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Investigating the links between ozone and organic aerosol chemistry in a biomass burning plume from a prescribed fire in California chaparral

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Abstract. Within minutes after emission, complex photochemistry in biomass burning smoke plumes can cause large changes in the concentrations of ozone (O_3) and organic aerosol (OA). Being able to understand and simulate this rapid chemical evolution under a wide variety of conditions is a critical part of forecasting the impact of these fires on air quality, atmospheric composition, and climate. Here we use version 2.1 of the Aerosol Simulation Program (ASP) to simulate the evolution of O_3 and secondary organic aerosol (SOA) within a young biomass burning smoke plume from the Williams prescribed fire in chaparral, which was sampled over California in November 2009. We demonstrate the use of a method for simultaneously accounting for the impact of the unidentified intermediate volatility, semi-volatile, and extremely low volatility organic compounds (here collectively called "SVOCs") on the formation of OA (using the Volatility Basis Set – VBS) and O_3 (using the concept of mechanistic reactivity). We show that this method can successfully simulate the observations of O_3, OA, NO_x, ethylene (C_2H_4), and OH to within measurement uncertainty using reasonable assumptions about the average chemistry of the unidentified SVOCs. These assumptions were (1) a reaction rate constant with OH of $\sim 10^{-11}$ cm^3 s^{-1}; (2) a significant fraction (up to $\sim 50\%$) of the $RO_2 + NO$ reaction resulted in fragmentation, rather than functionalization, of the parent SVOC; (3) ~ 1.1 molecules of O_3 were formed for every molecule of SVOC that reacted; (4) $\sim 60\%$ of the OH that reacted with the unidentified non-methane organic compounds (NMOC) was regenerated as HO_2; and (5) that $\sim 50\%$ of the NO that reacted with the SVOC peroxy radicals was lost, presumably to organic nitrate formation. Additional evidence for the fragmentation pathway is provided by the observed rate of formation of acetic acid (CH_3COOH), which is consistent with our assumed fragmentation rate. However, the model overestimates peroxyacetyl nitrate (PAN) formation downwind by about 50%, suggesting the need for further refinements to the chemistry. This method could provide a way for classifying different smoke plume observations in terms of the average chemistry of their SVOCs, and could be used to study how the chemistry of these compounds (and the O_3 and OA they form) varies between plumes.

1 Introduction

Biomass burning is a major source of atmospheric trace gases and particles that impact air quality and climate (e.g., Crutzen and Andreae, 1990; van der Werf, 2010; Akagi et al., 2011). Within minutes after emission, rapid and complex photochemistry within the young biomass burning smoke plumes can lead to significant increases in the concentrations of secondary pollutants such as ozone (O_3; e.g., Mauzerall et al., 1998; Goode et al., 2000; Hobbs et al., 2003; Pfister et al., 2006; Lapina et al., 2006; Val Martin et al., 2006; Yokelson et al., 2009; Jaffe and Widger, 2012; Akagi et al., 2012, 2013), peroxyacetyl nitrate (PAN; e.g., Jacob et al., 1992; Alvarado et al., 2010, 2011; Fischer et al., 2014), and organic aerosol (OA; e.g., Hobbs et al., 2003; Grieshop et al., 2009a, b; Yokelson et al., 2009; Hennigan et al., 2011; Heringa et al., 2011; Vakkari et al., 2014), while other smoke plumes can show little to no formation of O_3 (e.g., Alvarado et al., 2010, Zhang et al., 2014) or OA (e.g., Akagi et al., 2012). Understanding the atmospheric chemistry of these young smoke plumes, conditions of which can lead to the secondary formation of O_3, PAN, and OA, is thus critical to understanding the impact of these plumes on atmospheric composition and the resulting impacts on air quality, human health, and climate. However, global- and regional-scale Eulerian models of atmospheric chemistry artificially dilute biomass burning emissions into large-scale grid boxes, which can result in large errors in the predicted concentrations of O_3 and aerosol species downwind (e.g., Alvarado et al., 2009; Zhang et al., 2014). In contrast, plume-scale Lagrangian models allow us to examine the chemical and physical transformations within these concentrated plumes in detail and can be used to develop parameterizations for this aging process for coarser models (e.g., the parameterizations of Vinken et al., 2011, and Holmes et al., 2014, for ship plumes).

Our understanding of the formation of ozone within biomass burning plumes is still poor, due both to the limited observational data available on O_3 formation in smoke plumes and the highly variable results seen in the available observations. Several aircraft and surface studies of the chemistry of young biomass burning smoke plumes have found significant formation of O_3 within smoke plumes. For example, Baylon et al. (2014) reported $\Delta O_3 / \Delta CO$ from 0.4 to 11 %, corresponding to O_3 enhancements of 3.8 to 32 ppbv in 19 wildfire plumes samples at Mt. Bachelor Observatory. They note that plumes that have low values of $\Delta O_3 / \Delta CO$ can still correspond to significant O_3 enhancements in concentrated plumes, with one event with a $\Delta O_3 / \Delta CO$ value of 0.81 % corresponding to an O_3 enhancement of 17 ppbv. Akagi et al. (2013) found significant O_3 formation ($\Delta O_3 / \Delta CO$ from 10 to 90 %) within 2 h for all of the South Carolina prescribed fires studied; Parrington et al. (2013) found values of $\Delta O_3 / \Delta CO$ increased from 2.0 ± 0.8 % in boreal biomass burning plumes less than 2 days old over eastern Canada to 55 ± 29 % in plumes that were more than 5 days old. Sim-

ilarly, Andreae et al. (1994) found that aged plumes (over 10 days old) from the biomass burning regions of South America and Africa had $\Delta O_3 / \Delta CO$ values between 20 and 70 %. However, other studies, mainly in boreal regions, have found little formation or even depletion of O_3 in some young biomass burning plumes (e.g., Alvarado et al., 2010). This low-O_3 formation is likely due to a combination of low emissions of NO_x from the boreal fires (Akagi et al., 2011), sequestration of NO_x in PAN and other organic nitrates (e.g., Jacob et al., 1992; Alvarado et al., 2010, 2011), and reduced rates of photochemical reactions due to aerosol absorption and scattering (e.g., Jiang et al., 2012). Similarly, some studies have shown that fires can contribute to high surface O_3 events that exceed the USA air quality standard for O_3 (e.g., Jaffe et al., 2013), while other studies suggest that this enhanced surface O_3 is only present when the biomass burning emissions mix with anthropogenic pollution (Singh et al., 2012; Zhang et al., 2014). However, even given the observed variability among fires, it is likely that biomass burning has an impact on the concentrations of tropospheric O_3. For example, the recent review of Jaffe and Widger (2012) estimated that biomass burning could contribute 170 Tg of O_3 per year, accounting for 3.5 % of all global tropospheric O_3 production. However, Sudo and Akimoto (2007) estimated that over a third of tropospheric O_3 came from free tropospheric chemical production due to biomass burning outflow from South America and South Africa.

