the phases, $I(\phi_1, \phi_2)>0$ holds. For reliable detection of the phase synchronization in experimental data, it is necessary to establish that $I(\phi_1, \phi_2)>0$ with a statistical significance. Therefore, we compute both the MRL and $I(\phi_1, \phi_2)$ for a large number (5000) of realizations of the surrogate data, construct histograms and cumulative histograms of the MRL and $I(\phi_1, \phi_2)$ surrogate distributions (Fig. 4) and compare them with the MRL and $I(\phi_1, \phi_2)$ values obtained for the NAO and temperature data (solid vertical lines in Fig. 4). The cumulative histograms can be used for estimation of the significance of the test – the p value which is in fact the probability, that such a value of the MRL or $I(\phi_1, \phi_2)$, as obtained from the NAO and temperature data, can be obtained by chance from unsynchronized or independent processes with the same frequency spectra. The obtained values

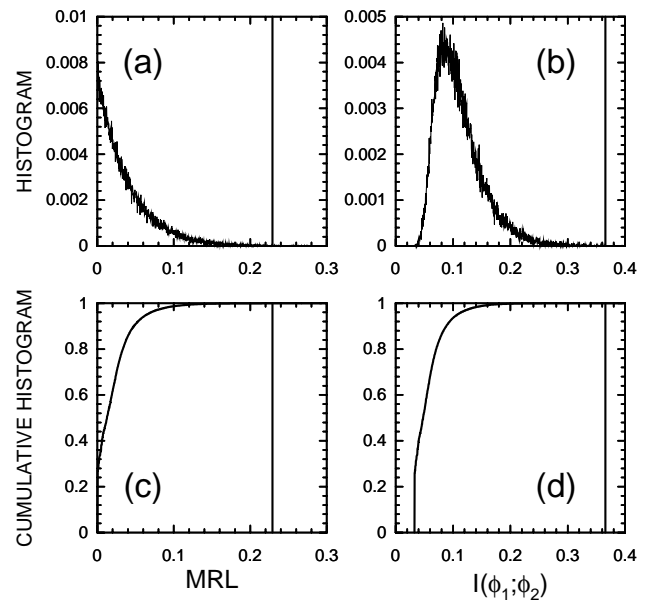


Fig. 4. Histograms (a, b) and cumulative histograms (c, d) for the surrogate distribution of the mean resultant length (a, c) and the mutual information $I(\phi_1, \phi_2)$ (b, c). The related values obtained from the QBO phases extracted from the monthly NAO index and the monthly Prague near-surface air temperature series are presented as the solid vertical lines. The wavelet transform was used for the phase estimation.

are $p < 0.0005$ for the MRL and $p < 0.00002$ for the $I(\phi_1, \phi_2)$ test. We can conclude that phase synchronization of the NAO and temperature QBO modes has been proven with a high statistical significance.

6 Discussion of results

An introduction of relatively new, nonlinear methods to analysis of experimental time series always evokes two types of questions: 1) Does the new method bring more information than common, already verified linear methods? 2) Does the the new nonlinear method indeed reflect a nonlinear phenomenon or is a linear model/explanation sufficient? In this case, cannot a simple cross-correlation mimic the phenomenon of phase synchronization? Searching for answers to both questions, we perform the same testing as above (Fig. 4), but instead of (independent) surrogate data which preserved spectral properties of individual time series, now we construct bivariate surrogate data (Prichard and Theiler, 1994; Paluš, 1996b) preserving also the cross-correlations of the original (NAO and temperature) time series. Looking at the resulting histograms (Fig. 5) we can see that bivariate FFT surrogate data can reject the null hypothesis of independence used in the previous test (Fig. 4) and, indeed, even cross-correlated linear processes could be erroneously considered as phase-synchronized nonlinear oscillations. In

