



Impacts of cloud heterogeneities on cirrus optical properties retrieved from space-based thermal infrared radiometry

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Received: 20 May 2014 – Published in Atmos. Meas. Tech. Discuss.: 27 August 2014

Revised: 2 January 2015 – Accepted: 12 January 2015 – Published: 9 February 2015

Abstract. This paper presents a study, based on simulations, of the impact of cirrus cloud heterogeneities on the retrieval of cloud parameters (optical thickness and effective diameter) for the Imaging Infrared Radiometer (IIR) on board CALIPSO. Cirrus clouds are generated by the stochastic model 3DCLOUD for two different cloud fields and for several averaged cloud parameters. One cloud field is obtained from a cirrus observed on 25 May 2007 during the airborne campaign CIRCLE-2 and the other is a cirrus uncinus. The radiative transfer is simulated with the 3DMCPOL code. To assess the errors due to cloud heterogeneities, two related retrieval algorithms are used: (i) the split-window technique to retrieve the ice crystal effective diameter and (ii) an algorithm similar to the IIR operational algorithm to retrieve the effective emissivity and the effective optical thickness. Differences between input parameters and retrieved parameters are compared as a function of different cloud properties such as the mean optical thickness, the heterogeneity parameter and the effective diameter. The optical thickness heterogeneity for each $1\text{ km} \times 1\text{ km}$ observation pixel is represented by the optical thickness standard deviation computed using $100\text{ m} \times 100\text{ m}$ subpixels. We show that optical thickness heterogeneity may have a strong impact on the retrieved parameters, mainly due to the plane-parallel approximation (PPA assumption). In particular, for cirrus clouds with ice crystal diameter of approximately $10\text{ }\mu\text{m}$, the averaged er-

ror on the retrieved effective diameter and optical thickness is about $2.5\text{ }\mu\text{m}$ ($\sim 25\%$) and -0.20 ($\sim 12\%$), respectively. Then, these biases decrease with increasing effective size due to a decrease of the cloud absorption and, thus, the PPA bias. Cloud horizontal heterogeneity effects are greater than other possible sources of retrieval errors such as those due to cloud vertical heterogeneity impact, surface temperature or atmospheric temperature profile uncertainty and IIR retrieval uncertainty. Cloud horizontal heterogeneity effects are larger than the IIR retrieval uncertainty if the standard deviation of the optical thickness, inside the observation pixel, is greater than 1.

1 Introduction

In the context of global climate change, the representation and role of clouds are still uncertain. For example, ice clouds play an important role in the climate and on the Earth's radiation budget (Liou, 1986). Cirrus clouds lead mainly to a positive radiative forcing due to their high temperature contrast with respect to the surface. However, the cirrus radiative forcing could depend on the cirrus optical thickness, altitude and ice crystal effective size (Katagiri et al., 2013). Consequently, to improve our knowledge, it is essential to assess the feedback and climate effects of these clouds (Stephens,

1980). Global satellite observations are well suited to monitoring and investigating cloud evolution and characteristics because passive top-of-atmosphere (TOA) radiometric measurements allow for retrievals of cloud properties such as optical thickness and ice crystal effective diameter. In this work, we focus on infrared measurements obtained by the Imaging Infrared Radiometer (IIR; Garnier et al., 2012, 2013) onboard the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observations (CALIPSO).

Because of operational constraints (lack of information regarding the 3-D structure of the atmosphere, time constraints, etc.), satellite-based cloud retrieval algorithms assume that clouds are homogeneous and infinite between two planes. This assumption of 1-D radiative transfer is called the homogeneous independent pixel approximation (Cahalan et al., 1994) or independent column approximation (Stephens et al., 1991). However, in a real atmosphere, clouds have 3-D structures, i.e. horizontal and vertical heterogeneities and the simplified 1-D atmosphere assumption may lead to biased cloud property retrievals (Fauchez et al., 2014). Many studies have been conducted to determine the impact of cloud heterogeneities on cloud products derived from solar spectral measurements. These studies primarily focused on warm clouds such as stratocumulus (Varnai and Marshak, 2001; Zinner and Mayer, 2006; Kato and Marshak, 2009, etc.) and showed that the sign and amplitude of retrieval errors depend on numerous factors, such as the spatial resolution, wavelength, geometry of observation and cloud morphology. Concerning cirrus clouds, Fauchez et al. (2014) showed that cirrus cloud heterogeneities lead to non-negligible effects on brightness temperatures (BT) and that these effects mainly depend on the standard deviation of the optical thickness inside the observation pixel. The retrieval of cloud properties using radiances or BT may thus be impacted by the heterogeneity effects. In this work, we extend the study of Fauchez et al. (2014) to investigate the impacts of cirrus heterogeneities on cloud optical property (optical thickness and ice crystal effective size) retrievals using simulations of radiometric measurements of IIR in three typical spectral bands, namely 8.65, 10.60 and 12.05 μm .

In the thermal infrared atmospheric window (8–13 μm), cloud optical properties (optical thickness and ice crystal effective size) are retrieved using the split-window technique (SWT) (Inoue, 1985; Parol et al., 1991; Dubuisson et al., 2008). This method is generally limited to thin cirrus clouds (optical thickness less than approximately 3 at 12 μm) and small crystals (effective diameters smaller than approximately 40 μm). In the visible and near-infrared spectra, cloud optical properties are commonly retrieved using the Nakajima and King method (Nakajima and King, 1990) that combines measurements in visible and near-infrared channels for optically thicker cirrus clouds and larger ice crystals. Cooper et al. (2007) combined these two methods for MODIS measurements to treat thin and thick cirrus simultaneously.

The paper is organized as follows. In Sect. 2, we present a short description of the modeling tools used in this study: (i) the cloud generator 3DCLOUD (Szczap et al., 2014), (ii) the radiative transfer code 3DMCPOL (Cornet et al., 2010; Fauchez et al., 2014) and (iii) two related retrieval algorithms. In Sect. 3, we present possible retrieval errors due to the 1-D approximation. In Sect. 4, we compare heterogeneity effects with other possible error sources considered in this paper such as those due to cloud vertical heterogeneity, surface temperature or atmospheric temperature profile uncertainty, as well as the IIR retrieval uncertainty. Conclusions and perspectives are given in Sect. 5.

2 Numerical models

2.1 3-D ice water content generation

The stochastic model 3DCLOUD (Cornet et al., 2010; Szczap et al., 2014) is employed to generate realistic 3-D cirrus clouds. This model uses a simplified dynamical and thermodynamical approach to generate heterogeneous 3-D clouds as well as a Fourier transform framework to constrain scale invariant properties (Hogan and Kew, 2005; Szczap et al., 2014). Two different cirrus fields were simulated (Fig. 1) in a mid-latitude summer (MLS) atmosphere.

The first cirrus field has been modeled from meteorological profiles presented by Starr and Cox (1985) coupled with a wind profile to form virgas. The cloud layer is defined by the mean optical thickness τ_c , the standard deviation of the optical thickness on the entire field σ_{τ_c} , the cirrus heterogeneity parameter $\rho_{\tau} = \sigma_{\tau_c}/\tau_c$ (Szczap et al., 2000) and the ice crystal effective diameter D_{eff} for an aggregate crystal shape (Yang et al., 2005). D_{eff} is defined as

$$D_{\text{eff}} = \frac{3 \int V(L)n(L)dL}{2 \int A(L)n(L)dL}, \quad (1)$$

where L is the maximum crystal size, $V(L)$ is the volume of the crystal, $A(L)$ is the projected area and $n(L)$ is the size distribution (Yang et al., 2000).

Eight cirrus clouds are generated (Table 1) by varying the above parameters to cover the characteristics of typical cirrus clouds (Sassen and Cho, 1992; Szczap et al., 2000; Carlin et al., 2002; Lynch et al., 2002). Note that the effective diameter of cirrus cases 3 to 5 ($D_{\text{eff}} = 9.95 \mu\text{m}$) is probably too small for cirrus with a mean optical thickness of 1.80 because aggregations processes tend to increase the effective size (Fig. 12 of Garnier et al., 2013). Cloud heterogeneity effects are probably slightly overestimated due to the too-small crystal effective size (heterogeneity effects are larger for small effective sizes) with respect to the mean cirrus optical thickness. Nevertheless, cirrus cases 3 to 5 are useful for understanding how heterogeneity effects increase with the optical thickness heterogeneity parameter (ρ_{τ} in Table 1), which increases from 0.7 to 1.1 and 1.5 with other

