

# High- and mid-latitude quasi-2-day waves observed simultaneously by four meteor radars during summer 2000

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**Abstract.** Results from the analysis of MLT wind measurements at Dixon (73.5° N, 80° E), Esrange (68° N, 21° E), Castle Eaton (UK) (53° N, 2° W), and Obninsk (55° N, 37° E) during summer 2000 are presented in this paper. Using S-transform or wavelet analysis, quasi-two-day waves (QTDWs) are shown to appear simultaneously at high- and mid-latitudes and reveal themselves as several bursts of wave activity. At first this activity is preceded by a 51–53 h wave with S=3 observed mainly at mid-latitudes. After a short recess (or quiet time interval for about 10 days near day 205), we observe a regular sequence of three bursts, the strongest of them corresponding to a QTDW with a period of 47–48 h and S=4 at mid-altitudes.

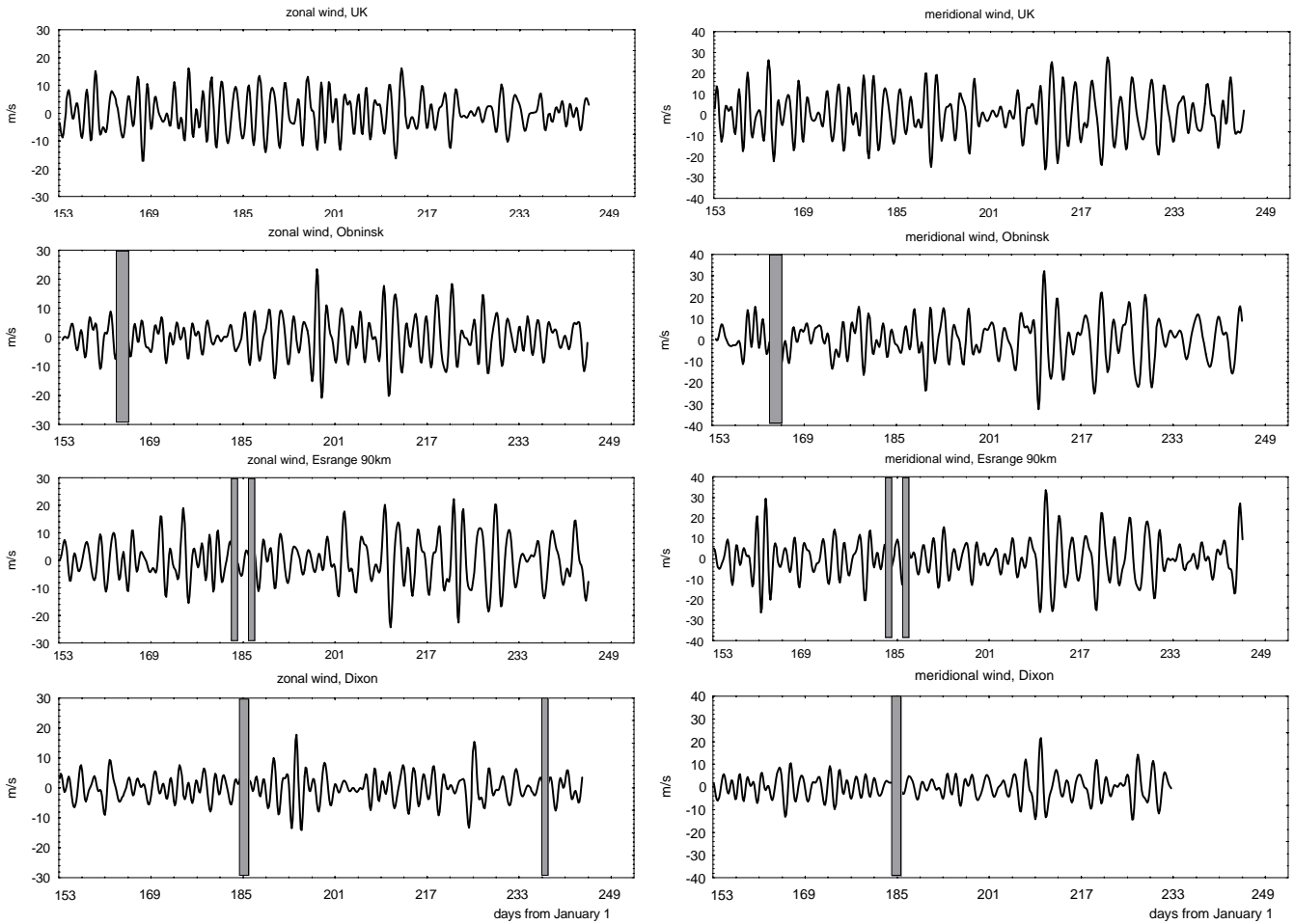
We hypothesize that these three bursts may be the result of constructive and destructive interference between several spectral components: a 47–48 h component with S=4; a 60-h component with S=3; and a 80-h component with S=2. The magnitudes of the lower (higher) zonal wave-number components increase (decrease) with increasing latitude. The S-transform or wavelet analysis indicates when these spectral components create the wave activity bursts and gives a range of zonal wave numbers for observed bursts from about 4 to about 2 for mid- and high-latitudes. The main spectral component at Dixon and Esrange latitudes is the 60-h oscillation with S=3. The zonal wave numbers and frequencies of the observed spectral components hint at the possible occurrence of the nonlinear interaction between the primary QTDWs and other planetary waves. Using a simple 3-D nonlinear numerical model, we attempt to simulate some of the observed features and to explain them as a consequence of the nonlinear interaction between the primary 47–48 h and the 9–10 day waves, and the resulting linear superposition of primary and secondary waves. In addition to the QTDW bursts, we also infer forcing of the 4-day wave with S=2 and the 6–7 day wave with S=1, possibly arising from nonlinear decoupling of the 60-h wave with S=3. The starting mechanism for this decoupling is the Rossby wave instability (e.g.

Baines, 1976). This result is consistent with the day-to-day wind variability during the observed QTDW events. An interesting feature of the final stage of the observed QTDW activity in summer 2000 is the occurrence of strong 4–5 day waves with S=3.

**Key words.** Meteorology and atmospheric dynamics (middle atmosphere dynamics; waves and tides; general or miscellaneous)

## 1 Introduction

There are not very many planetary-scale waves in the vicinity of the summer mesopause that have amplitudes of ~tens m/s and are regularly observed from one year to another. First of all, there are diurnal and semidiurnal tides, and secondly, there are QTDWs (quasi-two-day waves). Herein, we consider QTDWs to include those oscillations with frequencies between about  $0.5 \pm 0.1$  cycles day<sup>-1</sup> or periods of 40–60 h. Since the early 1970s (Muller, 1972; Kal'chenko and Bulgakov, 1973) these waves have been revealed in mesopause wind variations at low- and mid-latitudes. Glass et al. (1975) estimated the zonal wave number of this oscillation to be S=3. At present two main mechanisms have been proposed to explain the occurrence of the QTDW at levels of the mesosphere/lower thermosphere (MLT). The first one considers these waves as a manifestation of the Rossby-gravity (3,0) normal mode forced by the lower atmosphere (Salby, 1981). The second mechanism considers the waves as a result of the instability above the summer stratospheric westward jet as was proposed by Plumb (1983) and extended to two-dimensions by Pfister (1985). Salby and Callaghan (2001) explored the relationship between the Rossby-gravity mode and the instability. Their analysis recovered major features of the QTDW, including twice-yearly amplification around solstice through interaction with easterlies. Their calculations also revealed a wave number 4 component that exhibits less of the modal structure in comparison to wave number 3 component, and that should be more prevalent during July than during January.



**Fig. 1.** (left) Zonal and (right) meridional components of the winds observed from day 153 (1 June) to day 244 (31 August) in 2000. Data have been band-pass filtered to retain periods between 28 and 100 h. The longest gaps are shaded.