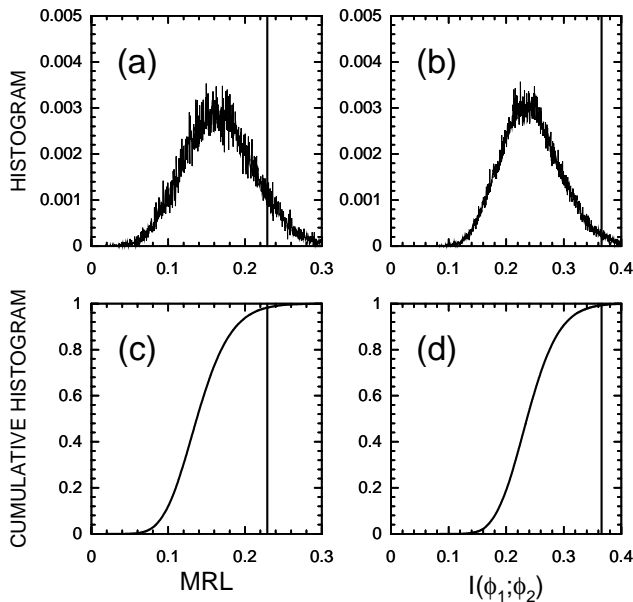


Fig. 5. Histograms (a, b) and cumulative histograms (c, d) for the distribution of the mean resultant length (a, c) and the mutual information $I(\phi_1, \phi_2)$ (b, c) obtained from the bivariate surrogate data preserving the crosscorrelation between the NAO and temperature series. The related values obtained from the QBO phases extracted from the monthly NAO index and the monthly Prague near-surface air temperature series are presented as the solid vertical lines. The wavelet transform was used for the phase estimation.

our case, however, the values of the mean resultant length and the mutual information $I(\phi_1, \phi_2)$ show that the association of the phases of the QBO modes in the NAO and temperature data is still significantly stronger than that obtained from bivariate surrogate data. The obtained significance values are $p < 0.02$ for the MRL and $p < 0.008$ for the $I(\phi_1, \phi_2)$ test (Fig. 5). So we can conclude that we have detected a nonlinear interaction of the QBO modes in the NAO and temperature records which neither can be explained by a linear model, nor can be sufficiently described by linear measures. On the other hand, the question whether we have detected the phase synchronization *sensu stricto* is more subtle. Phase synchronization is a phenomenon occurring in an interaction of two autonomous (possibly chaotic and/or noisy) oscillatory processes. The atmospheric temperature and pressure (giving the NAO index) are two variables of one large nonlinear spatio-temporal system – the atmosphere. It is a question of the level of description and/or approximation whether some modes of atmospheric variability can or cannot be considered as autonomous processes interacting through a nonlinear and time variable coupling. It is known that the relationship between the NAO and the NHT is time-variable (Polyakova et al., 2006) and can be influenced by solar (Gimeno et al., 2003) and geomagnetic activity (Bochníček and Hejda, 2005). The long-term dynamics reflected in long-

range correlations or long-term memory have different properties in the NAO and in the temperature records (Fraedrich, 2003). Recently, Mokhov and Smirnov (2006) have demonstrated that the NAO interacts, or is influenced by the other global atmospheric oscillatory process – the El Niño Southern Oscillation. We believe that phase synchronization can be at least a good working model and the above utilized tools can be developed in the context of phase synchronization to measure the association of phases. We also believe that measures characterizing the coupling direction (causality) in the interaction of phases (Rosenblum and Pikovsky, 2001; Paluš and Stefanovska, 2003) can help in the understanding of the interaction of modes of atmospheric variability (Mokhov and Smirnov, 2006). This does not mean that, in particular studies, the new, nonlinear methods should not be accompanied by standard approaches and their extensions, especially by those that can cope with nonstationarity, e.g. cross-wavelet or wavelet coherence approaches (Grinsted et al., 2004; Maraun and Kurths, 2004).

7 Conclusions

In this paper, we demonstrate a sophisticated combination of methods from nonlinear dynamics with an established linear method in an application to multivariate geophysical data. Monte Carlo Singular System Analysis has been extended by evaluating and testing the regularity of the dynamics of the SSA modes against the colored noise null hypothesis in addition to the test based on variance (eigenvalues). The nonlinear approach to the measurement of regularity and predictability of the dynamics, based on a coarse-grained estimate of the mutual information, increases the MC SSA test sensitivity and reliability in the detection of dynamical modes which are relatively more regular than those obtained by decomposition of colored noise.

Enhanced MC SSA has been applied to records of monthly average near-surface air temperature from several European locations as well as to the monthly NAO index and several significant oscillatory modes have been detected by testing the regularity of the modes. Then, we focused on the oscillatory modes with the average period 27 months which can be considered as realizations of the quasi-biennial oscillations (QBO) in both the NAO index and temperature data. Another method from nonlinear dynamics, synchronization analysis (Pikovsky et al., 2001; Paluš, 1997b) gives the possibility to establish existence of coupling between the QBO modes present in the NAO and in the temperature time series. Applying the concept of surrogate data, we proved with high statistical significance that the QBO modes in the NAO and in the temperature time series are phase synchronized. We also showed that the phenomenon of phase synchronization is probably a real physical phenomenon and its signatures do not occur in phase relations of formally extracted, non significant SSA modes, although they represent a very narrow-

band filtered colored noise which could erroneously be considered as real oscillations. The statistical evidence for the phase synchronization between the QBO modes can be considered as additional evidence for the real existence of these oscillatory modes in the dynamics of NAO and air temperature and thus, in return, it confirms the validity of our enhanced MC SSA test.