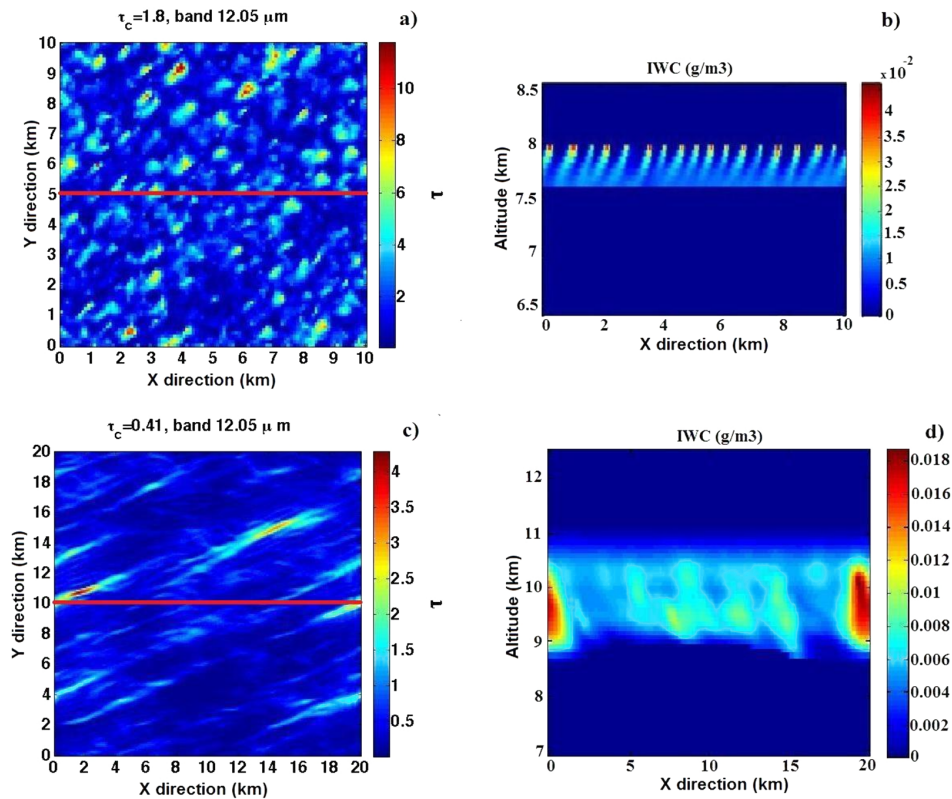


Figure 1. Top figures: Cirrus generated from realistic meteorological conditions (Starr and Cox, 1986; Hogan and Kew, 2005) with (a) the $10 \text{ km} \times 10 \text{ km}$ optical thickness field simulated at $12.05 \mu\text{m}$ with a horizontal spatial resolution of 100 m and (b) the x - z view through the red line of (a) of the cirrus IWC with a vertical spatial resolution of 58 m . Bottom figures: CII cirrus simulation based on optical and microphysical properties of the cirrus observed during the CIRCLE-2 campaign on 25 May 2007: (c) the $20 \text{ km} \times 20 \text{ km}$ optical thickness field at $12.05 \mu\text{m}$, with a horizontal spatial resolution of 100 m and with a mean optical thickness $\tau_c = 0.41$ observed by IIR at $12.05 \mu\text{m}$, and (d) the x - z view through the red line of (c) the cirrus IWC with a vertical resolution of 58 m .

cloud properties held constants. Two cirrus cloud cases are presented in Fig. 1. The first cloud structure is presented in Fig. 1a and b. Figure 1a presents the $10 \text{ km} \times 10 \text{ km}$ optical thickness field at $12.05 \mu\text{m}$ with a spatial resolution of 100 m , and Fig. 1b presents the x - z view of the ice water content (IWC) of cirrus case 3.

Figure 1c and d show cirrus generated from measurements obtained on 25 May 2007 during the CIRCLE-2 airborne campaign (Mioche et al., 2010). In situ measurements provided by the aircraft, as well as IIR radiometric measurements (mean optical thickness and mean heterogeneity parameter), are used as input for 3DCLoud. In addition, meteorological data from the European Center for Medium-Range Weather Forecasts are used to constrain the meteorological profiles (wind speed and orientation, temperature, humidity, etc.). The scale invariant properties of every cirrus case presented in Table 1 are controlled by a constant spectral slope ($-5/3$) for all scales and altitude levels. This agrees with the spectral slope of the backscattering coefficient measured at 532 nm at different altitudes by the Cloud-Aerosol Lidar with Orthogonal Polarization (CALIOP) on-board CALIPSO and

the extinction coefficient measured by the polar nephelometer at the aircraft altitude (Fauchez et al., 2014).

2.2 Optical property parametrization

Cirrus optical properties are difficult to characterize because of the diversity of crystal sizes, shapes and orientations in a cirrus cloud. Several parametrizations were developed for visible and infrared wavelengths (Magono, 1966; Labonnote et al., 2000; Yang et al., 2001, 2005; Baum et al., 2005b, 2011; Baran and Labonnote, 2007; Baran et al., 2009; Baran, 2012; Baran et al., 2013). For cirrus cases 1 to 8 we employ the aggregate ice crystal model (Yang et al., 2001, 2005) with a monodisperse distribution used in the IIR retrieval algorithm (Garnier et al., 2013) that provides an extinction coefficient, a single-scattering albedo and an asymmetry factor (Yang et al., 2001, 2005). Note that Dubuisson et al. (2008) have shown that the IIR thermal infrared channels are weakly sensitive to the ice crystal shape and almost insensitive to the size distribution. The IIR retrieval algorithm uses three ice crystal shapes (Garnier et al., 2012, 2013), namely a solid column, aggregate and plate. The phase functions of these

Table 1. Mean cloud properties of the cirrus generated by 3DCLLOUD. “CTA” corresponds to the cirrus-top altitude; “OP” corresponds to the optical properties parametrization; “Yal” represents the model of ice crystals developed by Yang et al. (2001, 2005) for aggregates ice crystals; and “Bal” represents the parametrization of ice crystals’ optical properties developed by Baran et al. (2009), Baran (2012), and Baran et al. (2013); τ_c is the cloud mean optical thickness; σ_τ is the cloud standard deviation of the optical thickness estimated from the optical thickness of the subpixels at the scale of $100\text{ m} \times 100\text{ m}$; ρ_τ is the cloud heterogeneity parameter defined as the ratio of σ_τ by τ_c ; and D_{eff} is the ice crystal effective diameter.

Cirrus	CTA (km)	τ_c	σ_τ	ρ_τ	D_{eff} (μm)	OP
1	7.97	0.45	0.32	0.7	9.95	Yal
2	7.97	0.90	0.63	0.7	9.95	Yal
3	7.97	1.80	1.26	0.7	9.95	Yal
4	7.97	1.80	1.98	1.1	9.95	Yal
5	7.97	1.80	2.70	1.5	9.95	Yal
6	7.97	1.80	1.26	0.7	20.09	Yal
7	7.97	1.80	1.26	0.7	40.58	Yal
8	11.06	0.90	0.63	0.7	9.95	Yal
CII-1	11.06	0.41	0.32	0.77	heterogeneous	Bal
CII-2	11.06	0.81	0.62	0.77	heterogeneous	Bal
CII-3	11.06	0.90	0.63	0.70	9.95	Yal

particles are relatively smooth in the thermal infrared with a small forward peak (asymmetry factor g usually below 0.9) and can be approximated by the Henyey–Greenstein phase function. While this assumption is certainly problematic for irregular crystal shapes, as shown by Baum et al. (2005a, b), we use the Henyey–Greenstein phase function to remain consistent with the official IIR retrieval algorithm (Garnier et al., 2012, 2013). For these cirrus cases, the optical properties are constant over the entire cloud.

In order to generate 3-D and heterogeneous cloud optical properties fields for the CII-1 and CII-2 cirrus cases, we used the parametrization of Baran et al. (2009, 2013) and Baran (2012). This parametrization, derived from in situ measurements of more than 20 000 particle-size distributions (Field et al., 2005, 2007), gives the optical coefficients as a function of IWC and temperature.

2.3 TOA brightness temperature simulations

TOA brightness temperatures in the three IIR thermal infrared channels (8.65, 10.60 and $12.05\ \mu\text{m}$) are simulated with the 3DMCPOL code developed in the visible range by Cornet et al. (2010) and extended to the infrared range by Fauchez et al. (2014). 3DMCPOL is a forward Monte Carlo algorithm using the local estimate method (Marshak and Davis, 2005; Mayer, 2009) and is able to simulate radiances and brightness temperatures from the visible to the infrared range, including the polarization. The atmosphere is subdivided in voxels (3-D pixels), with a constant horizontal size (dx , dy) and a variable vertical size (d_z). Each voxel is described by the extinction coefficient σ_e , the single-

scattering albedo ϖ_0 , the phase function and the cloud temperature T_c .

3-D BT are first simulated at $100\text{ m} \times 100\text{ m}$ spatial resolution and are then averaged to the IIR spatial resolution of $1\text{ km} \times 1\text{ km}$ ($\text{BT}_{1\text{ km}}^{3\text{-D}}$). 1-D BT are obtained by averaging the optical property field to $1\text{ km} \times 1\text{ km}$ spatial resolution before simulating the BT ($\text{BT}_{1\text{ km}}^{1\text{-D}}$).

Note that the statistical uncertainty of these simulations is below 0.5 K, which is less than the IIR accuracy of about 1 K. Comparisons between 3DMCPOL statistical uncertainty, IIR accuracy and heterogeneity effects can be found in Fauchez et al. (2014) (Figs. 8 and 10) for the same cloud scenes. This statistical uncertainty is reached by simulating between 5 and 10 billion photons for each case.

2.4 Retrieval algorithms of cloud parameters

Two related algorithms are used to retrieve cloud products: the split-window technique (Inoue, 1985; Parol et al., 1991; Dubuisson et al., 2008) to retrieve the effective diameter and an algorithm similar to the IIR operational algorithm to retrieve the effective emissivity and the effective optical thickness.

In the thermal infrared atmospheric window, the SWT is one of the most used methods to retrieve the effective diameter and the cloud optical thickness using the difference of brightness temperatures between two thermal infrared channels (Parol et al., 1991; Radcliff et al., 2003; Dubuisson et al., 2008; Garnier et al., 2012, 2013). Figure 2 shows brightness temperature difference (BTD) for varying optical thickness (0–50 at $12.05\ \mu\text{m}$) and eight effective diameters (D_{eff}) as a function of the $12.05\ \mu\text{m}$ BT (BT_{12}). Each “arch” corresponds to a single effective size, with BTD decreasing with increasing particle size and optical thickness decreasing along each arch from opaque cloud (low BT) to clear sky (high BT). It is evident that the sensitivity of the SWT to large particles ($D_{\text{eff}} > 40\ \mu\text{m}$) is weak, one of the main disadvantages of this method that can only accurately determine the effective size of particles smaller than approximately $40\ \mu\text{m}$ for cirrus clouds with an optical thickness approximately between 0.5 and 3 (Dubuisson et al., 2008; Sourdeval et al., 2012). Dubuisson et al. (2008) also show that the SWT retrieval accuracy for ice crystal effective diameter is between 10 and 25 % and for the optical thickness is about 10 %. We note that the amplitude of the BTD_{8-10} arches is significantly smaller than the two others because its sensitivity to D_{eff} is weaker. Consequently, this channel pair will not be used in this study.