The NO_x emitted by biomass burning is rapidly converted into a wide variety of inorganic nitrate (i.e., $HNO_{3(g)}$ and total aerosol inorganic nitrate, or $NO_{3(p)}$) and organic nitrate species (i.e., alkyl nitrates ($RONO_2$) and peroxy nitrates (RO_2NO_2), including PAN; Jacob et al., 1992; Yokelson et al., 2009; Alvarado et al., 2010, 2011; Akagi et al., 2012). The rate at which this conversion occurs and the relative production of inorganic nitrate, alkyl nitrates, and peroxy nitrates are a key control of the impact of the biomass burning on O_3 production and atmospheric composition. Once NO_x is converted to inorganic or organic nitrate, it is generally unavailable for further O_3 formation near the fire source. Furthermore, while conversion of NO_x into inorganic nitrate ($HNO_{3(g)} + NO_{3(p)}$) is generally irreversible (except for the slow reaction of $HNO_{3(g)}$ with OH), peroxy nitrate species like PAN can act as thermally unstable reservoirs of NO_x, allowing transport of NO_x in the upper atmosphere far from the original source and then producing NO_x via thermal decomposition as the air mass descends to the surface (e.g., Fischer et al., 2010). This regenerated NO_x can thus impact O_3 formation far from the original source.

In addition, photochemistry within the smoke plume can rapidly oxidize non-methane organic compounds (NMOCs), both those that were emitted in the gas phase and those emitted in the particle phase, lowering their vapor pressure and thus leading to the formation of secondary organic aerosol (SOA). As with O_3 and PAN formation, the formation of SOA in smoke plumes is highly variable, with the ratio of

OA to CO_2 increasing by a factor of 2–3 downwind of some biomass burning fires (e.g., Hobbs et al., 2003; Grieshop et al., 2009a, b; Yokelson et al., 2009; Hennigan et al., 2011; Heringa et al., 2011; Vakkari et al., 2014), while in others it can stay constant or even decrease (e.g., Capes et al., 2008; Akagi et al., 2012). For cases where little net SOA formation was observed, it is likely that the NMOCs were still being oxidized. However, in these cases the fragmentation of the organic species after oxidation (leading to higher volatility products) is likely more common than functionalization (i.e., the addition of oxygen to the organic species, leading to lower volatility products).

Plume-scale Lagrangian parcel models can be used to investigate the evolution of O_3, PAN, and OA in smoke plumes in detail, as their relatively simple parameterizations of plume dispersion and transport allow detailed simulation of the chemical and microphysical processes taking place within the young smoke plumes (e.g., Mauzerall et al., 1998; Mason et al., 2001, 2006; Jost et al., 2003; Trentmann et al., 2005; Alvarado and Prinn, 2009; Heilman et al., 2014; Arnold et al., 2015). Previous plume-scale modeling studies have greatly advanced our understanding of these transformations. Mauzerall et al. (1998) found that O_3 production within biomass burning plumes was limited by the concentration of NO_x and that the formation and subsequent degradation of PAN helped to maintain NO_x concentrations. Mason et al. (2001) and Trentmann et al. (2003) showed that oxygenated volatile organic compounds (OVOCs) were critical to the formation of O_3 within the smoke plumes. More recent work has suggested heterogeneous chemistry and currently unidentified organic species as potential explanations for the rapid formation of O_3 and organic aerosol seen within some smoke plumes (Trentmann et al., 2005; Mason et al., 2006; Alvarado and Prinn, 2009).

The Aerosol Simulation Program (ASP) was developed to simulate the formation of ozone and SOA within young biomass burning plumes (Alvarado, 2008). ASP v1.0 was used to simulate several African and North American plumes (Alvarado and Prinn, 2009) and to simultaneously simulate the chemistry, dynamics, and radiative transfer within a smoke plume using a high-resolution three-dimensional plume model (Alvarado et al., 2009). Alvarado and Prinn (2009) showed while their initial ASP v1.0 simulations underestimated the formation of both OH and O_3 in the Timbavati savannah fire smoke plume (Hobbs et al., 2003), if the OH concentration in ASP v1.0 was fixed at the estimated value of 1.7×10^7 molecules cm^{-3} then the model reproduced the observed concentrations of O_3. This suggested that the model was missing an important source of OH, and they proposed a heterogeneous reaction of NO_2 on aerosol particles producing HONO, followed by the photolysis of HONO into NO and OH, as a candidate for the missing source of OH within the smoke plume. Alvarado and Prinn (2009) also found that including only SOA formation from known SOA precursors (mainly aromatic species like toluene) underesti-

mated the concentrations of organic aerosol observed downwind by $\sim 60\%$, suggesting that the model was missing a large source of SOA. They proposed that the large amount of gas-phase organic compounds that were unidentified by the then current measurement techniques (Christian et al., 2003; Warneke et al., 2011) could include the precursors for the missing SOA. Assuming these compounds had SOA yields similar to monoterpenes gave the observed SOA formation.

In this paper, we describe recent updates to the gas-phase chemistry and SOA formation modules in ASP. We use this updated version (ASP v2.1) to simulate the chemical evolution of a young biomass burning smoke plume sampled over California in November near San Luis Obispo (the Williams fire; Akagi et al., 2012). The analysis of the O_3, PAN, and OA evolution in biomass burning plumes is complicated by the fact that a large fraction (30–50 % by carbon mass; Christian et al., 2003; Warneke et al., 2011) of the NMOCs present in smoke plumes are unidentified, and thus their oxidation chemistry is not well known.

Furthermore, while there was clear secondary formation of O_3 and PAN within the Williams fire plume, the dilution-corrected amount of OA in the plume decreased slightly (Akagi et al., 2012). Most current methods for modeling the OA evolution in smoke plumes lead to significant secondary growth of the OA (e.g., Grieshop et al., 2009a), but we need instead to modify ASP v2.1 to simulate both this slight loss of OA and the chemical formation of O_3, PAN, and other gas-phase species.