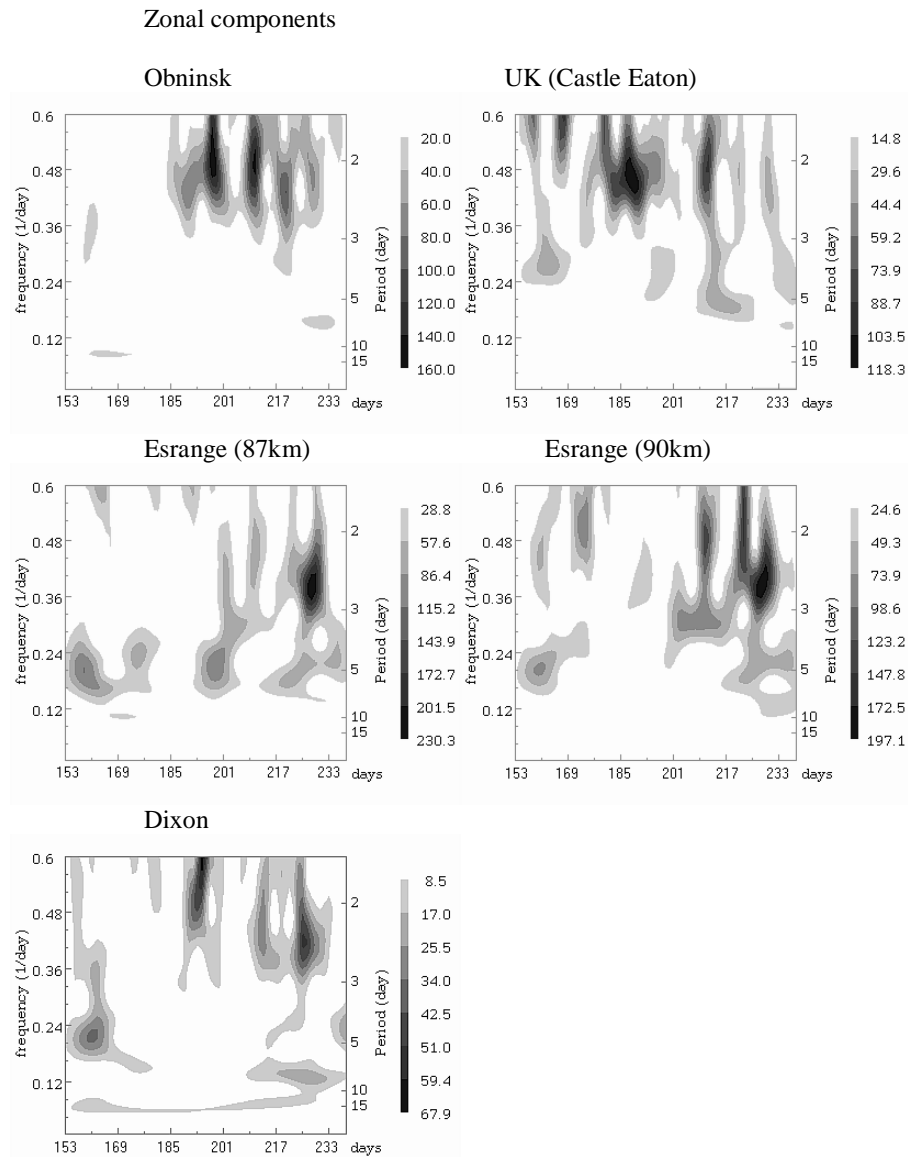
Analysis of wind data taken from several sites situated in a narrow latitudinal belt provides an opportunity to estimate the zonal wave number  $S$  of the investigated waves. Some recent results (Meek et al., 1996; Thayaparan et al., 1997; Jacobi et al., 2001) demonstrate that the range of this number is from 3 to 5 and more frequently, it is 3–4 for the Northern Hemisphere.

Numerical simulations using linear models (Hagan et al., 1993; Salby and Callaghan, 2001) and some global circulation models (Palo et al., 1999) suggest the QTDW with  $S=3$  to be the Rossby-gravity mode. However, the waves with  $S=4-5$  with slower phase speeds cannot as readily propagate from the lower to upper atmospheric levels and attain significant amplitudes at MLT heights. Simulations with global circulation models by Norton and Thuburn (1996), and Mayr et al. (2001) demonstrate that strong QTDWs with large wave numbers (4–5) can be forced in situ by baroclinic instability in the presence of strong radiative and dissipative damping and full nonlinear dynamics.

The presence of the QTDW in the high-latitude MLT region has been realized since the early investigations of the

wind circulation there (Rüster et al., 1988; van Eyken et al., 2000). However, to date there have been no estimates of the zonal wave numbers for these waves at high latitudes. The waves have been considered mainly at lower and mid-latitudes and the correspondence between the high- and mid-latitude waves has not been studied yet.

A characteristic feature of the QTDW is its amplitude modulation and the occurrence of several bursts of the wave activity. During recent years some nonlinear features of the QTDW were emphasized in relation to the amplitude modulation. For example, Jacobi et al. (1998) demonstrated regular occurrence of the nonlinear interaction between the propagating 2-day wave and other planetary waves in the upper mesosphere over Central Europe. This interaction results in an amplitude variability of the 2-day wave and forcing of secondary waves with periods near two days. Similar results were obtained by Pancheva et al. (2000), where an attempt was made to support the hypothesis of nonlinear coupling between the QTDW and the other planetary waves through bispectral analysis.



**Fig. 2.** S-transform spectra of the zonal wind variations with periods near 2 days.

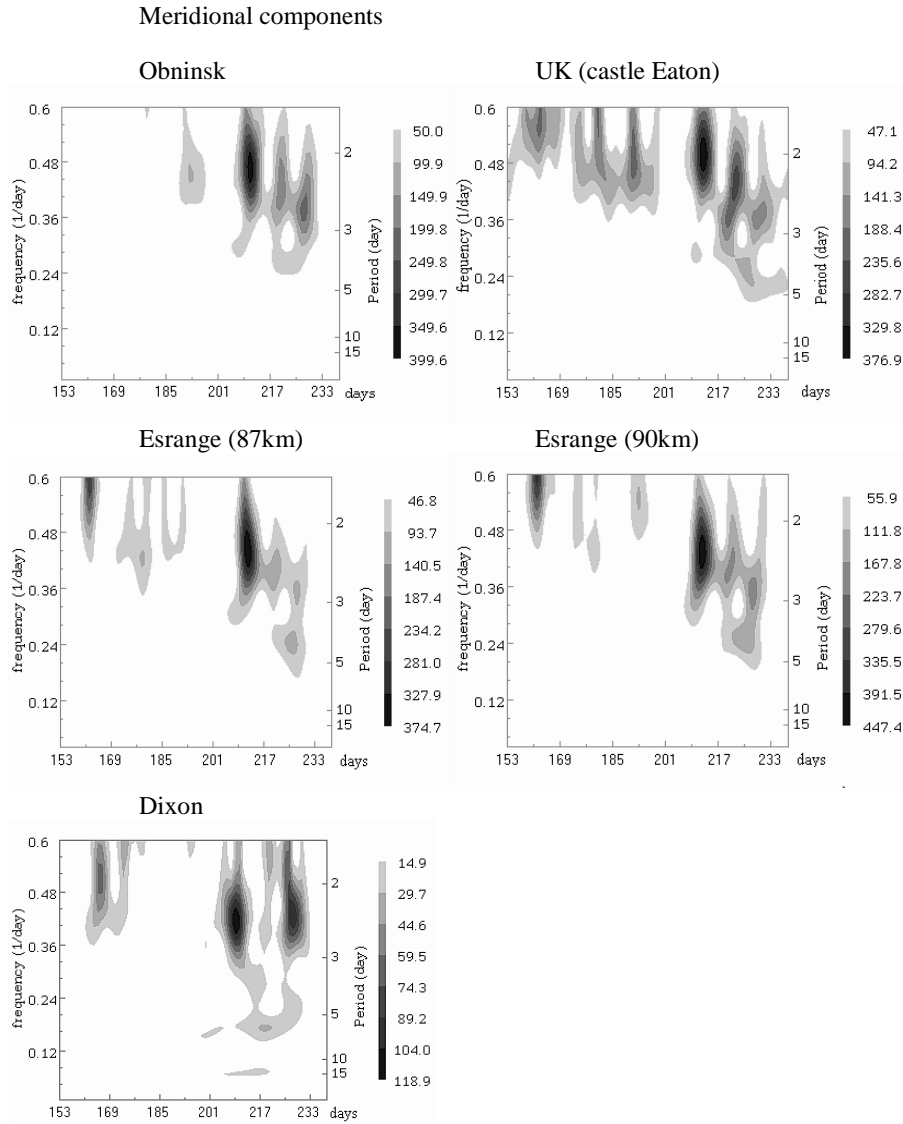
The zonal wave numbers of the secondary waves are equal to the sum and difference of the zonal wave numbers of the primary waves. The secondary wave with a lower zonal wave number is more probably observed at higher latitudes than the primary 2-day wave (just due to decreasing of wave amplitudes as  $\sin^{S-1} \theta$  when co-latitude  $\theta$  tends to 0). An analysis of wind observations during summer 2000 showed that QTDWs had periods of about 58 h at Dixon and Esrange (Middleton et al., 2001; Merzlyakov et al., 2001). This value is greater than the typical values observed at mid-latitudes and greater than the QTDW periods at the same time at Obninsk and UK. It hints at a possible occurrence of the nonlinear interaction between planetary waves during summer 2000.