Similar to the SWT, the IIR operational algorithm (Garnier et al., 2012) uses radiance differences between channels, though in a different way. Intermediate products (effective emissivity, effective optical thickness and microphysical indices) are computed to retrieve the ice crystal effective diameter and shape. The effective emissivity refers to the contribution of scattering in the retrieved emissivity, especially for

small ice crystals in the band at 8.65 μm . One of the major advantages of using the effective emissivity is its independence of cloud-top altitude or geometrical thickness, contrary to the brightness temperature differences used in the SWT. The effective emissivity, $\varepsilon_{\text{eff},k}$, for the channel k is defined as

$$\varepsilon_{\text{eff},k} = [R_k - R_{k,\text{BG}}] / [B_k(T_c, Z_c) - R_{k,\text{BG}}], \quad (2)$$

where R_k is the measured (or simulated) radiance in the channel k , $R_{k,\text{BG}}$ is the measured (or simulated) radiance at TOA for clear sky and $B_k(T_c, Z_c)$ is the radiance of an opaque cloud (black body) located at the centroid altitude Z_c and at the centroid temperature T_c , provided by the GEOS-5 model (Rienecker et al., 2008). The layer centroid altitude is a weighted average altitude based on the attenuated backscattered intensity of the LIDAR signal at 532 nm (Vaughan et al., 2009). Note that, in this study, we set the centroid altitude to the geometrical middle of the cloud.

The effective optical thickness $\tau_{\text{eff},k}$ is then calculated as

$$\tau_{\text{eff},k} = -\ln(1 - \varepsilon_{\text{eff},k}). \quad (3)$$

From $\tau_{\text{eff},k}$, the microphysical indices $\text{MI}^{12/8}$ and $\text{MI}^{12/10}$ are defined as the ratio of $\tau_{\text{eff},k}$ between 12.05 and 8.65 μm channels and 12.05 and 10.60 μm channels, respectively:

$$\text{MI}_{12/8} = \tau_{\text{eff},12} / \tau_{\text{eff},8}, \quad \text{MI}_{12/10} = \tau_{\text{eff},12} / \tau_{\text{eff},10}. \quad (4)$$

These microphysical indices strongly depend on the microphysical and optical properties of the cloud layer, namely the effective diameter and shape of the ice crystals. From a look-up table (LUT) of the microphysical indices as a function of the effective diameter and shape precalculated by the FASDOM code (Dubuisson et al., 2005), two values of effective diameters ($D_{\text{eff},1\text{ km}}(10, 8)$ and $D_{\text{eff},1\text{ km}}(12, 8)$) are obtained for each particle shape (aggregates, plates and solid columns) considered in the IIR retrieval algorithm. The shape corresponding to the smallest difference between the two $D_{\text{eff},1\text{ km}}$ is selected. For the computation of optical properties, the IIR operational algorithm uses the Yang et al. (2001, 2005) model with a monomodal effective diameter distribution.

The uncertainty of the retrieval algorithm was checked by comparing optical properties retrieved from simulated radiances with the optical properties used as input in the radiative transfer. For this, we perform a 1-D retrieval from 1-D-simulated radiances. The algorithm uncertainties are less than 2% for effective diameters retrieved with the SWT (test not shown here) and 4% for effective optical thickness retrieved with the algorithm similar to the IIR operational algorithm (test not shown here).

3 Impact of cirrus heterogeneities on the retrieved parameters

In this section, we present the heterogeneity effects on the retrieved products at the 1 km IIR spatial resolution as a

function of different cloud optical properties (i.e., optical thickness, effective diameters, extinction coefficients, single-scattering albedo and asymmetry factor) and microphysical (IWC) properties, cirrus-top altitude and geometrical thickness. The heterogeneity effects on the retrieved parameters are assessed by using the difference between products retrieved from modeled 3-D ($\text{BT}_{1\text{ km}}^{3\text{-D}}$) and 1-D ($\text{BT}_{1\text{ km}}^{1\text{-D}}$) 1 km brightness temperatures.

In order to estimate the heterogeneity effects on the retrieved cloud products, we define the following errors due to cloud heterogeneities:

$$\Delta\varepsilon_{\text{eff}} = \varepsilon_{\text{eff}}^{3\text{-D}} - \varepsilon_{\text{eff}}^{1\text{-D}} \quad (5)$$

for effective emissivities calculated by the Eq. (2);

$$\Delta\tau_{\text{eff}} = \tau_{\text{eff}}^{3\text{-D}} - \tau_{\text{eff}}^{1\text{-D}} \quad (6)$$

for effective optical thicknesses calculated by Eq. (3);

$$\begin{aligned} \Delta\text{MI}^{12/8} &= \text{MI}^{3\text{-D},12/8} - \text{MI}^{1\text{-D},12/8} \text{ and} \\ \Delta\text{MI}^{12/10} &= \text{MI}^{3\text{-D},12/10} - \text{MI}^{1\text{-D},12/10} \end{aligned} \quad (7)$$

for microphysical indices calculated from Eq. (4);

$$\Delta D_{\text{eff},1\text{ km}} = D_{\text{eff},1\text{ km}}^{3\text{-D}} - D_{\text{eff},1\text{ km}}^{1\text{-D}} \quad (8)$$

for ice crystal effective diameter retrieved with the SWT.

The “3-D” exponent corresponds to optical properties retrieved from $\text{BT}_{1\text{ km}}^{3\text{-D}}$ and the “1-D” exponent corresponds to those retrieved from $\text{BT}_{1\text{ km}}^{1\text{-D}}$. $D_{\text{eff},1\text{ km}}^{1\text{-D}}$ corresponds either to the effective diameter used in the radiative transfer simulation when it is known (cirrus cases 1 to 8 and CII-3) or to the effective diameter retrieved from $\text{BT}_{1\text{ km}}^{1\text{-D}}$ when the ice crystal effective diameters used in the radiative transfer simulation are unknown (cirrus cases CII-1 and CII-2).

Heterogeneity impacts due to the optical thickness variability are discussed in Sect. 3.1 and those due to optical and microphysical property variabilities in Sect. 3.2.

3.1 Heterogeneity impacts due to the optical thickness variability

Fauchez et al. (2014) show that $\text{BT}_{1\text{ km}}^{3\text{-D}}$ are larger than $\text{BT}_{1\text{ km}}^{1\text{-D}}$ and that their difference is well correlated with the standard deviation of the optical thickness inside the 1 km \times 1 km observation pixel $\sigma_{\tau_{1\text{ km}}}$. This brightness temperature difference is due to the plane-parallel approximation caused by the non-linearity of the relationship between brightness temperature and optical thickness. The impact of the PPA bias ($|\text{BT}_{1\text{ km}}^{3\text{-D}} - \text{BT}_{1\text{ km}}^{1\text{-D}}|$) is greater (in absolute value) for highly absorbing bands because the increase of cloud absorption leads to a larger brightness temperature contrast between the cirrus top and the clear sky atmosphere and, thus, to a stronger averaging effect. Figure 3 illustrates how cirrus heterogeneities affect the retrieval of the effective

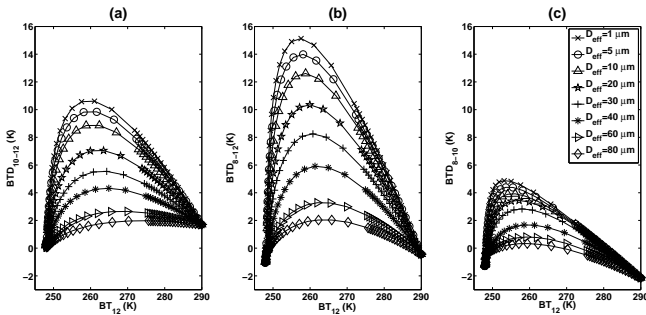


Figure 2. Brightness temperatures differences (BTD) as a function of the 12.05 μm brightness temperature (BT_{12}) for eight effective diameters (D_{eff}) and different optical thickness between 0 and 50 at 12.05 μm : (a) BTD_{10-12} between 10.60 and 12.05 μm channels, (b) BTD_{8-12} between 8.65 and 12.05 μm channels and (c) BTD_{8-10} between 8.65 and 10.60 μm channels.