Here we present a method for simultaneously accounting for the impact of the unidentified organic compounds (here collectively called "SVOCs") on the formation of OA and O_3, drawing on the volatility basis set (VBS) approach (e.g., Robinson et al., 2007) for modeling OA and the concept of the mechanistic reactivity of a mixture of organic compounds (e.g., Carter, 1994). We show that this method can successfully simulate the Williams fire plume observations using reasonable assumptions about the chemistry of the unidentified SVOCs.

Section 2 describes the updates to the gas-phase chemistry and SOA formation modules of ASP for version 2.1. Section 3 discusses our validation of the gas-phase chemistry in ASP v2.1 against data from a smog chamber (Carter et al., 2005). Section 4 describes the Williams fire and summarizes the available observations of the smoke plume from Akagi et al. (2012). Section 5 discusses the results of the ASP simulation of the Williams fire, including sensitivity tests to investigate the chemistry of the unidentified SVOCs and their impacts on O_3, PAN, other trace gases, and OA, while Sect. 6 gives the conclusions of our study and directions for future work.

2 Updates to the Aerosol Simulation Program

An overview of ASP v1.0 is given by Alvarado and Prinn (2009), and the routines are described in detail in Alvarado (2008). Here we briefly discuss the modules of ASP that have not changed since Alvarado and Prinn (2009) in Sect. 2.1 before describing the updates to the gas-phase chemistry (Sect. 2.2) and SOA formation (Sect. 2.3) routines for ASP v2.1.

2.1 ASP modules

Aerosols are represented in ASP by a single moving-center sectional size distribution, where the aerosol concentrations are distributed over increments in radius space (Jacobson, 1997, 2002, 2005). ASP includes modules to calculate aerosol thermodynamics, gas-to-aerosol mass transfer (condensation/evaporation), and coagulation of aerosols. The thermodynamics module in ASP uses the mass flux iteration (MFI) method of Jacobson (2005) to calculate the equilibrium concentration of gas and aerosol species. Equilibrium constants for the inorganic electrolyte reactions match those of Fountoukis and Nenes (2007). Binary activity coefficients of inorganic electrolytes are calculated using the Kusik–Meissner method (Kusik and Meissner, 1978; Resch, 1995), as are the mean activity coefficients. The water content of inorganic aerosols is calculated with an iterative routine that calculates water activities for aqueous solutions of a single electrolyte using a formula based on the Gibbs–Duhem equation (Steele, 2004). Steele (2004) and Alvarado (2008) found this approach compares well with other inorganic aerosol thermodynamics models such as ISORROPIA (Nenes et al., 1998; Fountoukis and Nenes, 2007).

Mass transfer between the gas and aerosol phases is calculated in ASP using a hybrid scheme where the flux-limited kinetic equations governing the condensation/evaporation of H_2SO_4 and organic species are integrated, whereas NH_3, HNO_3, and HCl are assumed to be in equilibrium (Alvarado, 2008). Aerosol coagulation is calculated using the semi-implicit scheme of Jacobson (2005) with a Brownian coagulation kernel.

2.2 Gas-phase chemistry updates

The gas-phase chemistry within the ASP model for Version 2.1 has been completely revised from ASP v1.0, which used the CalTech Atmospheric Chemistry Mechanism (CACM; Griffin et al., 2005). The revised ASP v2.1 gas phase chemical mechanism includes 1608 reactions between 621 species. Examples of the gas-phase species input file and the reaction mechanism input file for ASP v2.1, along with other key chemical input files, are included in the Supplement.

All inorganic gas-phase chemistry within ASP v2.1 was updated to follow the International Union of Pure and Applied Chemistry (IUPAC) recommendations (Atkinson et al., 2004; updated data downloaded from http://iupac.pole-ether.fr/, accessed June 2012). We also tested the Jet Propulsion Laboratory (JPL) recommendations (evaluation no. 17; Sander et al., 2011) for these rate constants, but found that the differences between the recommendations generally made little difference to the model simulations, and as the IUPAC values were closer to those in ASP v1.0, these values were used.

All gas-phase chemistry for organic compounds containing four carbons or less has been "unlumped", i.e., the chemistry for each individual organic compound is explicitly resolved. This was done by following the reactions of the Leeds master chemical mechanism (MCM) v3.2 (http://mcm.leeds.ac.uk/MCM/, accessed June 2012; Jenkin et al., 1997, 2003; Saunders et al., 2003; Bloss et al., 2005) for these species.

The chemical mechanism of isoprene within ASP v2.1 has been updated to follow the Paulot et al. (2009a, b) isoprene scheme, as implemented in GEOS-Chem and including corrections based on more recent studies (e.g., Crounse et al., 2011, 2012). The (lumped) chemistry for all other organic compounds in ASP has been updated to follow the regional atmospheric chemistry mechanism (RACM) v2 (Goliff et al., 2013). We chose RACM2 over the [California] Statewide Air Pollution Research Center (SAPRC-07) (Carter, 2010) and Carbon Bond (CB05) (Yarwood et al., 2005) lumped chemical mechanisms as the treatment of peroxy radicals in the RACM2 mechanism was more similar to the treatment in the Leeds MCM and the Paulot isoprene scheme, resulting in a more consistent chemical mechanism for ASP v2.1.

Photolysis rates are calculated offline using the tropospheric ultraviolet and visible (TUV) radiation model version 5.0 (Madronich and Flocke, 1998) for 15 min increments, which are then linearly interpolated in ASP. Alvarado and Prinn (2009) assumed a "clear-sky" radiation field that ignored the effect of aerosol absorption and scattering on the calculated photolysis rates. Here we instead estimate the time-dependent aerosol, O_3, SO_2, and NO_2 concentrations within the smoke plumes and calculate their effect on the photolysis rates at different heights within the plume. In the TUV simulations, we assume no clouds and that the initial smoke plume AOD at 330 nm decreases due to dilution assuming a background concentration of ~ 0, and the aerosol is assumed to have a constant (both with time and wavelength) single scattering albedo (SSA) of 0.9 based on the review of Aerosol Robotic Network (AERONET) biomass burning smoke optical property retrievals by Reid et al. (2005a). We also dilute the initial plume concentrations of the trace gases NO_2 and SO_2 assuming a background concentration of ~ 0, as these species can also absorb ultraviolet and visible (UV–VIS) light and thus can impact photolysis rates. For the photolysis rate calculations only, O_3 is assumed to be 0 initially and increased after 15 min to a constant value based on the observed formation of O_3 within the smoke plume.