The introduced set of methods will be applied to temperature data from other stations, as well as to other meteorological measurements and to the NCEP/NCAR reanalysis series. We believe that the results will help us to understand how global atmospheric circulation processes influence European weather and climate changes on various temporal scales.

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References

- Allefeld, C. and Kurths, J.: Testing for phase synchronization, *Int. J. Bif. Chaos*, 14(2), 406–416, 2004.
- Allen, M. R. and Smith, L. A.: Investigating the origins and significance of low-frequency modes of climate variability, *Geophys. Res. Lett.* 21, 883–886, 1994.
- Allen, M. R. and Smith, L. A.: Monte Carlo SSA: Detecting irregular oscillations in the presence of colored noise, *J. Climate*, 9(12), 3373–3404, 1996.
- Bochníček, J. and Hejda, P.: The winter NAO pattern changes in association with solar and geomagnetic activity, *J. Atmos. Solar-Terrestrial Phys.*, 67, 17–32, 2005.
- Bojariu, R. and Gimeno, L.: Predictability and numerical modelling of the North Atlantic Oscillation, *Earth-Sci. Rev.*, 63(1–2), 145–168, 2003.
- Cover, T. M. and Thomas, J. A.: *Elements of Information Theory*, J. Wiley & Sons, New York, 1991.
- Eden, C., Greatbatch, R. J., and Lu, J.: Prospects for decadal prediction of the North Atlantic Oscillation (NAO), *Geophys. Res. Lett.*, 29(10), 1466–1469, 2002.
- Fernandez, I., Hernandez, C. N., and Pacheco, J. M.: Is the North Atlantic Oscillation just a pink noise? *Physica A*, 323, 705–714, 2003.
- Fraedrich, K.: Short- and long term memory of the atmosphere, in: *Chaos in Geophysical Flows*, edited by: Boffetta, G., Larcotta, S., Visconti G., and Vulpiani, A., Otto Editore, Torino, 63–104, 2003.
- Gabor, D.: *Theory of Communication*, J. IEE London, 93, 429–457, 1946.
- Gámiz-Fortis, S. R., Pozo-Vázquez, D., Esteban-Parra, M. J., and Castro-Díez, Y.: Spectral characteristics and predictability of the NAO assessed through Singular Spectral Analysis, *J. Geophys. Res.*, 107(D23), 4685, 2002.
- Gao, J. B., Cao, Y., and Lee, J.-M.: Principal component analysis of $1/f^\alpha$ noise, *Phys. Lett. A*, 314, 392–400, 2003.
- Ghil, M. and Vautard, R.: Interdecadal oscillations and the warming trend in global temperature time series, *Nature*, 350(6316), 324–327, 1991.
- Gimeno, L., de la Torre, L., Nieto, R., García, R., Hernández, E., and Ribera, P.: Changes in the relationship NAO–Northern hemisphere temperature due to solar activity, *Earth Planet. Sci. Lett.*, 206, 15–20, 2003.
- Grinsted A., Moore, J. C., and Jevrejeva, S.: Application of the cross wavelet transform and wavelet coherence to geophysical time series, *Nonlin. Processes Geophys.*, 11, 561–566, 2004, <http://www.nonlin-processes-geophys.net/11/561/2004/>.
- Hurrell, J. W., Kushnir, Y., and Visbeck, M.: Climate – The North Atlantic oscillation, *Science*, 291(5504), 603, 2001.
- Keppenne, C. L. and Ghil, M.: Adaptive filtering and the Southern Oscillation Index, *J. Geophys. Res.*, 97, 20 449–20 454, 1992.
- Maraun, D. and Kurths, J.: Cross wavelet analysis: significance testing and pitfalls, *Nonlin. Processes Geophys.*, 11, 505–514, 2004, <http://www.nonlin-processes-geophys.net/11/505/2004/>.
- Mokhov, I. I. and Smirnov, D. A.: El Niño–Southern Oscillation drives North Atlantic Oscillation as revealed with nonlinear techniques from climatic indices, *Geophys. Res. Lett.*, 33, L03708, doi:10.1029/2005GL024557, 2006.
- Paluš, M.: Testing for nonlinearity using redundancies: Quantitative and qualitative aspects, *Physica D*, 80, 186–205, 1995.
- Paluš, M.: Coarse-grained entropy rates for characterization of complex time series, *Physica D*, 93, 64–77, 1996a.
- Paluš, M.: Detecting nonlinearity in multivariate time series, *Phys. Lett. A*, 213, 138–147, 1996b.
- Paluš, M.: Kolmogorov entropy from time series using information-theoretic functionals, *Neural Network World*, (<http://www.uivt.cas.cz/~mp/papers/rd1a.ps>), 3/97, 269–292, 1997a.
- Paluš, M.: Detecting phase synchronization in noisy systems, *Phys. Lett. A*, 235, 341–351, 1997b.
- Paluš, M. and Novotná, D.: Detecting modes with nontrivial dynamics embedded in colored noise: Enhanced Monte Carlo SSA and the case of climate oscillations, *Phys. Lett. A*, 248, 191–202, 1998.
- Paluš, M. and Stefanovska, A.: Direction of coupling from phases of interacting oscillators: An information-theoretic approach, *Phys. Rev. E*, 67, 055201(R), 2003.
- Paluš, M. and Novotná, D.: Enhanced Monte Carlo Singular System Analysis and detection of period 7.8 years oscillatory modes in the monthly NAO index and temperature records, *Nonlin. Processes Geophys.*, 11, 721–729, 2004, <http://www.nonlin-processes-geophys.net/11/721/2004/>.
- Paluš, M., Novotná, D., and Tichavský, P.: Shifts of seasons at the European mid-latitudes: Natural fluctuations correlated with the North Atlantic Oscillation, *Geophys. Res. Lett.*, 32, L12805, doi:10.1029/2005GL022838, 2005.
- Pikovsky, A., Rosenblum, M., and Kurths, J.: *Synchronization. A Universal Concept in Nonlinear Sciences*, Cambridge University Press, 2001.
- Polyakova, E. I., Journel, A. G., Polyakov, I. V., and Bhatt, U. S.: Changing relationship between the North Atlantic Oscillation and key North Atlantic climate parameters, *Geophys. Res. Lett.*, 33, L03711, doi:10.1029/2005GL024573, 2006.