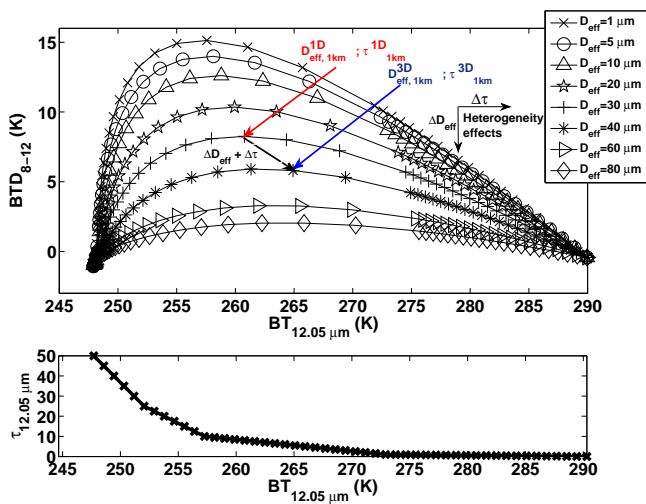


Figure 3. Top panel: brightness temperature differences between 8.65 and 12.05 μm (BTD_{12-8}) as a function of the brightness temperature at 12.05 μm (BT_{12}). The red arrow shows an example of effective diameter, $D_{\text{eff}, 1 \text{ km}}^{1-D}$, and optical thickness, $\tau_{1 \text{ km}}^{1-D}$, retrieved in 1-D without heterogeneity effects, and the blue arrow shows the corresponding effective diameter, $D_{\text{eff}, 1 \text{ km}}^{3-D}$, and optical thickness, $\tau_{1 \text{ km}}^{3-D}$, retrieved with heterogeneity effects. Each point of the arches corresponds to an optical thickness represented in the bottom panel, with $\tau_{12.05 \mu\text{m}}$ as the optical thickness at 12.05 μm . Using the plane-parallel approximation (PPA) leads to an overestimation of the effective diameter and to underestimation of the optical thickness, respectively, compared to a 3-D retrieval.

diameter and the optical thickness. The tip of the red arrow represents the BTD and BT values obtained with a homogeneous cloud with $D_{\text{eff}, 1 \text{ km}}^{1-D}$ and $\tau_{1 \text{ km}}^{1-D}$. Using 3-D radiative transfer inside a heterogeneous cloud with the same mean properties, we obtained the BTD and BT values represented by the tip of the blue arrow in Fig. 3. As heterogeneity effects are larger at the 12.05 μm channel than at the 8.65 μm

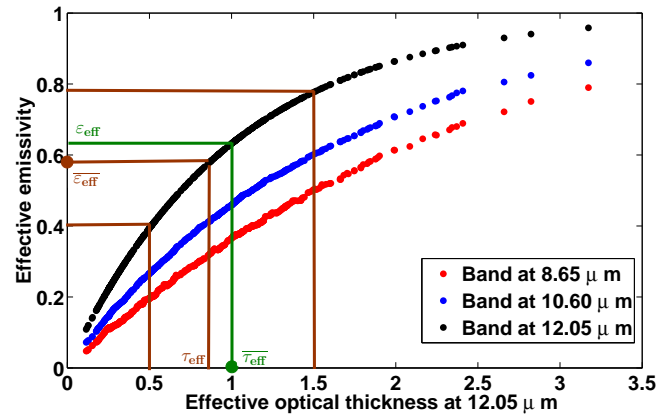


Figure 4. Variation of the effective emissivity as a function of the effective optical thickness at 12.05 μm , estimated in 1-D at the spatial resolution of 100 m for the three IIR channels and for cloudy pixels belonging to cirrus cases 1 to 5. τ_{eff} represents the effective optical thickness corresponding to the averaged effective emissivity $\bar{\varepsilon}_{\text{eff}}$; $\bar{\tau}_{\text{eff}}$ represents the averaged effective optical thickness and ε_{eff} is its corresponding effective emissivity. Brown and green lines show the effective emissivity and effective optical thickness values on the x axis and y axis, respectively, corresponding to a particular point on, for instance, the black arch. The mathematical formulation of the PPA is expressed by the Jensen inequality $\bar{\varepsilon}_{\text{eff}} < \varepsilon_{\text{eff}}(\bar{\tau}_{\text{eff}})$.

channel, brightness temperature differences (BTD_{8-12}), first simulated at the 100 m spatial resolution and then averaged to the 1 km IIR spatial resolution, are smaller than those retrieved from radiances directly simulated at 1 km spatial resolution. Consequently, as effective diameters increase with the decrease of BTD, the retrieved $D_{\text{eff}, 1 \text{ km}}^{3-D}$ is larger than the mean value $D_{\text{eff}, 1 \text{ km}}^{1-D}$ and the retrieved optical thicknesses $\tau_{1 \text{ km}}^{3-D}$ is smaller than the mean optical thickness $\tau_{1 \text{ km}}^{1-D}$.

In addition, Fig. 4 shows the effective emissivity as a function of the effective optical thickness estimated at 100 m spatial resolution. The relationship between effective emissivities and effective optical thickness is nonlinear, as it is between brightness temperatures and optical thickness. Because of the PPA bias, the average effective emissivity is smaller than the effective emissivity of the average of the effective optical thickness $\bar{\tau}_{\text{eff}}$. Similar to brightness temperatures, effective emissivities and effective optical thickness retrieved from radiances, first simulated at 100 m spatial resolution of and then averaged to the IIR spatial resolution of 1 km, are smaller than those retrieved from radiances directly simulated at the spatial resolution of 1 km.

Figure 5 presents $\Delta\varepsilon_{\text{eff}}$ (a, b and c) and $\Delta\tau_{\text{eff}}$ (d, e and f) as a function of the standard deviation of the optical thickness inside the 1 km \times 1 km observation pixel ($\sigma_{\tau_{1 \text{ km}}}$) for cirrus cases 1 to 5 and for 8.65, 10.60 and 12.05 μm channels, respectively. We notice, first of all, that $\Delta\varepsilon_{\text{eff}}$ and $\Delta\tau_{\text{eff}}$ are correlated with $\sigma_{\tau_{1 \text{ km}}}$ at more than 94 % except for cirrus case 1 at 8.65 μm , where the horizontal transport smooths the slight

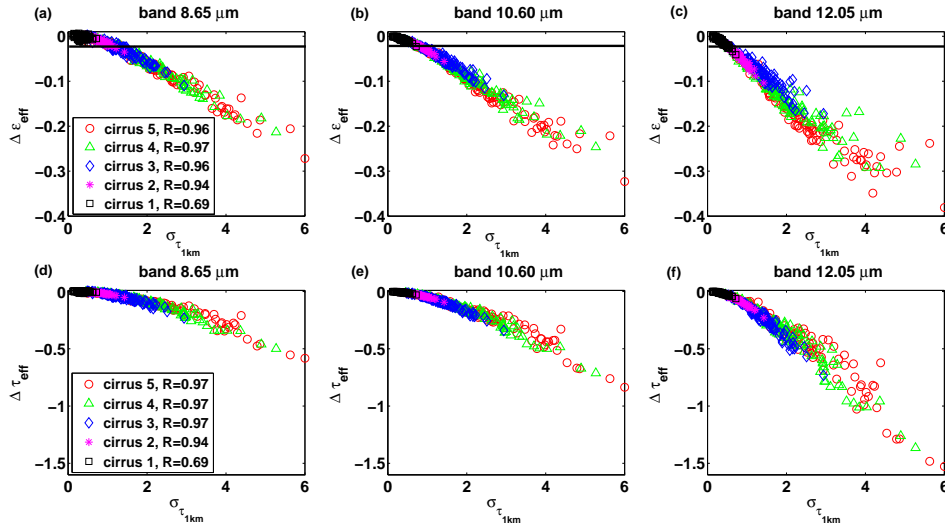


Figure 5. Errors on the effective emissivity $\Delta \varepsilon_{\text{eff}}$ (a, b and c) and on the effective optical thickness $\Delta \tau_{\text{eff}}$ (d, e and f) at 8.65, 10.60 and 12.05 μm , respectively, as a function of the optical thickness standard deviation, $\sigma_{\tau_{1\text{km}}}$, for cirrus cases 1 ($\tau_c = 0.45$, $\rho_\tau = 0.7$), 2 ($\tau_c = 0.90$, $\rho_\tau = 0.7$), 3 ($\tau_c = 1.80$, $\rho_\tau = 0.7$), 4 ($\tau_c = 1.80$, $\rho_\tau = 1.1$) and 5 ($\tau_c = 1.80$, $\rho_\tau = 1.5$) with $D_{\text{eff}} = 9.95 \mu\text{m}$ for the five cirrus. The black lines correspond to the IIR operational algorithm uncertainty on the effective emissivity.

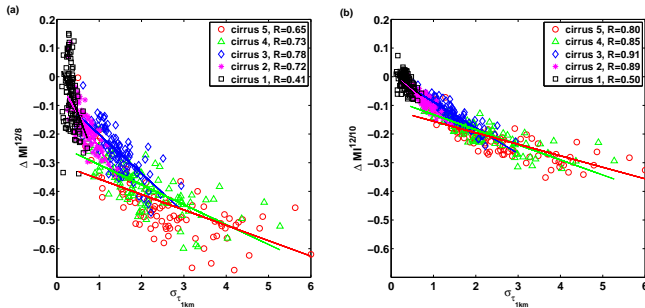


Figure 6. Microphysical index differences $\Delta \text{MI}^{12/8}$ (a) and $\Delta \text{MI}^{12/10}$ (b) as a function of the standard deviation of the optical thickness, $\sigma_{\tau_{1\text{km}}}$, for cirrus cases 1 ($\tau_c = 0.45$, $\rho_\tau = 0.7$), 2 ($\tau_c = 0.90$, $\rho_\tau = 0.7$), 3 ($\tau_c = 1.80$, $\rho_\tau = 0.7$), 4 ($\tau_c = 1.80$, $\rho_\tau = 1.1$) and 5 ($\tau_c = 1.80$, $\rho_\tau = 1.5$) with $D_{\text{eff}} = 9.95 \mu\text{m}$ for the five cirrus. R represents the correlation coefficient between ΔMI and $\sigma_{\tau_{1\text{km}}}$.