Section 5.1 has more details on the specific approach and quantitative values used for the Williams fire.

2.3 SOA formation updates

We have updated the SOA formation module to follow the semi-empirical VBS model of Robinson et al. (2007). Our implementation of this scheme followed the approach used by Ahmadov et al. (2012) to link the VBS scheme with the RACM chemical mechanism within WRF-Chem. We use nine surrogates or "bins" for SVOCs as in Dzepina et al. (2009), rather than only four as in Ahmadov et al. (2012). The saturation mass concentration at 300 K (C^*; see Robinson et al., 2007) of each SVOC bin differs by a factor of 10, and covers the range from 0.01 to $1.0 \times 10^6\ \mu g\,m^{-3}$. Note that "SVOC" as defined in this paper includes both semi-volatile organic compounds (C^* between 10^{-2} and $10^3\ \mu g\,m^{-3}$) and intermediate volatility organic compounds (species with C^* between 10^4 and $10^6\ \mu g\,m^{-3}$) as defined in Dzepina et al. (2009), but we refer to both of these species classes collectively as "SVOCs" rather than as "S/IVOCs" as in Dzepina et al. (2009) for simplicity. Following the Model to Predict the Multiphase Partitioning of Organics (MPMPO) of Griffin et al. (2003, 2005) and Pun et al. (2002), we assumed that an aqueous phase and a mixed hydrophobic organic phase are always present in the aerosol. Partitioning of organics between the gas and hydrophobic phase is governed by Raoult's law (assuming that all hydrophobic-phase Organic matter (OM) is quasi-liquid and can dissolve organics as in Pankow, 1994a, b), while partitioning of organics into the aqueous phase is governed by Henry's law. Following Pun et al. (2002), we assumed that (1) there is no interaction between the aqueous-phase inorganic ions and the aqueous-phase organics, and thus no organic salt formation, and (2) the activity coefficients of the organic ions (formed by the dissociation of organic acids) are equivalent to those of the corresponding molecular solute. We further assumed that the pH of the aqueous phase is dominated by the strong inorganic acids and bases, and that the pH effects of the dissociating organic acids are negligible.

Like most organic compounds, SVOCs will react with OH. Most mechanisms for this chemistry (e.g., Robinson et al., 2007; Dzepina et al., 2009; Grieshop et al., 2009a, b; Ahmadov et al., 2012) parameterize this chemistry by assuming that the SVOCs react with OH to form a lower volatility SVOC, as in the reaction:
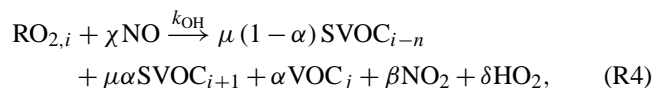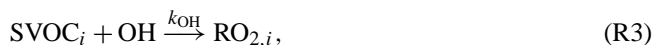
$$SVOC_i + OH \xrightarrow{k_{OH}} \mu SVOC_{i-n}, \tag{R1}$$
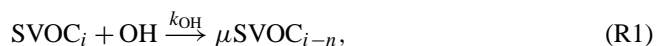
where μ is the relative mass gain due to oxidation (e.g., via O addition), k_{OH} is the reaction rate with OH, and n is the "volatility shift", or by how many factors of 10 to lower the C^* of the product with each OH reaction. This simplified chemistry can be extended to account for the fact that the SVOCs could fragment during oxidation, leading to higher volatility products:

$$SVOC_i + OH \xrightarrow{k_{OH}} \mu(1-\alpha)SVOC_{i-n}$$
$$+ \mu\alpha SVOC_{i+1} + \alpha VOC_j, \tag{R2}$$
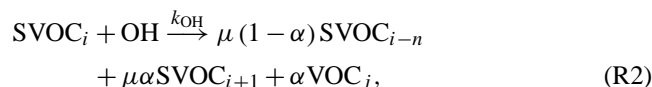
where α is the fraction of $SVOC_i$ that fragments into $SVOC_{i+1}$ and VOC_j. Shrivastava et al. (2013) used a similar approach to show that adding SVOC fragmentation to WRF-Chem simulations of the Mexico City plateau improved the model's ability to simulate the observed concentrations of SOA. However, the highly simplified chemistry of Reactions (R1) or (R2) is not appropriate for situations where reactions with the SVOC compounds are a potentially significant sink of OH, such as in a concentrated smoke plume. Thus, in ASP v2.1, the average, lumped chemistry of the SVOCs is instead parameterized in a more realistic manner for a generic organic species, following the idea of "mechanistic reactivity" (e.g., Carter, 1994; Bowman and Seinfeld, 1994a, b; Seinfeld and Pandis, 1998). After reaction with OH SVOCs produce peroxy radicals (RO_2), which can react with NO to form NO_2 and HO_2, thereby regenerating OH and forming O_3. Reactions (R3) and (R4) show this more general chemical mechanism for the SVOCs:

$$SVOC_i + OH \xrightarrow{k_{OH}} RO_{2,i}, \tag{R3}$$

$$RO_{2,i} + \chi NO \xrightarrow{k_{OH}} \mu(1-\alpha)SVOC_{i-n}$$
$$+ \mu\alpha SVOC_{i+1} + \alpha VOC_j + \beta NO_2 + \delta HO_2, \tag{R4}$$

where $k_{RO_{2,i}}$ is assumed to be $4.0 \times 10^{-12}\ cm^3\ molecule^{-1}\ s^{-1}$ based on the reaction rate for the peroxy radicals from long-chain alkanes and alkenes with NO in RACM2 (Goliff et al., 2013). We can see that $\chi - \beta$ is the number of NO_x lost (implicitly via the addition of a nitrate group to the product SVOCs), 1-δ is the number of HO_x lost, and $\beta + \delta$ is the number of O_3 made per reaction (by subsequent reactions of NO_2 and HO_2 to generate O_3). For example, the values for long-chain alkanes (HC8) in the RACM$_2$ mechanism (Goliff et al., 2013) would be $\chi = 1$, $\delta = 0.63$, and $\beta = 0.74$, such that $0.26\ NO_x$ and $0.37\ HO_x$ are lost and $1.37\ O_3$ are formed per reaction. Note that the mechanism of Reactions (R3) and (R4) is still highly simplified: we assume that reaction of SVOC with OH always produces a RO_2 radical, and that the RO_2 produced does not react with HO_2 or another RO_2. Also note that Reactions (R3) and (R4) represent the average chemistry of the unknown species collectively, and may not apply to any individual species in that mixture. Our purpose is less to detail all the possible reactions of the unidentified SVOCs and more to explore how their average chemistry might affect O_3 and OA evolution in smoke plumes. The specific combinations of parameters for Reactions (R3) and (R4) that were evaluated in this study are shown in Table 2.