The first aim of this work is to present results obtained from analysis of neutral winds measured by meteor radars at

Dixon (73.5° N, 80° E), Esrange (68° N, 21° E), Castle Eaton (UK) (53° N, 2° W), and Obninsk (55° N, 37° E) during summer 2000. We demonstrate a simultaneous manifestation of the QTDW events at high- and mid-latitudes. Secondly, we carry out an analysis of the zonal wave numbers and frequencies for the observed waves and point out a possible occurrence of the nonlinear interaction between waves, including the QTDW as well as other planetary waves. The third aim of this work is to provide some explanation of the observed features with the help of numerical simulations using a simple nonlinear 3-D model. We note the importance of the nonlinear interaction mechanism for this explanation.

## 2 Measurements and data analysis

This study is based on hourly mean wind measurements, carried out during summer 2000 at stations Dixon (73.5° N,



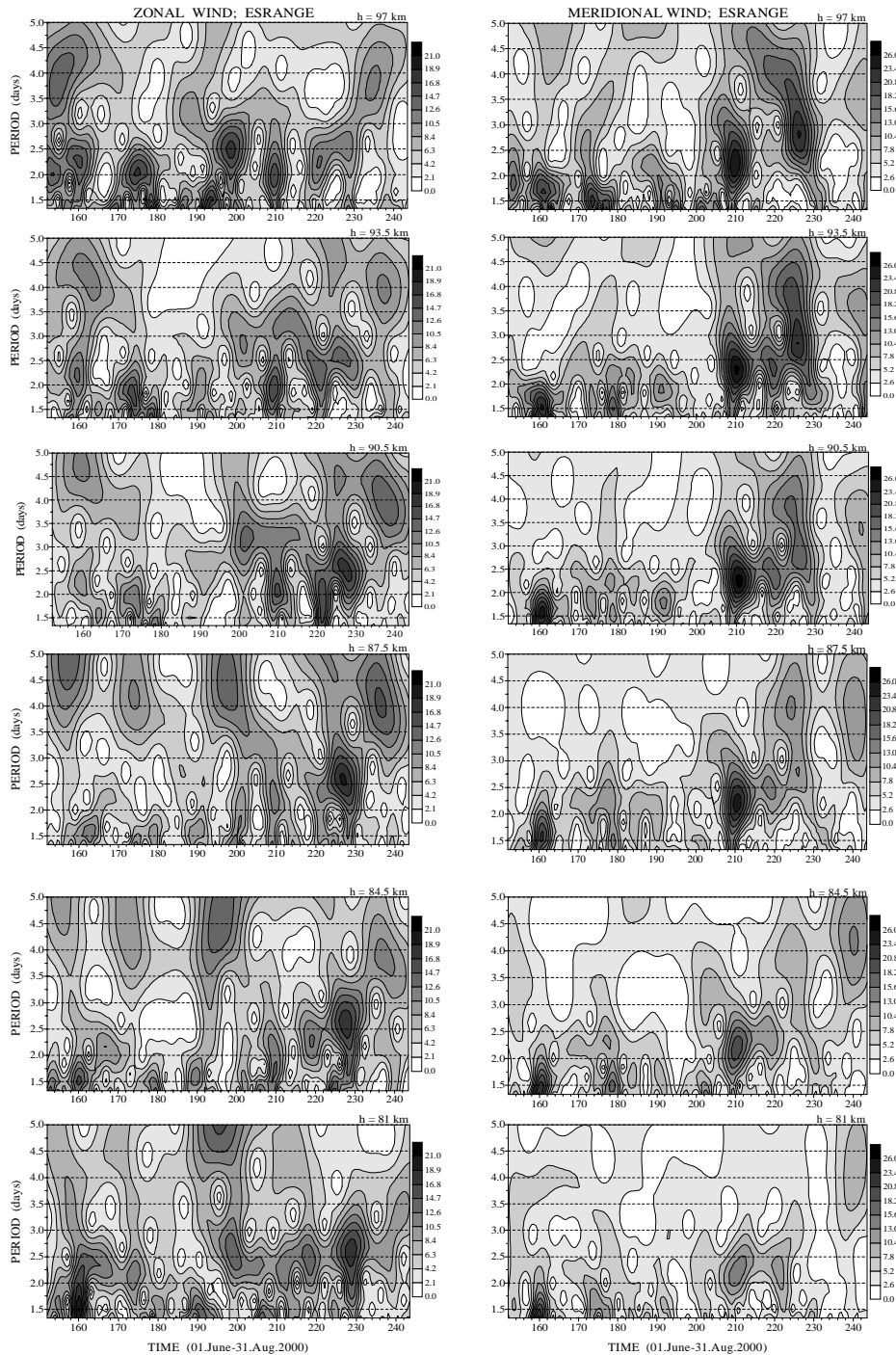
**Fig. 3.** As for Fig. 2, but for the meridional wind variations.

80° E), Esrance (68° N, 21° E), UK (53° N, 2° W) and Obninsk (55° N, 37° E). The data series cover the time interval from 1 June to 31 August. Measurements for two opposite meridional directions are available at Obninsk for this time interval. This gives two independent data realizations of the meridional wind at one site. Meteor radars are installed at all sites and the measurements at Esrance are performed with height finding. The average height of the measurements for the other three stations is estimated to be 88–90 km. The number of gaps in the Obninsk data is about 14% for June and the first half of July, and then the data are practically continuous. There is one long gap of 63 h in length for both components. Each of the UK data gaps does not cover more than several hours. The gaps are randomly distributed and constitute less than 1% of the data length. The Dixon data gaps are also of a few hours in length, randomly distributed, and form 12% of the data length. There is a long gap of 50 h in length for both components. Also, the last four days are

absent for the meridional wind data. The longest gaps of the Esrance data are an interval of 30 h in length at 90 km and an interval of 22 h in length at 87.5 km for both components in July. The rest of the data are practically continuous.

The gaps were interpolated using a least-squares fit with a polynomial of the second degree. This kind of interpolation was carefully investigated by Portnyagin et al. (1999), and it was found that it does not have any significant influence on the results.

Data processing was performed with the S-transform (Stockwell et al., 1996) and the continuous Morlet wavelet transform (e.g. Pancheva and Mukhtarov, 2000) to account for nonstationary features of the QTDWs in our analysis. Both methods showed identical spectra of wind variations, so to illustrate them we use mainly the results from the S-transform and only sometimes those from the wavelet transform.