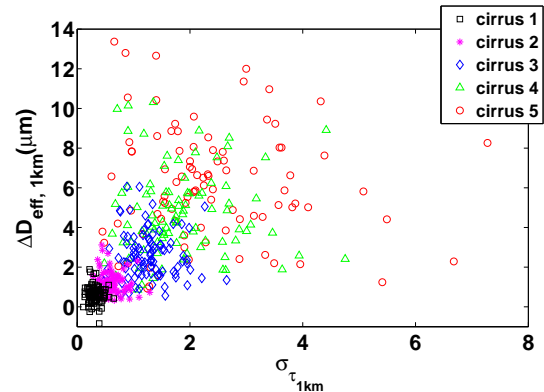


Figure 7. Errors on the retrieved effective diameter $\Delta D_{\text{eff}, 1\text{km}}$ as a function of the standard deviation of the optical thickness, $\sigma_{\tau_{1\text{km}}}$, for cirrus cases 1 ($\tau_c = 0.45$, $\rho_\tau = 0.7$), 2 ($\tau_c = 0.90$, $\rho_\tau = 0.7$), 3 ($\tau_c = 1.80$, $\rho_\tau = 0.7$), 4 ($\tau_c = 1.80$, $\rho_\tau = 1.1$) and 5 ($\tau_c = 1.80$, $\rho_\tau = 1.5$) with $D_{\text{eff}} = 9.95 \mu\text{m}$ for the five cirrus. Effective diameters are estimated using the split-window technique.

heterogeneity of the radiative field. $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ are negative, meaning that the 3-D effective emissivities and effective optical thickness are smaller than those in 1-D. Indeed, as explained by Fauchez et al. (2014), heterogeneity effects lead to an increase of radiances or brightness temperatures. As radiances decrease with the cloud extinction, larger radiances lead then to smaller cloud effective emissivity and effective optical thickness. In addition, $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ are shown to depend on the wavelength. For example, at $\sigma_{\tau_{1\text{km}}} = 1$, $\Delta \varepsilon_{\text{eff}}$ is equal to -0.01 at 8.65 μm , -0.03 at 10.60 μm and -0.05 at 12.05 μm . This is due to the increase of absorption from 8.65 to 12.05 μm that leads to an increase of the contrast between cloud and clear sky pixels, and thus to an increase of the PPA bias. For comparison, Garnier et al. (2012) have shown that

the effective emissivity error due to the retrieval method is about 0.03 for the 12.05 μm band (black lines of the Fig. 5) assuming a 1 K clear sky atmosphere temperature uncertainty for an ocean scene. This uncertainty is smaller than the average error due to cloud heterogeneity $\overline{\Delta \varepsilon_{\text{eff}}}$. We can note that, at $\sigma_{\tau_{1\text{km}}} \sim 1$, $\Delta \varepsilon_{\text{eff}}$ is equal to or larger than 0.03 for the 10.60 and 12.05 μm bands. $\sigma_{\tau_{1\text{km}}} \sim 1$ corresponds also to the limit where the heterogeneity effects on brightness temperatures become larger than the IIR instrumental accuracy of 1 K (Fauchez et al., 2014).

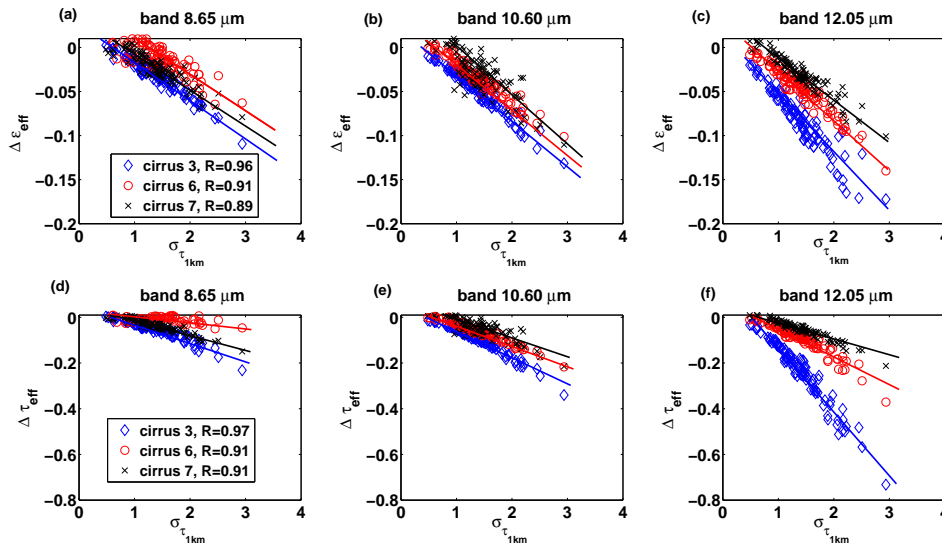


Figure 8. Errors on the effective emissivity $\Delta \varepsilon_{\text{eff}}$ (a, b and c) and on the effective optical thickness $\Delta \tau_{\text{eff}}$ (d, e and f) at 8.65, 10.60 and 12.05 μm , respectively, as a function of the optical thickness standard deviation, $\sigma_{\tau_{1\text{km}}}$, for three identical cirrus fields but for different ice crystal effective diameters: cirrus cases 3 ($D_{\text{eff}} = 9.95 \mu\text{m}$), 6 ($D_{\text{eff}} = 20.09 \mu\text{m}$) and 7 ($D_{\text{eff}} = 40.58 \mu\text{m}$), with $\tau_c = 1.80$ and $\rho_\tau = 0.7$ for the three cirrus. R represents the correlation coefficient between $\Delta \varepsilon_{\text{eff}}$ (a, b and c) and $\sigma_{\tau_{1\text{km}}}$ and between $\Delta \tau_{\text{eff}}$ (d, e and f) and $\sigma_{\tau_{1\text{km}}}$.

Figure 6a and b show the error on the microphysical indices $\Delta \text{MI}^{12/8}$ and $\Delta \text{MI}^{12/10}$, respectively, as a function of $\sigma_{\tau_{1\text{km}}}$ for cirrus cases 1 to 5. First note that the errors on the two microphysical indices are on average negative, except again for cirrus case 1 for $\Delta \text{MI}^{12/8}$, and they increase with the cirrus mean optical thickness (from cirrus cases 1 to 3) and heterogeneity parameter (from cirrus cases 3 to 5). The correlation with $\sigma_{\tau_{1\text{km}}}$ is better for $\Delta \text{MI}^{12/10}$ than for $\Delta \text{MI}^{12/8}$. Again, the strongest scattering in the band at 8.65 μm tends to smooth the radiative field heterogeneities and, therefore, to degrade the correlation between $\Delta \text{MI}^{12/8}$ and $\sigma_{\tau_{1\text{km}}}$. $\Delta \text{MI}^{12/8}$ is, on average, larger than $\Delta \text{MI}^{12/10}$ because the difference of absorption and effective emissivity is significantly greater for the 12.05 μm / 8.65 μm pair than for 12.05 μm / 10.60 μm . Effective diameters of ice crystals are estimated from the microphysical indices using a LUT and are thus also impacted by heterogeneity effects. As $\Delta \text{MI} = \text{MI}^{3\text{-D}} - \text{MI}^{1\text{-D}}$ is negative, the impact of cloud heterogeneities leads to an underestimation of the microphysical indices. This underestimation leads then to an overestimation of the retrieved effective diameters (smaller microphysical indices correspond to larger effective diameters).

Using the SWT, we are also able to simulate the impact of cirrus heterogeneities on the retrieved effective diameters of ice crystals. In Fig. 7, we plot the error on the effective diameter error $\Delta D_{\text{eff}, 1\text{km}}$, due to heterogeneities, as a function of $\sigma_{\tau_{1\text{km}}}$ for cirrus cases 1 to 5. We see that $\Delta D_{\text{eff}, 1\text{km}}$ is positive and generally increases with the cirrus mean optical thickness (from cirrus cases 1 to 3) and the heterogeneity parameter ρ_τ (from cirrus cases 3 to 5). Indeed, $\sigma_{\tau_{1\text{km}}}$ generally increases with τ_c and ρ_τ , as expected.

Figure 8 is the same as Fig. 5, except for different effective diameters: $D_{\text{eff}} = 9.95 \mu\text{m}$, $D_{\text{eff}} = 20.09 \mu\text{m}$ and $D_{\text{eff}} = 40.58 \mu\text{m}$ (cirrus cases 3, 6 and 7, respectively). Here $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ decrease with increasing D_{eff} (except at the 8.65 μm band where ω_0 increases between $D_{\text{eff}} = 9.95$ and 20.09 μm). Indeed, ω_0 increases with D_{eff} (except at the 8.65 μm band) and leads to a decrease of the absorption and, thus, of the PPA bias. The impact of the effective diameter on $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ is particularly marked for the 12.05 μm band where the absorption of ice crystals decreases strongly between $D_{\text{eff}} = 9.95 \mu\text{m}$ and $D_{\text{eff}} = 40.58 \mu\text{m}$ (cirrus cases 3 and 7, respectively).

In addition, we estimated the heterogeneity effects on the retrieved ice crystal effective diameters ($\Delta D_{\text{eff}, 1\text{km}}$) for the three D_{eff} . On average, $\Delta D_{\text{eff}, 1\text{km}} \sim +3 \mu\text{m}$ for cirrus case 3 ($D_{\text{eff}} = 9.95 \mu\text{m}$) and 6 ($D_{\text{eff}} = 20.09 \mu\text{m}$). Thus, there is no a significant increase of heterogeneity effects on retrieved effective diameters between these two effective sizes. For cirrus case 7 ($D_{\text{eff}} = 40.58 \mu\text{m}$), there is no real effect ($\Delta D_{\text{eff}, 1\text{km}} \sim \pm 0 \mu\text{m}$) due to the saturation of the SWT. Indeed, as noted above, effective diameters close to 40 μm lead to weak brightness temperature differences. This is illustrated in Fig. 2, where the amplitude of arches, and thus the sensitivity, decreases with the increase of the effective diameter.