We also adjusted the calculation of aerosol water content to use the "kappa" (κ) parameterization of organic hygroscopicity (Petters and Kreidenweis, 2007) for the lumped SVOCs. In this parameterization, the hygroscopicity parameter κ for the organic species is defined as

$$\frac{1}{a_{\mathrm{w}}} = 1 + \kappa \frac{V_{\mathrm{s},j}}{V_{\mathrm{w},i}}, \tag{1}$$

where a_{w} is the activity of water in the solution (equal to the relative humidity (RH) at equilibrium), $V_{\mathrm{s},i}$ is the volume of the dry organic solute i and $V_{\mathrm{w},i}$ is the volume of water in the solution. The water content calculated for each organic species, along with that calculated for the inorganic solution ($V_{\mathrm{w,inorg}}$; see Sect. 2.1 above) are then combined using the Zdanovskii, Stokes, and Robinson (ZSR) approximation (Zdanovskii, 1948; Stokes and Robinson, 1966):

$$V_{\mathrm{w}} = \frac{a_{\mathrm{w}}}{1 - a_{\mathrm{w}}} \sum_i \kappa_i V_{\mathrm{s},j} + V_{\mathrm{w,inorg}}. \tag{2}$$

3 ASP photochemistry evaluated with smog chamber data

To evaluate the performance of the updated photochemical mechanism in ASP v2.1 in predicting the formation of ozone, several test simulations were performed to compare the results of the mechanism to laboratory smog chamber data. This comparison provides us with a baseline for interpreting the results of our simulation of O_3 formation in the Williams fire in Sect. 5. The data used for the comparison came from the Environmental Protection Agency (EPA) chamber of Carter et al. (2005), which consists of two collapsible 90 m^3 FEP Teflon reactors (chambers A and B).

Table S1 in the Supplement shows the temperature and initial reactant concentrations used in our model to simulate each chamber study. All model simulations were performed at a pressure of 1000 mbar, a RH of 1 %, and a CH_4 concentration of 1800 ppbv. The temperature and concentration data were provided by William P. L. Carter (http://www.cert. ucr.edu/~carter/SAPRC/SAPRCfiles.htm, accessed March 2014). The EPA chamber runs used an eight compound surrogate for ambient VOC concentrations, consisting of formaldehyde (HCHO), ethylene (C_2H_4), propene, trans-2-butene, n-butane, n-octane, toluene, and m-xylene (Carter et al., 1995, 2005). The initial concentrations of HONO were extrapolated from CO–NO_x and n-butane–NO_x runs to account for the potential chamber radical source (Carter et al., 2005).

Table S2 in the Supplement presents the rates of offgassing (i.e., re-emission of HCHO and other species from the walls of the reaction chamber), wall reaction rates, and selected photolysis rates for the chamber experiments considered here. The off-gassing rate for HONO was determined as the rate that enabled the SAPRC-99 chemical mechanism
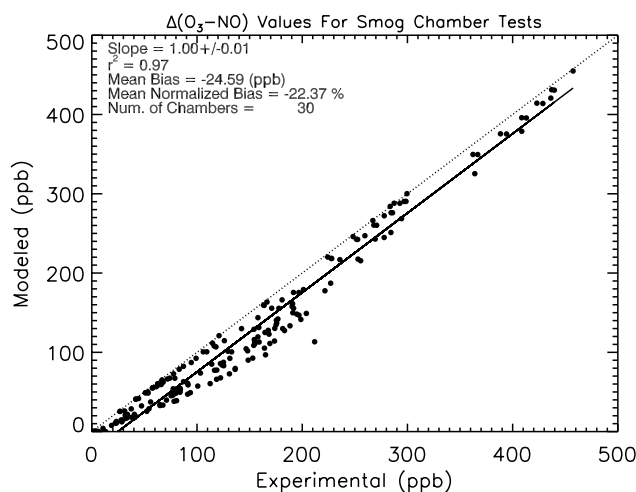

Figure 1. ASP calculated hourly values of $\Delta([O_3] - [NO]) \equiv ([O_3]_{\mathrm{final}} - [NO]_{\mathrm{final}}) - ([O_3]_{\mathrm{initial}} - [NO]_{\mathrm{initial}})$ vs. the values measured in the EPA chamber of Carter et al. (2005) for 30 "full surrogate" experiments. Note that all time points for the 30 chamber experiments are plotted, not just the final values.

(Carter, 2000) to best predict the O_3 formation observed in CO-, HCHO-, and CO–HCHO-air experiments performed within the chamber (Carter et al., 2005). The rate in Chamber A was found to be slightly higher than that in Chamber B, so different values are used for the chambers. The off-gassing rate of HCHO was chosen to match the low but measurable amount of HCHO found even in pure air and CO–NO_x experiments in the chamber. Heterogeneous wall loss reaction rates for O_3, NO_2, and N_2O_5 were also estimated from reactor observations (Carter et al., 2005). The photolysis rate of NO_2 in the chambers was measured directly, and scaling factors for the other photolysis rates were calculated by Carter et al. (2005) from the relative spectral intensity of the arc lamp.