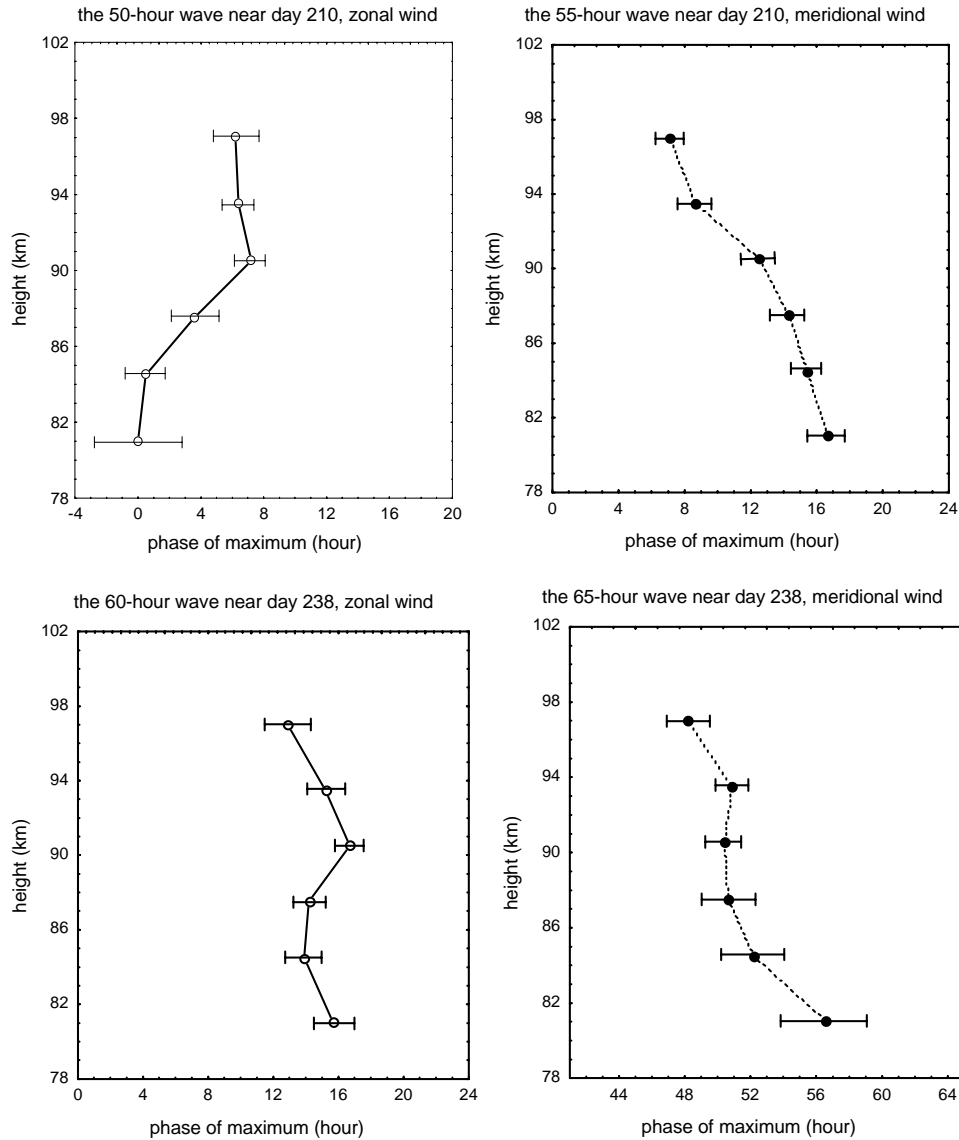


**Fig. 4.** Morlet wavelet spectra of the wind variations at Esrange.

Although we use data collected only by meteor radars, there are some differences in the results obtained by different radars. First of all, the amplitudes of oscillations observed at Esrange are greater than the values obtained at Dixon and even of the same order as amplitudes of oscillations observed at middle-latitudes. Mitchell et al. (2002) noted similar differences, except for the tides. Fortunately, the phases and frequencies of oscillations usually are very consistent for data

obtained by different techniques. All of these differences may in fact be real, and may not signal any experimental deficiency. We simply note these differences for completeness.

The day-to-day wind variations with periods greater than 4 days are considered separately because of the very large amplitude of the QTDWs. Therefore, the data were first filtered by a low-pass filter with a cutoff period of 72 h. The S-transform uses only neighboring spectral peaks of the Fast



**Fig. 5.** Phase profiles for two main bursts of the QTDW activity at Esrange.

Fourier Transform to obtain a response to a given frequency. Hence, we can treat the S-transform spectra of the day-to-day wind variations as if they had been calculated from hourly data. Then we can estimate the peak variance as Portnyagin et al. (2000), but should take for an estimation of the primary data variance only independent values of the filtered wind. Thus, we can find significance levels for every frequency in the S-transform spectra. We used a periodogram analysis to specify spectral content of the QTDWs.

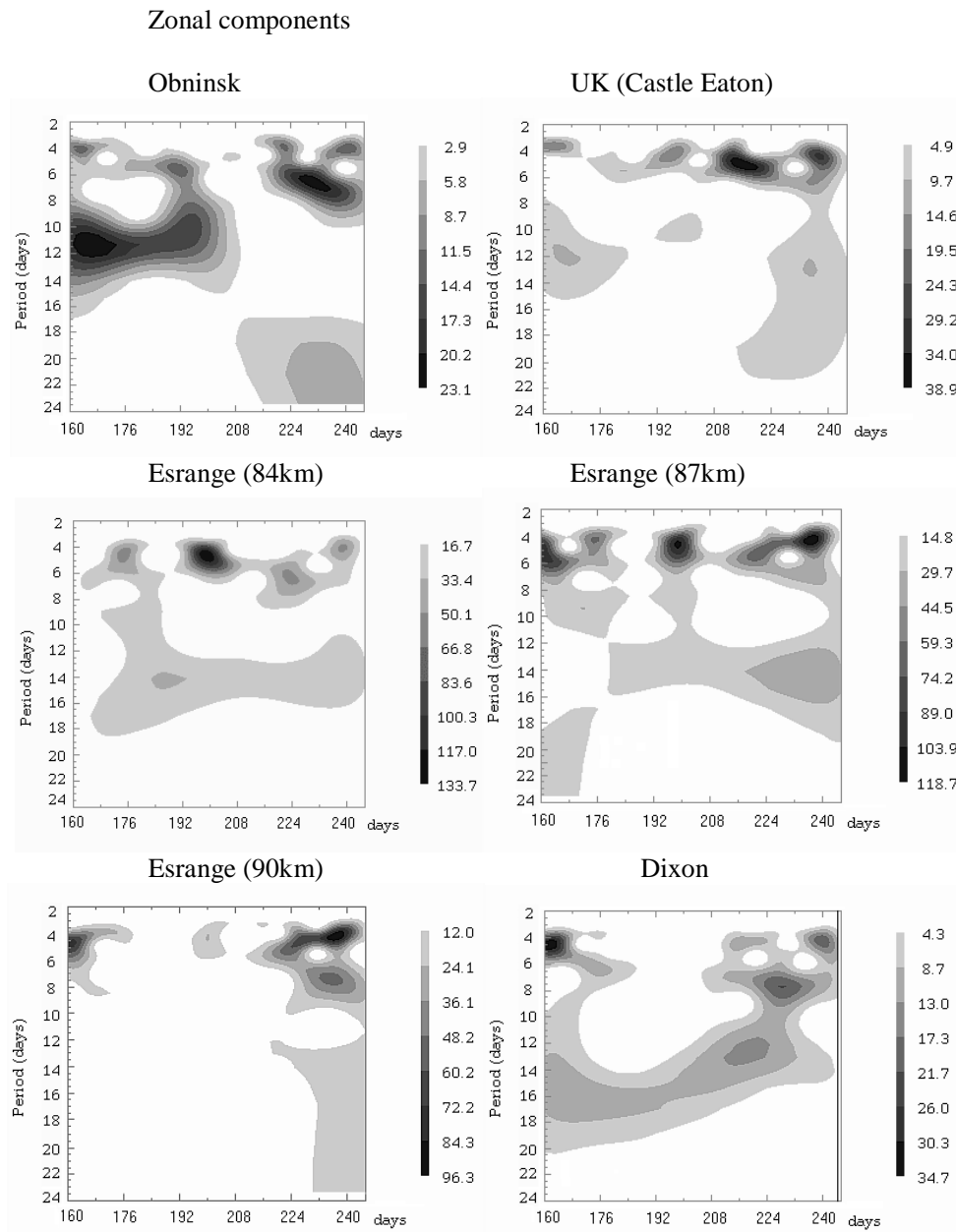
The zonal wave numbers of the oscillations were estimated from the phase values, directly obtained for each frequency from the S-transform and from a least-squares fit of the wind data with corresponding spectral components (a value of the period is taken from the wavelet spectra) and tides. The last analysis also gives values of phase errors, which were used for estimating the wave number errors. The difference

between latitudes of Dixon and Esrange (about  $5.5^\circ$ ) is neglected when comparing phases to interpret zonal wave numbers. As we shall see, such an approach is not always correct. Fortunately, latitudinal phase variations can be inferred from the numerical model, and used to assess the reasonableness of this assumption. A description of our approach to the numerical simulation is presented in the Appendix.

### 3 Results

#### 3.1 QTDW bursts: wavelet analysis

To illustrate a general look at the QTDWs in summer 2000, the data were subjected to a band-pass filter with cutoff periods of 28 h and 100 h. Figure 1 shows the resulting filtered winds. Days are numbered from 1 January. One can see a typical picture of a QTDW manifestation as a set of