3.2 Heterogeneity effects due to optical and microphysical property variabilities

As presented in Sect. 2.2, we use the parametrization developed by Baran et al. (2009, 2013) and Baran (2012) to sim-

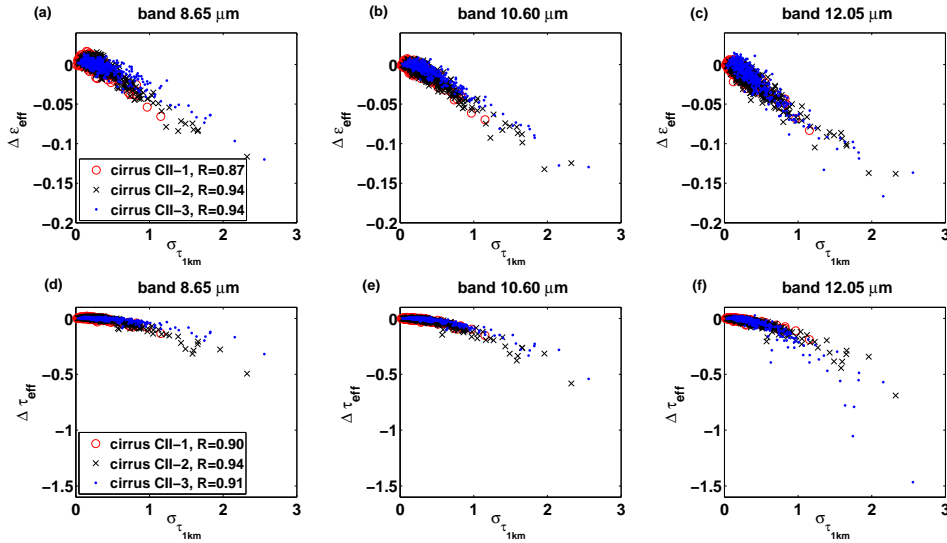


Figure 9. Errors on the effective emissivity $\Delta \varepsilon_{\text{eff}}$ (a, b and c) and on the effective optical thickness $\Delta \tau_{\text{eff}}$ (d, e and f) at 8.65, 10.60 and 12.05 μm , respectively, as a function of the optical thickness standard deviation, $\sigma_{\tau_{1\text{km}}}$, for cirrus cases CII-1, CII-2 and CII-3.

ulate a 3-D heterogeneous cloud optical property field from the 3-D distribution of the IWC and temperature. IWC values measured during the CIRCLE-2 campaign are coupled with a MLS temperature profile to generate a realistic 3-D optical property field for simulating of the CII-1 and CII-2 cirrus cases. In addition, to compare with the previous cirrus cases, the cirrus case CII-3 was generated from the CIRCLE-2 cloud field using optical properties identical to cirrus case 8.

Figure 9 shows the impact of cirrus heterogeneities on the retrieved effective emissivity and on the effective optical thickness as a function of the standard deviation of the optical thickness, $\sigma_{\tau_{1\text{km}}}$, for cirrus cases CII-1, CII-2 and CII-3. $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ are similar for the three cirrus cases, although some slight differences are evident as a function of the wavelength. Indeed, at 8.65 μm , $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ are smaller for the CII-3 cirrus case than for the two others cirrus cases. At 10.60 μm , this difference is close to 0. At 12.05 μm , $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$ are larger for the CII-3 cirrus than for CII-1 and CII-2 cirrus. This effect is due to the variability of the optical properties for the CII-1 and CII-2 cirrus. Indeed, cirrus case CII-3 contains only aggregate crystals of effective diameter $D_{\text{eff}} = 9.95 \mu\text{m}$ resulting from the model of Yang et al. (2001, 2005), while cirrus cases CII-1 and CII-2 contain crystal of various sizes. For CII-3 cirrus, small crystals have a single-scattering albedo maximum at 8.65 μm , leading to a lower PPA bias. At 12.05 μm , small particles are more absorbing and the PPA bias is larger. For the CII-1 cirrus, corresponding to the cirrus observed during the CIRCLE-2 campaign, the average effective emissivity error is within the limit of the method sensibility (Garnier et al., 2012) of about 0.03 in absolute value.

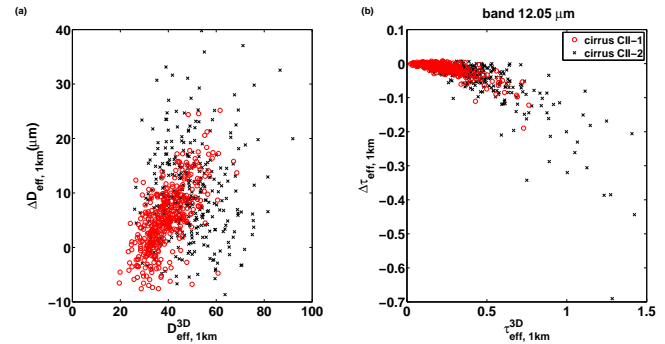


Figure 10. (a) Errors on the retrieved effective diameter $\Delta D_{\text{eff}, 1\text{km}}$ as a function of the effective diameter, $D_{\text{eff}, 1\text{km}}^{3\text{D}}$, and (b) error on the effective optical thickness, $\Delta \tau_{\text{eff}, 1\text{km}}$, as a function of the effective optical thickness, $\tau_{\text{eff}, 1\text{km}}^{3\text{D}}$, for cirrus CII-1 and CII-2.

To study heterogeneity effects on the retrieved ice crystals' effective diameters for the CII-1 and CII-2 cirrus, we compare effective diameters $D_{\text{eff}, 1\text{km}}^{3\text{D}}$ and $D_{\text{eff}, 1\text{km}}^{1\text{D}}$ retrieved from $\text{BT}_{1\text{km}}^{3\text{D}}$ and $\text{BT}_{1\text{km}}^{1\text{D}}$, respectively. $D_{\text{eff}, 1\text{km}}^{3\text{D}}$ and $\tau_{\text{eff}, 1\text{km}}^{3\text{D}}$ represent the cloud optical properties resulting from a 3-D radiative transfer simulation through a heterogeneous atmosphere ($\text{BT}_{1\text{km}}^{3\text{D}}$). The differences $\Delta D_{\text{eff}, 1\text{km}} = D_{\text{eff}, 1\text{km}}^{3\text{D}} - D_{\text{eff}, 1\text{km}}^{1\text{D}}$ and $\Delta \tau_{\text{eff}, 1\text{km}} = \tau_{\text{eff}, 1\text{km}}^{3\text{D}} - \tau_{\text{eff}, 1\text{km}}^{1\text{D}}$ correspond, therefore, to the heterogeneity effects on the retrieval of $D_{\text{eff}, 1\text{km}}^{3\text{D}}$ and $\tau_{\text{eff}, 1\text{km}}^{3\text{D}}$. For these two cirrus, the optical properties are heterogeneous. Therefore, Fig. 10a shows $\Delta D_{\text{eff}, 1\text{km}}$ as a function of $D_{\text{eff}, 1\text{km}}^{3\text{D}}$ and Fig. 10b shows $\Delta \tau_{\text{eff}, 1\text{km}}$ as a function of $\tau_{\text{eff}, 1\text{km}}^{3\text{D}}$. We see that $\Delta D_{\text{eff}, 1\text{km}}$ and $\Delta \tau_{\text{eff}, 1\text{km}}$ increase, in absolute value, with

Table 2. $\overline{D_{\text{eff,IIR}}}$ and $\overline{\tau_{\text{IIR}}}$: averaged effective diameter and optical thickness, respectively, retrieved by IIR on 25 May 2007 during the CIRCLE-2 campaign; $\overline{D_{\text{eff},1\text{ km}}^{3\text{-D}}}$ and $\overline{\tau_{1\text{ km}}^{3\text{-D}}}$: averaged effective diameter and optical thickness, respectively, retrieved for CII-1 and CII-2 cirrus and $\Delta D_{\text{eff},1\text{ km}}^{3\text{-D}}$ and $\Delta \tau_{\text{eff},1\text{ km}}^{3\text{-D}}$: averaged errors on the effective diameter and optical thickness, respectively, due to cloud heterogeneities in absolute value and in percentage.

Cirrus	$\overline{D_{\text{eff,IIR}}} (\mu\text{m})$	$\overline{D_{\text{eff},1\text{ km}}^{3\text{-D}}} (\mu\text{m})$	$\overline{\Delta D_{\text{eff},1\text{ km}}^{3\text{-D}}} (\mu\text{m})$	$\overline{\tau_{\text{IIR}}}$	$\overline{\tau_{\text{eff},1\text{ km}}^{3\text{-D}}}$	$\overline{\Delta \tau_{\text{eff},1\text{ km}}^{3\text{-D}}}$
CII-1	44.2	38.9	5.1 (13 %)	0.41	0.40	-0.02 (-5 %)
CII-2	-	48.7	9.7 (20 %)	-	0.74	-0.05 (-7 %)