Following Carter et al. (2005), we evaluated the ability of our mechanism to simulate the total amount of NO oxidized and O_3 formed in the experiments, measured as $\Delta([O_3] - [NO])_t \equiv ([O_3]_t - [NO]_t) - ([O_3]_{\mathrm{initial}} - [NO]_{\mathrm{initial}})$. The hourly results of the comparisons for $\Delta([O_3] - [NO])$ are presented in Fig. 1. We can see that the ASP v2.1 mechanism tends to underestimate $\Delta([O_3] - [NO])$, with a mean absolute bias of -24.6 ppbv and a mean normalized bias of -22.4 %. Comparisons of the ASP calculations for O_3, NO, and NO_x (not shown) show that this model underestimate of $\Delta([O_3] - [NO])$ is primarily due to the model underestimating O_3 formation, rather than underestimating the loss of NO or NO_x. Similarly, the ASP v2.1 calculations for the concentrations of the organic gas species matches well with the chamber measurements (not shown) except for HCHO, where the secondary formation of HCHO appears to be underestimated. Figure 2 shows the bias in $\Delta([O_3] - [NO])$ vs. the initial ratio of the mixing ratio of reactive organic gases (ROGs; e.g., the concentration of the
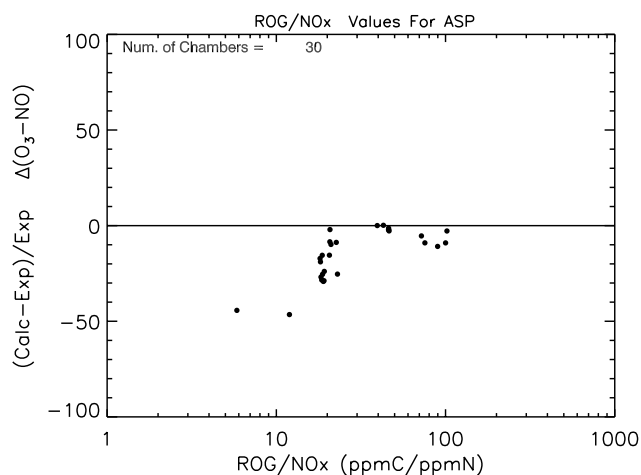
Figure 2. Percentage bias in final $\Delta([O_3] - [NO])$ vs. the initial ratio of reactive organic gases (ROG) to NO_x (ppm C / ppm N) for the chamber experiments.

surrogate gases in ppm C) to the mixing ratio of NO_x (in ppm N). We can see that the bias is between 0 and -10% for ROG / NO_x ratios greater than 30, but increases to -40 to -50% for "high NO_x" cases (ROG / NO_x ratios $<< 20$). For comparison, the initial ROG / NO_x ratio in biomass burning smoke can range from ~ 10 to 100 ppm C / ppm N (Akagi et al., 2011; assuming that the total NMOC mass is 1.6 times the mass of carbon in these compounds). Both the general underestimation of $\Delta([O_3] - [NO])$ and the increase of the negative bias at low ROG / NO_x concentrations is consistent with the behaviors of the SAPRC-99 (Carter et al., 2005), SAPRC-07 (Carter, 2010), and CB05 (Yarwood et al., 2005) mechanisms evaluated against the EPA chamber data. Carter (2010) noted that this underprediction of O_3 at low ROG / NO_x ratios was apparently linked to the presence of aromatics in the surrogate mixture, with comparisons of SAPRC-07 with EPA chamber runs with a non-aromatic surrogate mixture showing a positive bias of about $+25\%$ for cases with low ROG / NO_x ratios.

4 Williams fire data

The Williams fire ($34°41'45''$ N, $120°12'23''$ W) was sampled by the US Forest Service (USFS) Twin Otter aircraft from 10:50 to 15:20 LT on 17 November 2009 (Akagi et al., 2012). The fire burned approximately 81 hectares of scrub oak woodland understory and coastal sage scrub. Skies were clear all day and RH was low (11–26 %) with variable winds (2–5 m s^{-1}). The Williams fire smoke plume showed significant secondary production of O_3 and PAN, but the enhancement ratio of OA to CO_2 decreased slightly downwind (Akagi et al., 2012). In this study, we use the processed data from Akagi et al. (2012) that provided concentrations of several trace gases and OA measured during several quasi-

Lagrangian transects of the Williams fire. Full details on the measurements made and the processing of the data for the plume transects are given in Akagi et al. (2012); those used in this study are briefly described here.

4.1 Airborne Fourier transform infrared spectrometer

The University of Montana airborne Fourier transform infrared spectrometer (AFTIR) system and the instruments described below were deployed on a US Forest Service (USFS) Twin Otter aircraft. The AFTIR (Yokelson et al., 1999, 2003) was used to measure 21 gas-phase species, including carbon dioxide (CO_2), carbon monoxide (CO), nitric oxide (NO), nitrogen dioxide (NO_2), nitrous acid (HONO), PAN, ozone (O_3), glycolaldehyde ($HCOCH_2OH$), C_2H_4, HCHO, acetic acid (CH_3COOH), and formic acid (HCOOH). IR spectra were collected at 1 Hz. "Grab samples" of air were selected by closing the valves for 1–2 min to allow signal averaging, and the resulting IR spectra were analyzed to quantify all detectable compounds.

4.2 Aerosol mass spectrometer

An Aerodyne compact time-of-flight (CToF) aerosol mass spectrometer (herein referred to as AMS) measured aerosol chemical composition in a repeating cycle for four out of every 12 s during flight, including within the smoke plume. The AMS has been described in great detail elsewhere (Drewnick et al., 2005; Canagaratna et al., 2007). An isokinetic particle inlet sampling fine particles with a diameter cut-off of a few microns (Yokelson et al., 2007; Wilson et al., 2004) supplied the AMS. As the AMS does not measure super-micron particles, the inlet transmission should not have affected the results. A collection efficiency of 0.5 (Huffman et al., 2005; Drewnick et al., 2003; Allan et al., 2004) was applied to the AMS data, which were processed to retrieve the mass concentration for the major non-refractory particle species: OA, non-sea salt chloride, nitrate, sulfate, and ammonium, with $< 36\%$ uncertainty.

4.3 Other measurements

The ambient three-dimensional wind velocity, temperature, RH, and barometric pressure were measured at a frequency of 1 Hz with a wing-mounted AIMMS-20 probe (Aventech Research, Inc.; Beswick et al., 2008). A non-dispersive infrared instrument (NDIR) (LiCor model 7000) measured CO_2 (at 0.5 Hz) from the third channel on the isokinetic particle inlet that also supplied the AMS.