- Pozo-Vázquez, D., Esteban-Parra, M. J., Rodrigo, F. S., and Castro-Díez, Y.: A study of NAO variability and its possible nonlinear influence on European surface temperature, *Clim. Dyn.*, 17, 701–715, 2001.
- Press, W. H., Flannery, B. P., Teukolsky, S. A., and Vetterling, W. T.: *Numerical Recipes: The Art of Scientific Computing*. Cambridge Univ. Press, Cambridge, 1986.
- Prichard, D. and Theiler, J.: Generating surrogate data with several simultaneously measured variables, *Phys. Rev. Lett.*, 73, 951–954, 1994.
- Rosenblum, M. G., Pikovsky, A. S., and Kurths, J.: Phase synchronization of chaotic oscillators, *Phys. Rev. Lett.*, 76, 1804–1807, 1996.
- Rosenblum, M. G. and Pikovsky, A. S.: Detecting direction of coupling in interacting oscillators, *Phys. Rev. E*, 64, 045202(R), 2001.
- Theiler, J., Eubank, S., Longtin, A., Galdrikian, B., and Farmer, J. D.: Testing for nonlinearity in time series: the method of surrogate data, *Physica D*, 58, 77–94, 1992.
- Torrence, C. and Compo, G. P.: A practical guide to wavelet analysis, *Bull. Amer. Meteorol. Soc.*, 79(1), 61–78, 1998.
- Vautard, R. and Ghil, M.: Singular spectrum analysis in nonlinear dynamics, with applications to paleoclimatic time series, *Physica D*, 35, 395–424, 1989.
- Vautard, R., Yiou, P., and Ghil, M.: Singular spectrum analysis: a toolkit for short noisy chaotic signals, *Physica D*, 58, 95–126, 1992.
- Yiou, P., Ghil, M., Jouyel, J., Paillard, D., and Vautard, R.: Nonlinear variability of the climatic system, from singular and power spectra of Late Quarternary records, *Clim. Dyn.*, 9, 371–389, 1994.