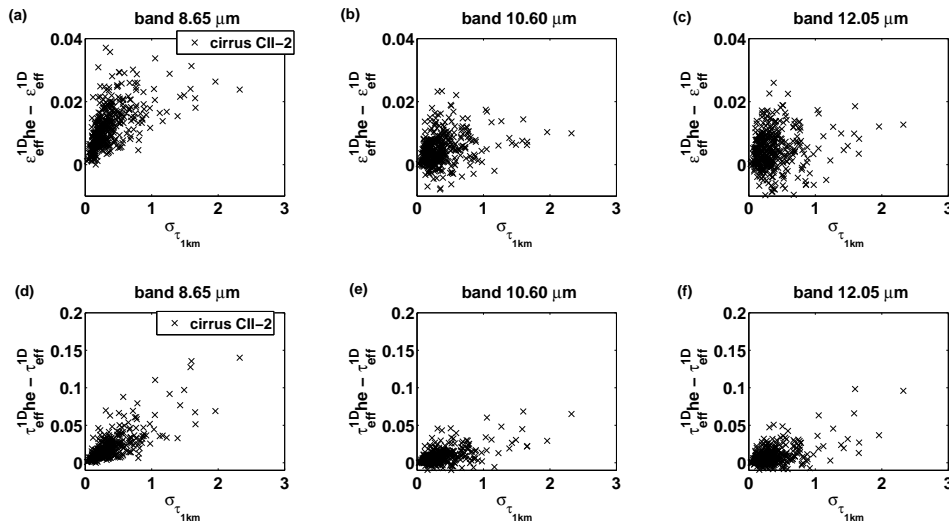


Figure 11. Effective emissivity differences (a, b, c) between $\varepsilon_{\text{eff}}^{1\text{-D}}_{\text{he}}$ and $\varepsilon_{\text{eff}}^{1\text{-D}}$ and effective optical thickness differences (d, e, f) between $\tau_{\text{eff}}^{1\text{-D}}_{\text{he}}$ and $\tau_{\text{eff}}^{1\text{-D}}$ retrieved from radiances calculated in the case of vertically heterogeneous and vertically homogeneous cloudy columns, respectively, as a function of the standard deviation of the optical thickness $\sigma_{\tau_{1\text{ km}}}$ for cirrus CII-2 for bands at 8.65, 10.60 and 12.05 μm , respectively.

$D_{\text{eff},1\text{ km}}^{3\text{-D}}$ and $\tau_{\text{eff},1\text{ km}}^{3\text{-D}}$, respectively. Table 2 summarizes the optical properties retrieved by IIR during CIRCLE-2 and those retrieved from our simulations as well as the estimated heterogeneity effects. First of all, for the CII-1 cirrus possessing the characteristics of the cirrus observed during the CIRCLE-2 campaign, the average value of the retrieved effective diameter ($\overline{D_{\text{eff},1\text{ km}}^{3\text{-D}}} \sim 38.9 \mu\text{m}$) and the mean effective optical thickness ($\overline{\tau_{\text{eff},1\text{ km}}^{3\text{-D}}} \sim 0.40$ at 12.05 μm) are close to those retrieved from the IIR measurements along the CALIOP/CALIPSO track ($\overline{D_{\text{eff,IIR}}} = 44.2 \mu\text{m}$ and $\overline{\tau_{\text{IIR}}} = 0.41$ without underlying liquid water cloud). Thus, there is a good agreement between optical properties retrieved by the IIR operational algorithm during the CIRCLE-2 campaign and those retrieved with our simulations. The mean error due to heterogeneity effects is approximately 5.1 μm (13 %) for retrieved effective diameter and approximately -0.02 (5 %) for effective optical thickness. On average, these relative errors due to heterogeneity effects are, thus, weak compared to the uncertainty estimate of Dubuisson et al. (2008) for the IIR retrieval (10 to 25 % for $D_{\text{eff},1\text{ km}}$ and 10 % for $\tau_{\text{eff},1\text{ km}}$).

However, at the observation pixel scale, some values can reach more than 40 % for effective diameter and 15 % for effective optical thickness, which is quite significant. Furthermore, errors due to cloud heterogeneities increase with the IWC or the cirrus mean optical thickness; the cirrus CII-2 case, for instance, with IWC twice as large as cirrus CII-1, has $\overline{\Delta D_{\text{eff},1\text{ km}}^{3\text{-D}}} \sim 9.7 \mu\text{m}$ (20 %) and $\overline{\Delta \tau_{\text{eff},1\text{ km}}^{3\text{-D}}} \sim -0.05$ (7 %).

3.3 Influence of the vertical variability of optical properties

To find the influence of the vertical variability of cirrus optical properties (σ_e , ϖ_0 and g) on the retrieval errors, we compare cloud products retrieved from $\text{BT}_{1\text{ km}}^{1\text{-D}}$ with vertically heterogeneous columns with those retrieved for vertically homogeneous columns obtained after a vertical averaging of the IWC, for the CII-2 cirrus case.

Figures 11 shows the effects of the vertical heterogeneity of the optical properties on the effective emissivity (a, b and c) and on the effective optical thickness (d, e and f).

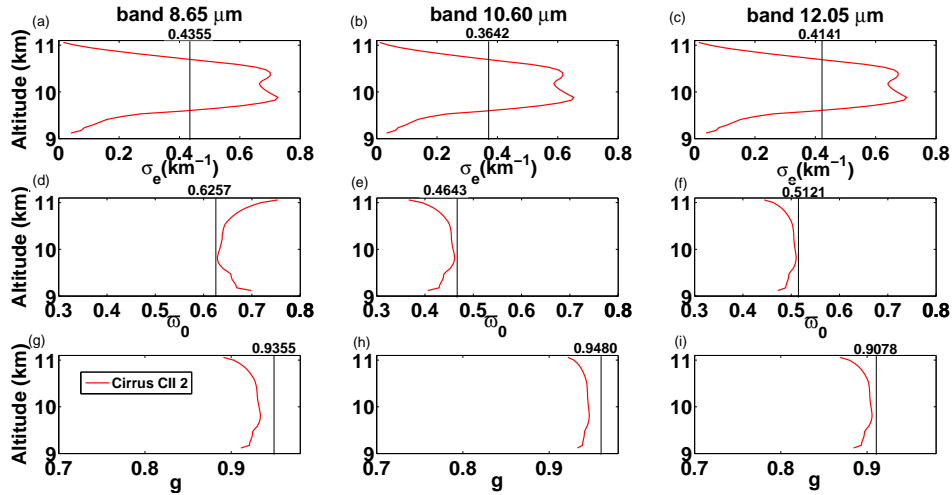


Figure 12. (a), (b) and (c): vertical variation of the mean extinction coefficient σ_e ; (d), (e) and (f): vertical variation of the mean single-scattering albedo ω_0 ; (g), (h) and (i): vertical variation of the asymmetry factor g for the three IIR channels at 8.65, 10.60 μm and 12.05 for cirrus CII-2. Vertical black lines correspond to the mean value of the optical coefficient obtained after vertical averaging of the IWC.

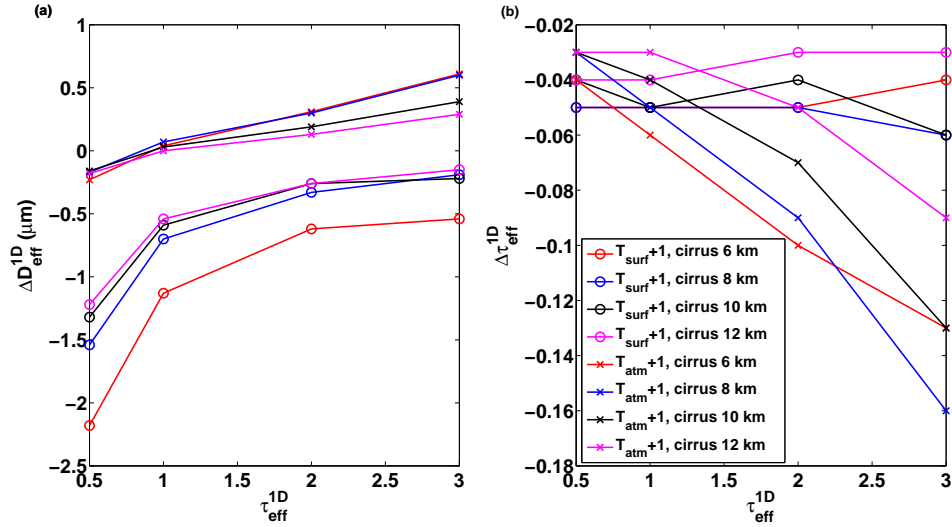


Figure 13. Errors on the (a) retrieved effective diameter, $\Delta D_{\text{eff}}^{1-D}$, and (b) on the effective optical thickness, $\Delta \tau_{\text{eff}}^{1-D}$, as a function of τ_{eff}^{1-D} due to an error of +1 K on the surface temperature ($T_{\text{surf}} + 1$) and +1 K on the atmospheric temperature profile ($T_{\text{atm}} + 1$) for cirrus with top altitudes of 6, 8, 10 and 12 km.