5 ASP simulation of Williams fire

5.1 ASP setup

As in Alvarado and Prinn (2009), we simulated the Williams fire smoke plume using ASP within a simple Lagrangian parcel model following Mason et al. (2001). We assume a Lagrangian parcel of fixed vertical extent (H; here assumed to be 1 km) and down-trajectory length (L), but variable cross-trajectory width $y(t)$. The temperature and pressure of the parcel are assumed to be constant. The full continuity equations for the Lagrangian parcel model are then

$$\frac{dC_q}{dt} = -\frac{4K_y}{\left(y_o^2 + 8K_y t\right)}\left(C_q - C_q^a\right) - \frac{v_d}{H}C_q$$
$$+ \left(\frac{dC_q}{dt}\right)_{cond} + \left(\frac{dC_q}{dt}\right)_{chem}, \quad (3)$$

$$\frac{dn_i}{dt} = -\frac{4K_y}{\left(y_o^2 + 8K_y t\right)}\left(n_i - n_i^a\right) - \frac{v_d}{H}n_i + \left(\frac{dn_i}{dt}\right)_{cond},$$
$$+ \left(\frac{dn_i}{dt}\right)_{chem} \quad (4)$$

$$\frac{dC_{q,i}}{dt} = -\frac{4K_y}{\left(y_o^2 + 8K_y t\right)}\left(C_{q,i} - C_{q,i}^a\right) - \frac{v_d}{H}C_{q,i}$$
$$+ \left(\frac{dC_{q,i}}{dt}\right)_{cond} + \left(\frac{dC_{q,i}}{dt}\right)_{coag}\left(\frac{dC_{q,i}}{dt}\right)_{chem}, \quad (5)$$

where C_q is the concentration of gas-phase species (molecules cm^{-3} air), n_i is the number concentration of particles in size bin i (particles cm^{-3} air), $c_{q,i}$ is the concentration of aerosol species q in size bin i (mol cm^{-3} air), y_o is the initial plume width (m), and K_y represents the horizontal diffusivity of the atmosphere (m^2 s^{-1}). The superscript a indicates the concentration of the given species in the atmosphere outside of the parcel (i.e., the background concentration).

The first term on the right-hand side of Eqs. (3)–(5) represents the effect of plume dispersion on the concentrations. Note that y_o and K_y can be reduced to a single parameter, the initial dilution timescale $\tau_{mix,o}$:

$$-\frac{4K_y}{\left(y_o^2 + 8K_y t\right)} = -\frac{1}{\frac{y_o^2}{4K_y} + 2t} = -\frac{1}{\tau_{mix,o} + 2t}. \quad (6)$$

The second term on the right-hand side of Eqs. (3)–(5) is the effect of deposition on the concentrations, where v_d is the deposition velocity (m s^{-1}). We set the dry deposition velocity equal to 0 for gas-phase species, as the plume did not touch the ground during the modeled period, and use the size-dependent terminal velocity of the aerosol particles as the deposition velocity for aerosol species assuming a 1 km thick plume. As submicron aerosol dominated the aerosol
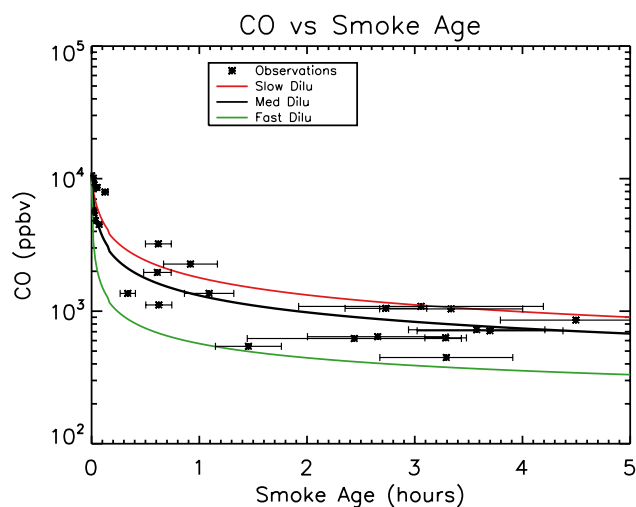

Figure 3. CO mixing ratio (ppbv) vs. smoke age. Red, black, and green are for the slow, best fit (medium), and fast plume dilution rates. Asterisks are the measured mixing ratios, with the horizontal error bars showing the uncertainty in the estimated age, which is much larger than the uncertainty in the CO mixing ratio.

mass in the smoke plume, this deposition of aerosol species has a negligible effect on the results, and given the low RH during the Williams fire, we also did not include wet deposition of particles or gases. The remaining terms represent the change in gas- and particle-phase concentrations due to net mass transfer between the gas and aerosol phases (cond), coagulation of particles (coag), and chemical production and loss (chem).

The observed changes in CO mixing ratio were used to determine the best-fit model initial dilution timescale ($\tau_{mix,o} = 106.9$ s) as well as upper and lower limits of the timescale ($\tau_{mix,o}(0) = 15.0$ and 212.2 s, respectively), as shown in Fig. 3. Note that this dilution-fitting procedure neglects the impact of the chemical production and loss of CO on the observed concentrations, but modeling results with an inert tracer suggest that this error is $< 4\%$, much smaller than the dilution uncertainty represented by our upper and lower limit estimates. The temperature of the plume was set at a constant value of 288.4 K, pressure of 880 hPa, and RH of 15.7 % based on the observations of Akagi et al. (2012). The parcel was assumed to be emitted at 11:00 Pacific standard time (PST) and the model was integrated for 5 h. The integration of the different terms of the continuity Eqs. (3)–(5) were operator split for computational efficiency. The chemistry and mixing time steps were 1 s for the first 10 min of model integration due to the rapid dilution and chemical changes during this period, and were 60 s thereafter. The aerosol thermodynamics, condensation, and coagulation time steps were 60 s throughout.

The initial and background concentrations for the gas-phase inorganic and NMOC species are in Table S3 of the

Table 1. Definition of SVOC species following Grieshop et al. (2009a).

Species	C^*	ΔH_{vap}	MW	POA volatility distributions[a]	
	(μg m^{-3} @ 300 K)	(kJ mol^{-1})	(g mol^{-1})	Grieshop et al. (2009a)	This study[b]
SVOC$_1$	10^{-2}	77	524	0	0
SVOC$_2$	10^{-1}	73	479	0	0
SVOC$_3$	10^{0}	69	434	0.1	0.038
SVOC$_4$	10^{1}	65	389	0.14	0.0532
SVOC$_5$	10^{2}	61	344	0.33	0.1254
SVOC$_6$	10^{3}	57	299	0.33	0.1254
SVOC$_7$	10^{4}	54	254	0.1	0.038
SVOC$_8$	10^{5}	50	208	0	0.62
SVOC$_9$	10^{6}	46	163	0	0

[a] Relative mass emissions in each volatility bin. [b] Where the relative amounts of SVOC$_i$ ($i = 1$ to 7) are kept as in Grieshop et al. (2009a), but additional organic mass is added to SVOC$_8$ to account for the unidentified NMOC mass reported by Akagi et al. (2011).

Table 2. SVOC chemistry parameters in the mechanisms studied here; see Reactions (R3) and (R4) for definitions of the parameters.