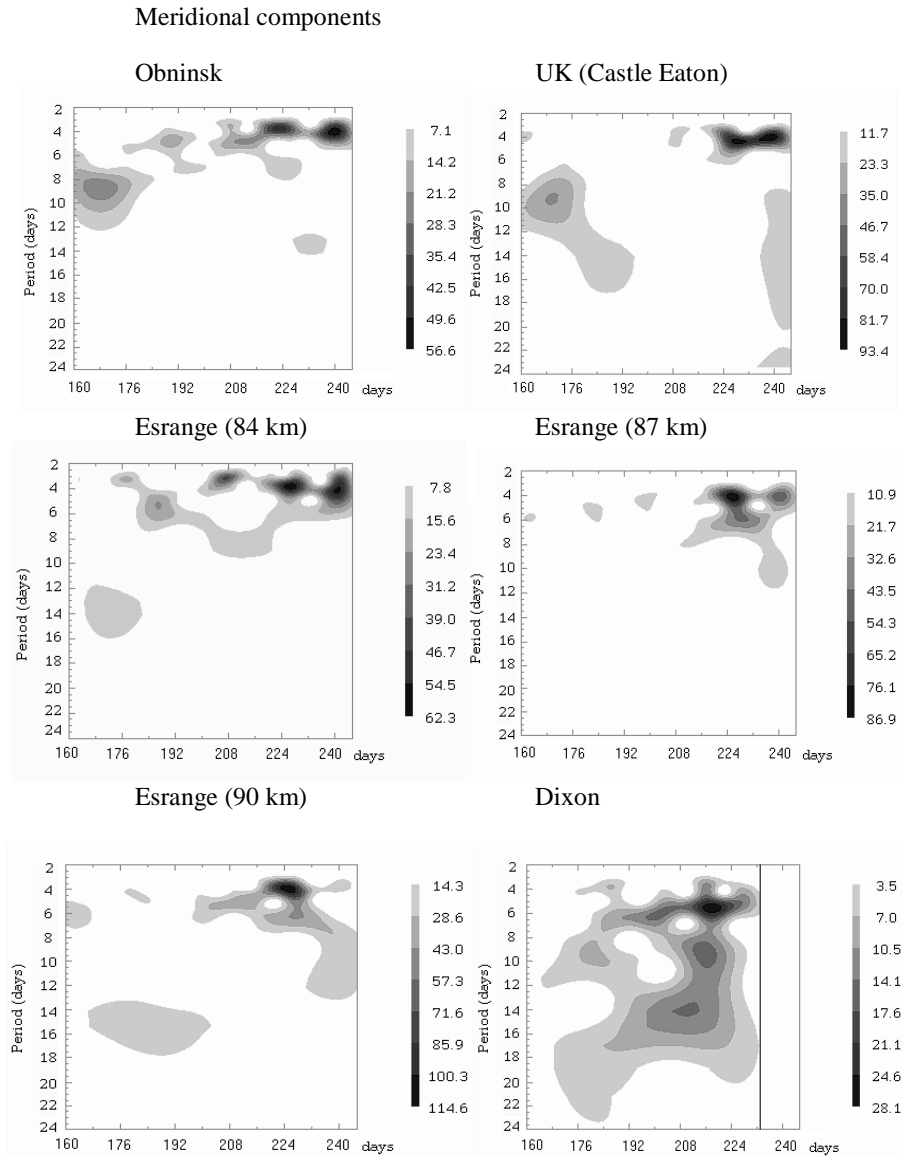
**Fig. 6.** S-transform spectra of day-to-day zonal wind variations. The data are low-pass filtered with a cut-off period of 72 h.

several wave-intensity bursts for both the components. The S-spectrograms of hourly mean wind data for Obninsk, UK, Dixon and Esrage (at 87.5 and 90.5 km) are presented in Figs. 2 and 3. The squares of amplitudes in  $\text{m}^2/\text{s}^2$  are shown by shades of gray. Results for the zonal component are in Fig. 2 and results for the meridional wind component are in Fig. 3.

At first we consider mid-latitudes. There are two intervals of wave activity: before day 205 and after it. A period of about ten days exists between these intervals and it is characterized by weak intensity of the QTDWs. We designate the first interval as I and in the second one we select three main bursts of wave intensity and designate them as II, III

and IV. Note that the zonal wave numbers estimated here are obtained from phase values at only two different points for a given latitudinal belt. So, these numbers are subject to uncertainty and should only be considered as tentative indicators of true longitudinal variability and only for the specified longitudinal sector.

At mid-latitudes the first interval I is characterized by (westward) zonal wave number  $S=2.8\pm 0.3$  (2.8), centered at day 189 for the zonal wind component and  $S=3.3\pm 0.2$  (3.3), centered at day 192 for the meridional one. The parentheses include values estimated from the S-transform. This QTDW has a period of 51–53 h, so the wave is very similar to the normal Rossby-gravity wave mode.

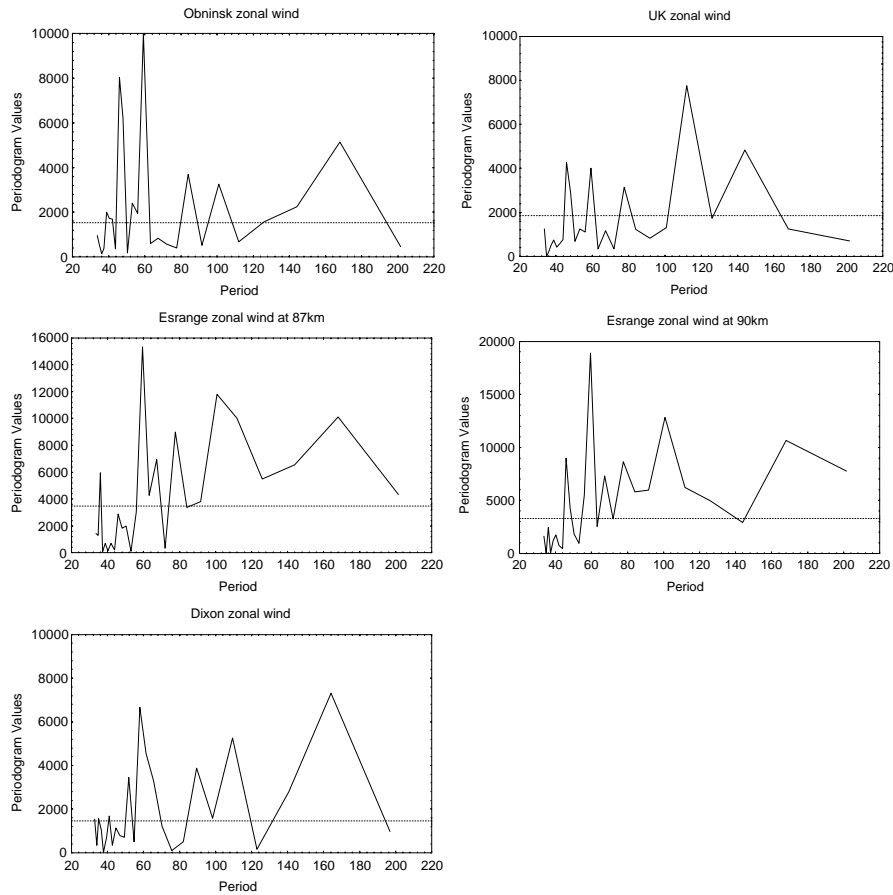


**Fig. 7.** As for Fig. 6, but for day-to-day meridional wind variations.

Burst II is characterized by zonal wave number  $S=4.2\pm 0.2$  (4.2) for the zonal component and  $S=4.1\pm 0.1$  (4.0) for the meridional one centered at day 211. For burst III we obtained  $S=2.8\pm 0.8$  (3.0), centered at day 221 for the zonal component and  $S=3.5\pm 0.2$  (3.2) for the meridional one, respectively. Finally, burst IV has  $S=4.1\pm 0.4$  (3.8) for the zonal component and  $S=2.5\pm 0.3$  (2.2) for the meridional component, centered at day 228. For this last case different bursts are observed in the zonal and meridional components with different zonal wave numbers and different periods of oscillations. As can be seen from the above, values of the zonal wave numbers obtained from both methods are very close, as was shown by Portnyagin et al. (1999). Similar values of the zonal wave numbers were obtained by the cross-wavelet analysis, too. Therefore, for the sake of simplicity we estimate the zonal wave numbers only from the S-transform in the next part of our study.

There are only two common peaks for Dixon and Esrance and both are after day 205. Possibly it is a result of height averaging at Dixon. The wavelet spectra for the Esrance data are presented in Fig. 4, where amplitudes in m/s are shown by shades of gray. Significant variability of spectra with height is well seen. It may be related to superposition of several waves with lower amplitudes than at mid-latitudes, as can be seen below. Among these waves there is an oscillation with a period of 40 h and  $S=2$ . Due to its lower zonal wave number than that of the QTDWs, this wave is more prominent at high-latitudes and introduces additional variability into the wavelet spectra. Another possible contributor to the observed difference could be latitudinal variation of phases and the difference between Dixon and Esrance latitudes. A wave activity burst is occurring in the wavelet spectra, when spectral components forming this burst are in phase. A small phase shift between the spectral components could destroy a burst considered at a given time moment.