Here $\varepsilon_{\text{eff}}^{1-D}$ and τ_{eff}^{1-D} are estimated from vertically heterogeneous cloudy columns and $\varepsilon_{\text{eff}}^{1-D}$ and τ_{eff}^{1-D} from vertically homogeneous cloudy columns, as a function of the standard deviation of the optical thickness $\sigma_{\tau_{1\text{ km}}}$ for the three IIR channels. Differences between retrieved products estimated from vertically heterogeneous and homogeneous cloudy columns are significantly weaker than those due to 3-D heterogeneities (horizontal and vertical heterogeneities). Furthermore, contrary to the 3-D heterogeneity effects $\Delta \varepsilon_{\text{eff}}$ and $\Delta \tau_{\text{eff}}$, the differences ($\varepsilon_{\text{eff}}^{1-D} - \varepsilon_{\text{eff}}^{1-D}$) and ($\tau_{\text{eff}}^{1-D} - \tau_{\text{eff}}^{1-D}$) are positive. These effects are particular to our simulations, where vertical heterogeneities tend thus to smooth the hor-

izontal heterogeneity effects. These observations can be explained with Fig. 12, which shows the vertical profiles of the optical properties of cirrus CII-2 in the vertically heterogeneous case (red curves) and the vertically homogeneous case (black lines) after vertical averaging of the IWC and temperature using the parametrization of Baran et al. (2009, 2013) and Baran (2012). In this way, values of the vertically homogeneous case are different from the average of the optical coefficients of the vertically heterogeneous case: ω_0 of the vertically homogeneous case is larger than the vertical averaging of the heterogeneous case for the 10.60 and 12.05 μm bands. In addition, the asymmetry parameter g of the verti-

Table 3. Averaged errors on the retrieved cirrus optical properties due to the 3-D cloud heterogeneity for three ice crystal effective diameters (columns 3, 4 and 5); the vertical heterogeneity of optical properties (column 6) with 1Dvhe and 1Dvho representing the -1 -D radiative transfer with vertically heterogeneous and homogeneous columns, respectively; an incertitude of 1 K of the surface temperature (column 7) and the temperature atmospheric profile (column 8); and the IIR retrieval uncertainty (Dubuisson et al., 2008, column 9). $\overline{\Delta D_{\text{eff}, 1 \text{ km}}}$ and $|\overline{\Delta D_{\text{eff}, 1 \text{ km}}}|$ correspond to the absolute error in micrometers and to the relative error in percent, respectively, on the retrieval of the effective diameter; $\overline{\Delta \tau_{1 \text{ km}}}$ and $|\overline{\Delta \tau_{1 \text{ km}}}|$ correspond to the absolute and relative error in percent, respectively, on the retrieval of the optical thickness.

$\sigma_{\tau_{1 \text{ km}}}$	$D_{\text{eff}} (\mu\text{m})$	Horizontal heterogeneity effects as a function of D_{eff}			Vertical heterogeneity (1Dvhe – 1Dvho)	Others uncertainties		IIR uncertainty
		40.58	20.09	9.95		surface 1 K	atmosphere 1 K	
1	$\overline{\Delta D_{\text{eff}, 1 \text{ km}}} (\mu\text{m})$	–0.5	2.0	2.5	2	1	0.2	–
	$ \overline{\Delta D_{\text{eff}, 1 \text{ km}}} (\%)$	~ 1	~ 10	~ 25	~ 4	~ 10	~ 2	~ 10 to ~ 25
	$\overline{\Delta \tau_{\text{eff}, 1 \text{ km}}}$	–0.02	–0.10	–0.20	0.03	0.04	0.08	–
	$ \overline{\Delta \tau_{\text{eff}, 1 \text{ km}}} (\%)$	~ 1	~ 6	~ 12	~ 4	~ 2	~ 4	~ 10
2	$\overline{\Delta D_{\text{eff}, 1 \text{ km}}} (\mu\text{m})$	1	3	3	2	–	–	–
	$ \overline{\Delta D_{\text{eff}, 1 \text{ km}}} (\%)$	~ 3	~ 15	~ 40	~ 4	–	–	~ 10 to ~ 25
	$\overline{\Delta \tau_{\text{eff}, 1 \text{ km}}}$	–0.10	–0.20	–0.50	0.10	–	–	–
	$ \overline{\Delta \tau_{\text{eff}, 1 \text{ km}}} (\%)$	~ 6	~ 12	~ 28	~ 12	–	–	~ 10

cally homogeneous case is larger than the average of the vertically heterogeneous case in the three bands. Consequently, the cirrus is less absorbent in the vertically homogeneous case and thus the effective emissivities and effective optical thicknesses are weaker. This vertical variability of optical properties, for the cirrus CII-2 case, impact the retrieval of the effective diameter of, on average, $4 \mu\text{m}$ (figure not presented here).

Note that effects of the vertical variability are discussed here for the structure of the cirrus observed during the CIRCLE-2 campaign. Effects could be different for other cirrus structure but they are not discussed here. For example, for old cirrus, sedimentation processes could be much larger, increasing differences between the cloud top and base. The impact of the vertical variability on cloud properties retrieved from satellite observations could thus be larger.

4 Other sources of uncertainty

We show in the previous sections that heterogeneity effects can be an important source of errors in the retrieved optical properties. To compare its importance on the retrieved cloud parameters with regard to other possible error sources for IIR measurements, we test the impact of a 1 K uncertainty in the surface temperature and in the atmospheric temperature profile measurements, an error that corresponds to that estimated by Garnier et al. (2012). Figure 13a and b show the error in the retrieved effective diameter ($\Delta D_{\text{eff}}^{1-D}$) and in the retrieved effective optical thickness ($\Delta \tau_{\text{eff}}^{1-D}$), respectively, as a function of τ_{eff}^{1-D} for cirrus with a top altitude of 6, 8, 10 and 12 km. The retrieval of the effective diameter and optical thickness is performed using the SWT on 1-D radiative trans-

fer simulations. We can see that $\Delta D_{\text{eff}}^{1-D}$ are less than $2.5 \mu\text{m}$ (25 %) and $\Delta \tau_{\text{eff}}^{1-D}$ less than 0.16 (5 %). By comparison, these errors are in the IIR retrieval uncertainty of 10–25 % for D_{eff} and about 10 % for τ_{eff} (Dubuisson et al., 2008). In addition, they are significantly smaller than those due to cloud heterogeneity effects (more than 50 % for D_{eff} and 10 to 15 % for τ_{eff}). In Fig. 13a, it is evident that increasing cirrus optical thickness or cloud-top altitude decreases the effective diameter retrieval error due to a 1 K surface-temperature uncertainty. In Fig. 13b, the effective diameter retrieval error due to atmospheric temperature profile uncertainty increases with increasing optical thickness because cloud emissivity also increases.

Fauchez et al. (2014) show that cloud-top altitude and geometrical thickness significantly influence the heterogeneity effects because the brightness temperature contrast between the surface and the cloud top increases with increasing cloud-top altitude and decreases with increasing vertical extension for a constant cloud top (as the cloud base is closer to the surface). For retrieved cloud products estimated with an algorithm similar to the IIR operational algorithm, the effective emissivity is independent of the cloud altitude and geometrical thickness; thus the impacts of altitude and geometrical thickness on the retrieval are weak.

5 Summary and conclusions

In this paper, we discussed the impact of cirrus heterogeneity effects in the retrieval of cloud parameters from thermal infrared radiometric measurements from space. We have focused on the IIR radiometer for which the operational algorithm estimates the cirrus effective emissivity, the effec-

tive optical thickness and the ice crystal effective diameter of the observation pixel. We show that errors due to the cirrus heterogeneity effects on the effective emissivity and the effective optical thickness are well correlated to the sub-pixel optical thickness standard deviation $\sigma_{\tau_{1 \text{ km}}}$ and generally increase with increasing optical thickness $\tau_{1 \text{ km}}$. These errors are greater than the precision of the retrieval method ($\Delta \varepsilon_{\text{eff}} \sim 0.03$) for $\sigma_{\tau_{1 \text{ km}}} \sim 1$, corresponding also to the value from which heterogeneity effects on brightness temperatures become larger than the IIR instrumental accuracy of 1 K (Fauchez et al., 2014).

Our results are summarized in Table 3. Heterogeneity effects for three effective diameters are compared with the retrieval errors caused by the vertical inhomogeneity of optical properties and with the impact of an error of 1 K, corresponding to the IIR accuracy, on the atmospheric temperature profile and on the surface temperature. Results are shown for pixels with $\sigma_{\tau_{1 \text{ km}}} = 1$ (medium heterogeneity) and 2 (large heterogeneity). The most important errors in the cloud optical property retrieval concern those due to the subpixel heterogeneity of the optical thickness, in particular for the smallest crystals ($\Delta D_{\text{eff}, 1 \text{ km}} = 2.5 \mu\text{m}$ ($\sim 25\%$) and $\Delta \tau_{\text{eff}, 1 \text{ km}} = -0.20$ ($\sim 12\%$) for $\sigma_{\tau_{1 \text{ km}}} = 1$). Indeed, the absorption is larger for small crystals and, thus, the PPA bias is greater. For $D_{\text{eff}} = 40.58 \mu\text{m}$, the ice crystal optical properties in the three IIR channels converge to similar values leading to smaller brightness temperature differences between channels and, thus, to a decrease of the retrieval accuracy. Errors due to the vertical inhomogeneity of the optical properties, an error of 1 K on the surface temperature or atmospheric temperature profile are smaller than the IIR retrieval errors (Dubuisson et al., 2008). Thus, the influence of these parameters appears negligible compared to optical thickness heterogeneity and IIR retrieval uncertainty.

The impacts of cirrus heterogeneities on the retrieved cloud parameters studied in this paper are for a 1 km spatial resolution. These biases could decrease with an increase of the spatial resolution although photon transport effects would increase. Fauchez et al. (2014) estimate that a 250 m spatial resolution could significantly reduce the PPA bias while photon transport effects remain weak. However, heterogeneity effects on the retrieved cloud products at this resolution require further investigation. This study also provides ways to potentially correct the heterogeneity errors using the subpixel measurements to estimate $\sigma_{\tau_{1 \text{ km}}}$. Furthermore, differences between heterogeneity effects in the visible/near-infrared and thermal infrared ranges for different spatial resolutions also require further investigation to estimate their impact on cloud products retrieved using a combination of the visible and near-infrared/short-wave infrared and infrared retrieval methods, as proposed by Cooper et al. (2007).

Acknowledgements. The authors acknowledge the Centre National de la Recherche Scientifique, the Programme National de Télé-détection Spatiale and the Direction Générale de l'Armement for their financial support. We also thank the use of resources provided by the European Grid Infrastructure. For more information, please refer to the EGI-InSPIRE paper (<http://go.egi.eu/pdnon>).

We thank Anthony Baran for the numerous fruitful discussions on the cirrus optical properties.

We also thank the two reviewers whose remarks largely improved the quality of this article.

Edited by: A. Macke

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