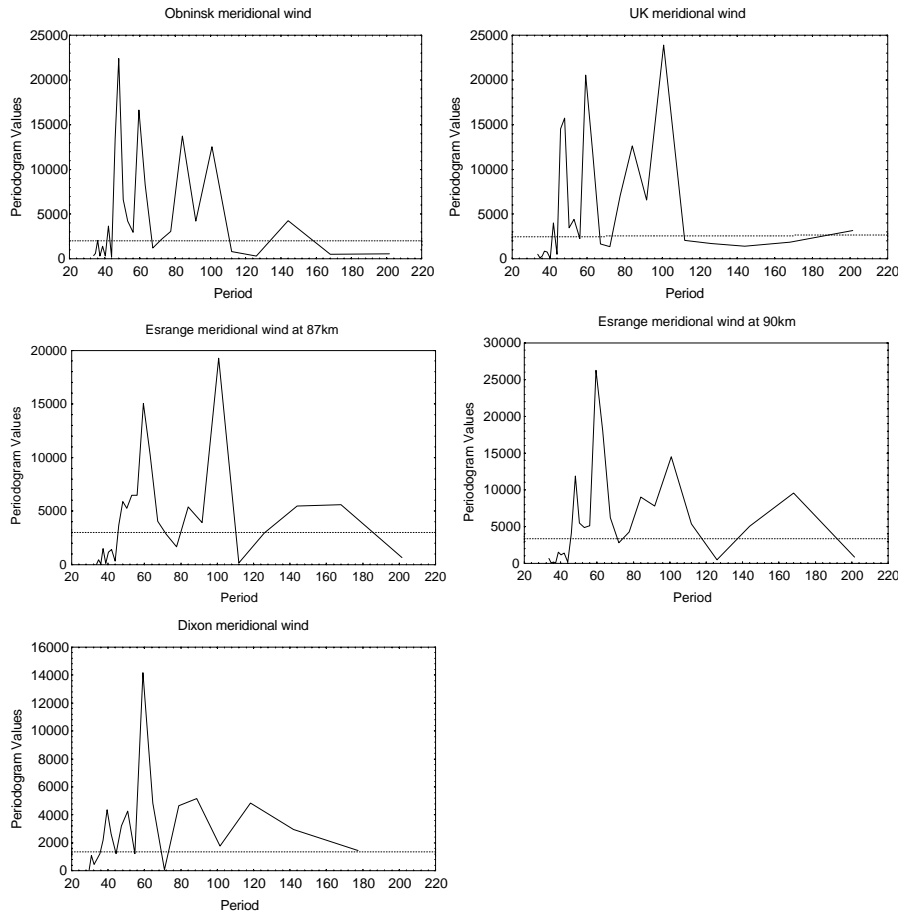
**Fig. 8. (a)** Spectra of zonal wind oscillations from 21 July 2000 to the end of measurements.

The first common burst of the high-latitude wave activity is near day number 211 and so it coincides with burst II at middle latitudes. We obtained zonal wave number  $S=3.6$  for the zonal component and  $S=2.5$  for the meridional one. The next common burst is near day 228 at high-latitudes. It has for the zonal wind component  $S=2$  and for the meridional wind component  $S=3.2$ . This burst coincides with burst IV at mid-latitudes. There is also a weak burst in the meridional wind component near day 221 with  $S=2.6$ . The vertical phase profiles for both high-latitude bursts are shown in Fig. 5. Their vertical wavelength is about 100 km. The oscillation periods are different for the meridional and zonal wind components, and there is a tendency towards the occurrence of the longer periods in the meridional wind.

### 3.2 Long-period ( $\geq 4$ days) oscillations

We now provide a short description of the long-term oscillations in the wind field in pursuit of our aim to concentrate on the nonlinear interaction between the QTDW and the planetary waves of longer periods. Note that the zonal wave numbers estimated here are obtained from phase values at two different points for a given latitudinal belt and for rather small amplitudes. So, these numbers are subject to uncertainty and should only be considered as tentative indicators

of true longitudinal variability and only for the specific longitudinal sector. Spectra of the long-period oscillations (with periods greater than 4 days) are presented in Figs. 6 and 7. It is worth remembering here that all periods and time moments of maximum activity are found by wavelet spectra with definite errors (see Portnyagin et al., 2000), and the significance of an oscillation depends on its duration. All of these wind variations have small amplitudes and are presented because they are features of the spectra, which are significant. Although they do not have large amplitudes, they are observed almost simultaneously at the different stations. While the spectral amplitudes may be significant at each station for a given period, the fact remains that the dominant fraction of spectral energy at one station is often not at that common period. However, the existence of an oscillation is not confined to contours of maximum amplitudes. For zonal wave-number estimation we use those time intervals when an oscillation exists at two stations of different longitude. Additional confidence in our wave-number estimations follows from the fact that the phases of considered oscillations practically do not change during the time of their existence at stations under study (i.e. including time intervals with small and maximum amplitudes). Therefore, there is a very small probability that these long-term neutral wind variations are stochastic



**Fig. 8. (b)** As for (a), but for meridional wind oscillations.

noise. Further, they may be indicative of the existence of larger magnitude waves with the same frequency and zonal wave number at lower latitudes. Incorporation of these oscillations and their zonal wave-number estimates into a theory of QTDW variability is of course strengthened by corroborating experimental, theoretical and modeling evidence. Some of this evidence is presented in the following.

At mid-latitudes one can observe in the zonal component a significant 5-day wind variation with  $S=1$  near day 190 (days 176–190) and a 6–7 day oscillation near day 230 with  $S=1.5$ , a 12-day wind variation with  $S=0.5$  near the beginning of the measurements, and a 20-day variation with  $S=0.43$ . The 6–7 day oscillations are important for our following consideration. Comparing zonal winds at Obninsk and UK one can observe a strong peak at Obninsk and no peak at UK. But a response at this frequency exists in the UK spectra and has significant amplitude of the same order as that at Obninsk. This oscillation is observed at high-latitudes, too (see below). For an explanation of the difference we could suggest the simultaneous existence of two oscillations with close periods and different zonal wave numbers 1 and 2. Data from other longitudes are needed for a comprehensive study of this feature. In the meridional mid-latitude prevailing wind we can

define a 9-day oscillation with  $S=2.2$  at the beginning of the measurements. Near day 224 for both components a 4-day wind oscillation is visible. We obtained 1.7 and 2.2 for zonal wave numbers of these oscillations for meridional and zonal components correspondingly.

At high-latitudes there exists a 12-day oscillation during the QTDW event and a 14-day oscillation at the end of this event in the zonal wind. These oscillations have the same zonal wave numbers,  $S\sim 1$ , and are possibly a manifestation of one and the same wave. It mainly occupies the lower part of the meteor zone. At the end of the studied interval there are 6–7-day oscillations observed in the zonal wind for both stations and in the meridional wind at Esrange. For the 6–7-day wave observed in the zonal component we obtained  $S=0.54$ . It should be noted that these 6–7-day oscillations are observed at the same time as those at mid-latitudes. In this case we possibly observe one wave at different latitudes. The wave number estimates at middle and high latitudes hint at a value of  $\sim 1$ .

Besides these waves it is easy to find in Fig. 4 oscillations with a period of about 4 days in the zonal and meridional wind near day 238 and an oscillation with a period of about 4 days in the meridional wind near day 225. These

oscillations are common for high- and mid-latitudes. The 4-day wave in the meridional wind is a wave number 2 westward propagating wave at high- and mid-latitudes near day 225. Day 238 is too close to the end of our data to obtain reliable wave-number estimations. We obtained  $S=3.4$  for the zonal component and  $S=3$  for the meridional component at mid-latitudes. Only the wave number for the zonal component can be estimated from the high-latitude data. It equals 3.3. These waves have periods that are slightly greater than 4 days near day 225. The wave observed near day 225 exists during the same time as the QTDWs.

Note that the burst II has  $S=4$  for both components, while the latest burst has this zonal wave number only for the zonal wind component. This can mean that the wave source is changing. This is confirmed by an excitation of the waves with a period of 4 days and  $S=3$  after the 2-day wave events. In particular, the velocity of the summer jet decreases and waves of longer periods and larger zonal wave numbers may be forced by the instability.

### 3.3 QTDW bursts: spectral content

The previous description is based on wavelet spectra and does not distinguish spectral contents of different QTDW bursts. Now we will consider this for the sequence of the QTDW bursts after day 205.

Spectra of the wind variations after day 205 are presented in Figs. 8a–b. Only part of the whole spectrum is shown for convenience. It is seen that the QTDW consists mainly of two spectral components: about 47 h and about 60 h. There are additional peaks with periods of 80 h and 100 h. From the cross-spectra for UK and Obninsk data we found  $S\sim 3$  for the 60-h oscillation,  $S\sim 2.5$  for the 80-h oscillation and  $S\sim 3$  for the 100-h oscillations. The 100-h wave is possibly due to the instability. It is clear that the nonlinear interaction between the 47-h wave and another wave with a longer period may force two secondary waves with peaks that are 60 h and 80 h. This interaction needs a wind oscillation with a period of 9–10 days and this oscillation is not visible at the heights of our observations but could exist at lower altitudes and/or at other latitudes. Indeed, this wave is revealed from the stratospheric data for the 2000th year (Pogoreltsev, private communications, 2002) and an average over ten years spectra for the Southern Hemisphere are presented by Fedulina et al. (2003), and this is a common normal-mode frequency of the atmosphere. Note that there is a peak with a period of 6–7 days and we have obtained  $1/60\text{ h}\approx 1/100\text{ h}+1/(6\text{--}7\text{ days})$ . It possibly means decoupling of the 60-h wave with  $S=3$  into the waves with  $S=1$  (about 6–7 days) and  $S=2$  (about 100 h), and these waves are forced in the mesosphere. So, the 4-day wave observed near day 225 might be a result of this decoupling process. The starting mechanism for this decoupling could be the Rossby wave instability (e.g. Baines, 1976) connected with nonlinear resonant wave triads composed of one finite and two infinitesimal components. Two infinitesimal components are provided by noise and conditions for the creation of the resonant triad select oscillations from the noise. Baines

(1976) examined the barotropic stability of Rossby waves on a sphere to small disturbances using numerical simulations.

Although the phase speed of the 4-day wave with  $S=2$  is the same as that of the 2-day wave with  $S=4$  and the wave may, therefore, be excited by the jet instability, this alternative source for the wave is also plausible.

It is interesting to note that other periods different from 60 h and 80 h are observed at middle and high latitudes. Indeed, we see the wave bursts, each of them containing several main spectral components (47 h and 60 h or 60 h and 80 h). Dixon is located at the highest latitudes and we obtain for its wind data mainly one QTDW spectral component with a period of about 60 hours and zonal wave number 3.

The first burst at Dixon has a period of 54 h for the zonal wind component and 57 h for the meridional wind component. At Esrange we obtained 50–52, 5 h and 54 h, respectively. That corresponds to the simultaneous occurrence of two spectral components with  $S=4$  and  $S=3$  and the influence of the first of them is greater for the zonal component and decreases gradually to the pole. And this is the explanation of our zonal wave number estimation pointed out above.

Burst III is significantly weaker at Dixon than at Esrange. At middle latitudes we observe these oscillations with a period of about 50–54 h, but at Esrange the period is about 58 h for the meridional wind component and has variable values as a function of height for the zonal component. It means again the coexistence of two spectral components. The meridional component of the wave with  $S=4$  or 3 is larger than the zonal one. Therefore, the period of the meridional component is less variable and that is why phase profiles of the zonal components differ from those of the meridional components in Fig. 5.

At middle latitudes Burst IV is characterized by periods of 55 h (UK) and 51 h (Obninsk) for zonal wind components. For meridional components we found a period of 64 h for both stations. These periods are changing with latitude: we found for the zonal component a period of 62 h at Esrange and a period of 57 h at Dixon; for the meridional component they were 66 h and 55 h, respectively. One can see that burst IV is formed at Dixon by two close oscillations. One of them has a period of 62–64 h. It is worthwhile to note an increasing influence of a 40-h oscillation on the S-spectra at Dixon, which is possibly a normal mode (2, 0) and appears near the QTDW activity.

To summarize our observational results, the QTDW activity observed during summer 2000 consists of several spectral components with periods from about 47 to 66 h and zonal wave numbers from about 2 to 4. This activity is started by the 51–53 h waves with  $S=3$ , interpreted as a manifestation of the normal Rossby-gravity mode (3,0), probably amplified by the instability (Salby and Callaghan, 2001). After a recess of the wave activity (for about 10 days) we observed a regular sequence of wave activity bursts. The strongest burst is connected to the spectral component with a period of 47–48 h and  $S=4$  and located at mid-latitudes. Although the bursts demonstrate different periods and zonal wave numbers at different latitudes for both wind components, their spectral

content is rather simple and includes mainly 3 spectral components. The frequencies of these spectral components satisfy the condition for nonlinear interaction with the 9–10 day planetary wave, and this wave was observed in the stratosphere of the Southern Hemisphere.

It should be emphasized here that the observed variability of periods and zonal wave numbers is based upon only two latitudinal points, and thus, open to uncertainty. Thus, we can only suggest possible hypotheses within the scope of the obtained results.

### 3.4 Model results

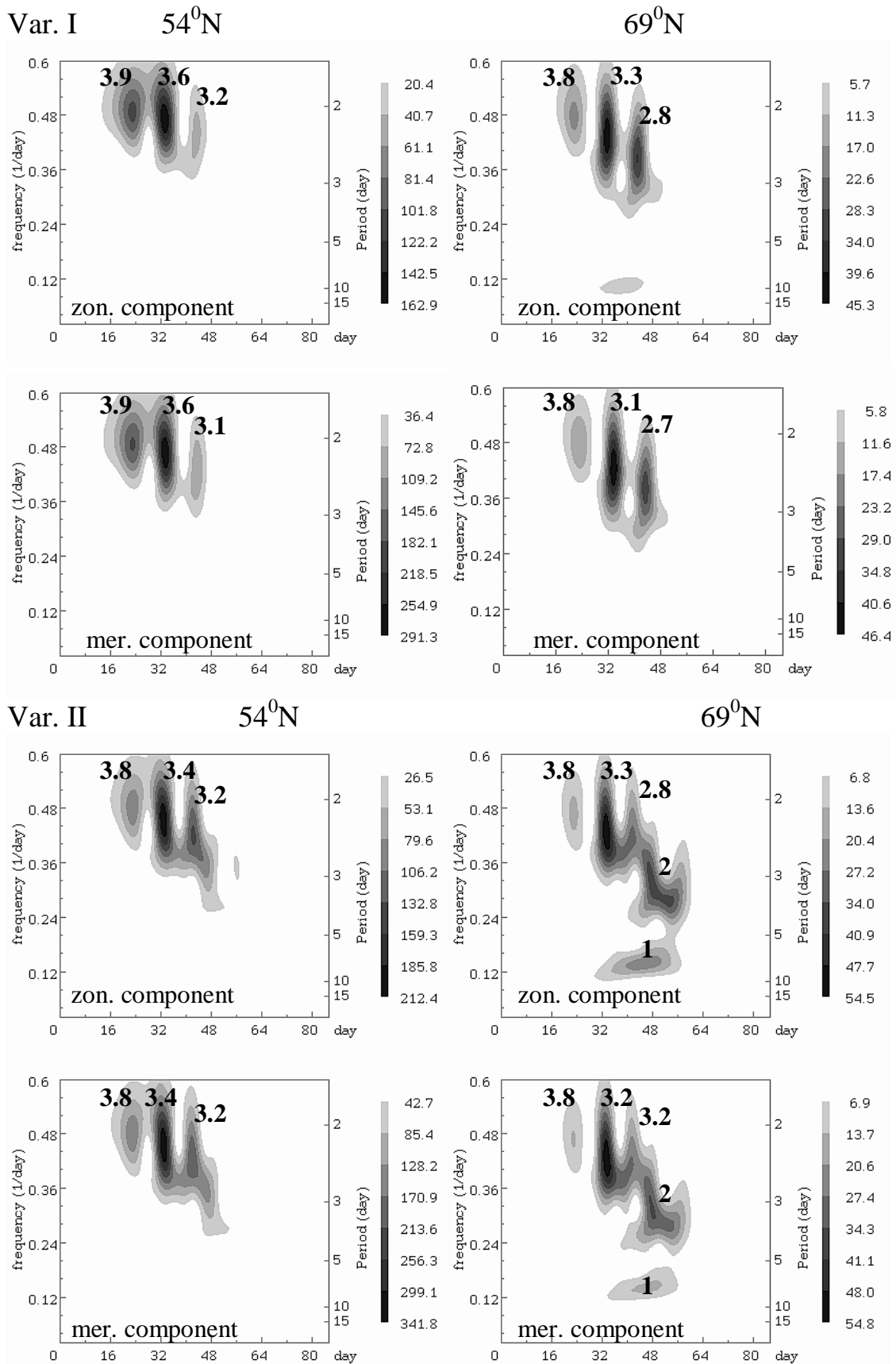
During summer the 2-day wave with  $S=4$  cannot readily propagate from heights of the troposphere and the stratosphere to the mesopause, due to its slow zonal phase speed and the occurrence of critical levels. One needs unrealistically large amplitudes at the bottom to obtain observable values at the mesopause (see, e.g. Salby and Callaghan, 2001). Unfortunately, only a few nonlinear numerical studies have been published recently that attempt to explain the wave occurrence at the summer mesopause of the Northern Hemisphere (i.e. Norton and Thuburn, 1996; Mayr et al., 2001). The numerical investigation of Mayr et al. showed simultaneous excitation of waves with  $S=2-4$ , consistent with our observations during the summer of 2000. However, we observe rather regular sequences of three bursts, and this feature possibly points to a role of the nonlinear interaction between waves. We now consider wave-wave nonlinear interaction as a source of QTDW variability within the context of a numerical model.

In our numerical simulations, the 2-day wave with  $S=4$  was excited artificially at about 60 km, i.e. near a level of maximum QTDW amplitudes in the simulations of Mayr et al. (2001) and Norton and Thuburn (1996). As sources of planetary waves we used thermal sources with amplitudes depending on time, like the Gauss function. For the case of the QTDW the latitudinal structure of the source was taken in such a way as to obtain the same latitudinal temperature structure as presented by Norton and Thuburn (1996). It implies indirectly that we consider the jet instability as a source of the 2-day wave with  $S=4$ . The wave activity before day 205 is considered independent of events after that day.

The amplitude and duration of the thermal forcing for the 2-day wave with  $S=4$  were tuned to the experimental results. Additionally, one or two planetary waves with zonal wave numbers 1 and 2 were forced near the bottom of the model. Periods of these waves were changed from 5 to 20 days. A corresponding Hough mode was taken as the latitudinal structure of the thermal source for these waves. Again, amplitudes and durations of the thermal forcing for these planetary waves were tuned to reproduce the experimental results. Amplitudes of wind variations in each burst of the QTDW activity, the width of the burst, and the number of the bursts and their distribution in the wavelet spectra at different latitudes significantly depend on values of all these parameters.

The first model run tested the possibility to excite the QTDW with  $S=4$  as a secondary wave resulting from the interaction between a long-period planetary Rossby wave (or a stationary wave, too) with  $S=1$  and the normal Rossby-gravity QTDW with  $S=3$ . Both waves are considered as propagating upward from the bottom. The reference distribution of temperature for stationary wave 1 was taken from Barnett and Corney (1985) for July. There was no strong response for the 2-day wave with  $S=4$ . This means that another source of wave 4 should be suggested. The jet instability is suggested through this work. It is interesting to note that the secondary wave from the considered nonlinear interaction could serve as an initial wind perturbation for such an instability development.

We now turn to the use of our numerical model to simulate some salient features of the observations, particularly how the nonlinear interaction between the QTDW and a planetary wave may create a sequence of three wave bursts that have different periods and different zonal wave numbers at high and middle latitudes. A lot of variants (more than 100) with different parameter values were considered. Comparison between numerical and experimental results was carried out by using the  $S$ -spectra for meridional wind components at a height of 90 km and latitudes of Obninsk and Esrange. Only a few of them demonstrated results for the QTDW similar to the observations. The simplest variant consists of nonlinear interaction between a transient 10-day wave with  $S=1$  and a transient 47-h wave with  $S=4$ . Note that we did not take into account the spectral content of the QTDW bursts known from the experimental results and just tried to reproduce the observations. An important numerical result is the appearance of a large-amplitude wave with 60 h period and  $S=3$ . In the first run we demonstrated that it is impossible for the 47-h wave with  $S=4$ . This result has rather a simple explanation for barotropic waves; see Baines (1976). In its turn the 60-h wave interacts with the 10-day wave, too, and then a wave of an 80-h period is forced. The last wave has a small amplitude and appears later than the two previous ones. The initial 47-h wave is practically absent by this time. In Fig. 9 the  $S$ -spectra of the numerically simulated series are presented for latitudes of Obninsk (left) and Esrange (right). Two cases are presented, which differ from one another by a latitudinal shift of the 2-day wave source. The two top panels show zonal and meridional components of the QTDW events for the case when a temperature source is near the equator, and the two bottom panels do it for the second case with a source shifted towards the pole. At the latitude of Dixon the  $S$ -spectra are similar to those of Esrange. The zonal wave numbers are marked in the figure above each burst. These numbers were obtained in the same way as those for the experimental data and contain errors due to the latitudinal effect that increases a zonal wave number by values from 0.2 to 0.3 for the last two bursts at high latitudes. Apparently, there is a strong dependence of the results on a source location. Irrespective of this fact, for both cases three wave bursts are obtained as a result of the nonlinear interaction between the QTDW and a planetary wave with a period of 9–10 days.



**Fig. 9.** Results of numerical oscillations for two different locations of the 2-day wave source.

Thus, using only the S-spectra from the experimental data we could (1) obtain the observed spectral content of the QTDW bursts; (2) suggest an explanation for their appearance and reproduce the number and spacing of the bursts; and (3) reproduce observed periods and zonal wave numbers for the meridional components.

#### 4 Summary

We conclude that the QTDW activity observed during the summer of 2000 is a complex event. It is possibly a result of the combined influence of the nonlinear interaction between different planetary waves and jet instability. The QTDW activity is observed simultaneously at high- and mid-latitudes and is manifested as several bursts of wave activity, revealed using S-transform or wavelet analysis. Initially, this activity is preceded by the 51–53 h wave with  $S=3$ , which is observed mainly at mid-latitudes. After a short recess (for about 10 days near day 205) we observed a regular sequence of three wave activity bursts, the strongest of them being connected to a QTDW wind oscillation (spectral component) with a period of 47–48 h and  $S=4$ , mainly at mid-latitudes. We hypothesize that these three bursts may be created by destructive and constructive interference between a few spectral components. The component set includes the primary QTDW and two secondary waves could arise from a sequence of the nonlinear interactions with a 9–10 day oscillation. A 9–10 day oscillation during the interval of our observations has been located in the stratosphere by A. Pogoreltsev (private communication, 2002). The superposition leads to an appearance of bursts with new frequencies and fractional zonal wave numbers in the observations. We do not explain all fractional zonal wave numbers only by the existence of a few spectral components. This result follows from numerical simulations and serves as an additional source for understanding the experimental results. It is also revealed that the zonal wave numbers of observed bursts range from about 4 to about 2 for mid- and high-latitudes. The main spectral component at Dixon and Esrange latitudes is the 60-h oscillation with  $S=3$ . Thus, it is demonstrated that the correspondence between the high- and mid-latitude 2-day events may be understood in the framework of nonlinear wave interaction, linear wave interference and instability. The wave coupling can explain the source of the new waves, but it also defines the time of the wave occurrence. It should be emphasized here that the observed variability of periods and zonal wave numbers is only based upon two latitudinal points, and thus, open to uncertainty. Thus, we can only suggest possible hypotheses within the scope of the obtained results. In addition to the QTDW bursts, we obtained as well a forcing of the 4-day wave with  $S=2$  and the 6–7-day wave with  $S=1$  due to the nonlinear decoupling of the 60-h wave with  $S=3$ . The starting mechanism for this decoupling could be the Rossby wave instability (e.g. Baines, 1976). This result corresponds to the observed day-to-day wind variability during the QTDW events. An interesting feature of the final stage of the two-day-wave events in

summer 2000 is an occurrence of the strong 4–5 day waves with  $S=3$ .

#### Appendix

The nonlinear time-dependent model used in our analysis employs a background wind field similar to the climatic empirical model of Fleming et al. (1988). The numerical model is based on that of Rose (1983), Jacobs et al. (1986). The horizontal momentum equation, the thermodynamic equation, the continuity equation and the hydrostatic equation in spherical log-pressure co-ordinates are solved by the explicit finite-difference method. Unlike the models referenced above, we used an expansion in Fourier harmonics in longitude, and instead of a gravity wave parameterization we used a body force like that proposed by Fritts and Luo (1995). The radiation processes are parameterized by the Newtonian cooling as  $\alpha(T-T_0)$ , where the rate coefficient  $\alpha$  was adopted from Zhu (1993) and  $T_0$  is the reference temperature from Fleming et al. (1988). The finite difference grid has a step of  $3^\circ$  in latitude and 0.25 in height from  $z=0$  to  $z=20$ . Here,  $z=-\ln(P/P_S)$ ,  $P$  is pressure, and  $P_S$  is a constant reference pressure (1000 mb). The expansion in longitude is performed in terms of  $\exp(im\lambda)$ , where  $m$  ( $=-6, \dots, +6$ ) is a zonal wave number and  $\lambda$  is longitude. The coefficients of dynamic molecular viscosity and molecular thermal heat conduction were taken from Forbes and Garrett (1979), eddy viscosity was adopted from Hagan et al. (1995) and hydromagnetic effects are included in a simple form as in Forbes and Garrett (1979). Some horizontal smoothing was applied to calculated fields that is equivalent to horizontal dissipation of the fourth order with a rate of about  $10^{15} \text{ m}^4/\text{s}$ . Background field distributions are obtained from a model run with initially motionless atmosphere and horizontally uniform temperature.

#### Boundary conditions

At the bottom we imposed the condition that the vertical velocity is 0. It means  $d\Phi/dt=0$  at  $z=0$ , where  $\Phi$  is geopotential. For velocity components and the nonzonal component of temperature ( $m \neq 0$ ) we take the conditions like those utilised by Forbes (1982) to simulate the surface interactions. For the mean zonal component of temperature we take a time independent temperature distribution from Fleming et al. (1988).

The log-pressure vertical velocity ( $dz/dt$ ), vertical gradients of velocities and nonzonal components of temperature ( $m \neq 0$ ) are set equal to 0 at the upper horizontal boundary. The mean zonal component of temperature ( $m=0$ ) does not depend on time at the upper boundary and was estimated from models of Fleming et al. (1988). As sources of planetary waves we used thermal sources. The source of the long-period wave was placed near the tropopause and had a latitudinal distribution corresponding to that of a normal mode. The source of the 2-day wave with  $S=4$  was placed at level  $z=8$  and had a latitudinal distribution which is like a hump

of the  $\cos(3 \cdot \theta)$ , where  $\theta$  is colatitude. The amplitudes of the primary waves and their time of existence are tuned to reproduce the features observed in the experiment.

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