Mechanism	$k_{OH} \times 10^{11}$ (cm^3 molecule^{-1} s^{-1})	μ	n	α	χ	β	δ
Grieshop et al. (2009)	2.0	1.4	2	0	0	0	0
Robinson et al. (2007)	4.0	1.075	1	0	0	0	0
Ahmadov et al. (2012)	1.0	1.075	1	0	0	0	0
Half fragmentation	1.0	1.075	1	.5	0	0	0
Optimized SVOC chemistry	1.0	1.075	1	.5	1	.5	.6

Supplement, and Table S4 gives the initial and background concentrations used for the aerosol species. Initial and background concentrations of trace gases and aerosols in the smoke were taken from observations of the Williams fire (Akagi et al., 2012), where available. Emission ratios for other species were calculated using the literature reviews of Akagi et al. (2011) and Andreae and Merlet (2001). Other background concentrations were taken from runs of the GEOS-Chem model (Bey et al., 2001), run for our period as in Fischer et al. (2014).

The volatility distribution for the POA was taken from the wood smoke study of Grieshop et al. (2009a, b). Table 1 shows the POA total mass fractions used for wood smoke in Grieshop et al. (2009a) and the values used in this study for the Williams fire. At the measured temperature (288.4 K) and initial concentration of organic aerosol in the Williams fire smoke plume (849 μg m^{-3}), the Grieshop et al. (2009a) POA volatility distribution implies that 81 % of the total mass of SVOC species SVOC$_1$ to SVOC$_7$ is in the aerosol phase, leaving 200 μg m^{-3} of SVOC species in the gas phase. Note that the May et al. (2013) POA volatility distribution (not shown in Table 1) is more volatile than Grieshop et al. (2009a), with 65 % of the total mass of SVOC species SVOC$_1$ to SVOC$_7$ in the aerosol phase, leaving about 460 μg m^{-3} of SVOC species in the gas phase.

However, Grieshop et al. (2009a) and May et al. (2013) were only able to measure species with a saturation mass concentration (C^*) of 10^4 μg m^{-3} or less. Furthermore, Akagi et al. (2011) provide emission factors for unidentified NMOCs from savannah/grassland and chaparral fires, with unidentified NMOCs estimated to be equal in mass to the identified species. The savannah/grassland estimate is about twice as large as the chaparral estimate, as fewer species have been identified in chaparral fires. Here we use the savannah/grassland estimate to calculate an emission ratio of 0.195 g unidentified NMOC g^{-1} CO, but assign this value an uncertainty of $\sim$ 50 %, consistent with the lower chaparral estimate. This implies that there is about 2000 $\pm$ 1000 μg m^{-3} of unidentified NMOCs in the gasphase of the smoke. So to be consistent with the EFs of Akagi et al. (2011) and the volatility distributions of Grieshop et al. (2009a) and May et al. (2013), there still needs to be another 1500–1800 μg m^{-3} of unidentified NMOCs initially in the plume with $C^* > 10^4$ μg m^{-3} which the techniques used by Grieshop et al. (2009) and May et al. (2013) would not have been able to measure. These remaining unidentified NMOCs were included as SVOC$_8$ ($C^* = 10^5$ μg m^{-3}), as shown in Table 1 and Tables S3 and S4 in the Supplement. Below we also discuss sensitivity tests that were performed to see how the results change if the remaining uniden-

tified NMOCs are considered as $SVOC_9$ ($C^* = 10^6 \, \mu g \, m^{-3}$) instead, as well as for an increase or decrease of the estimated unidentified SVOC concentrations by 50 %.

For all organic species, we assumed a constant $\kappa = 0.04$, corresponding to an O / C ratio of 0.25 (Jimenez et al., 2009) that is typical of biomass burning organic aerosol (Donahue et al., 2011). Since the RH in the Williams fire plume was very low, this assumption had little impact on our results.

The initial smoke aerosol size distribution was assumed to be a lognormal with a geometric mean diameter D_g of 0.10 µm and a standard deviation (σ) of 1.9 (unitless) based on Reid and Hobbs (1998) for flaming combustion of Brazilian cerrado, which structurally is a similar mix of shrubs and grasses as in the Williams fire. The initial total number concentration of aerosol particles (2.34×10^6 particles cm^{-3}) was calculated such that the initial total organic aerosol mass matched the $\Delta OA / \Delta CO_2$ emission ratio from Akagi et al. (2012). The evolution of the aerosol size distribution with time was simulated by ASP v2.1 using a center-moving sectional size distribution with 10 bins, 8 bins for particles with volume-equivalent spherical diameters between 0.05 and 2.0 µm, one for particles with diameters smaller than 0.05 µm, and one for particles with diameters greater than 2 µm.

Photolysis rates were calculated offline using TUV v5.0 (Madronich and Flocke, 1998) as noted in Sect. 2.2 above. The smoke aerosols were assumed to dilute with time according to the three dilution rates derived above (see Fig. 3). In the TUV simulations, we assumed no clouds and an initial AOD of 8.0 at 330 nm (consistent with the ASP v2.1 calculated initial extinction coefficient and the assumed plume thickness of 1 km), which decreases due to dilution assuming a background concentration of ~ 0. As noted above, we assumed a constant single scattering albedo of 0.9 based on Reid et al. (2005a). We also assumed initial plume and background concentrations of the trace gases NO_2 (initial 295 ppbv, background 0 ppbv) and SO_2 (initial 50.9 ppbv, background 0 ppbv). For the photolysis rate calculations only, O_3 was assumed to be 0 initially and increased after 15 min to a constant value of 100 ppbv to account for the observed formation of O_3 within the smoke plume. The overhead ozone column was assumed to be 278 Dobson units (DU), based on the average of values from the Ozone Monitoring Instrument (OMI) for 16 November 2009 (276 DU) and 18 November 2009 (280 DU) (accessed through Phttp:// jwocky.gsfc.nasa.gov/teacher/ozone_overhead.html on June 2012, now at http://ozoneaq.gsfc.nasa.gov/tools/ozonemap/). The surface albedo was assumed to be 0.035 based on the GEOS-Chem data file for the $0.5° \times 0.667°$ North American grid for November 1985, which is in turn based on data from the total ozone mapping spectrometer (TOMS). Photolysis rates were calculated for three altitudes: just above the plume (i.e., at 2.1 km altitude), near the middle of the plume (1.6 km), and near the bottom of the plume (1.1 km). This, combined with the three dilution rates, gave nine estimates of
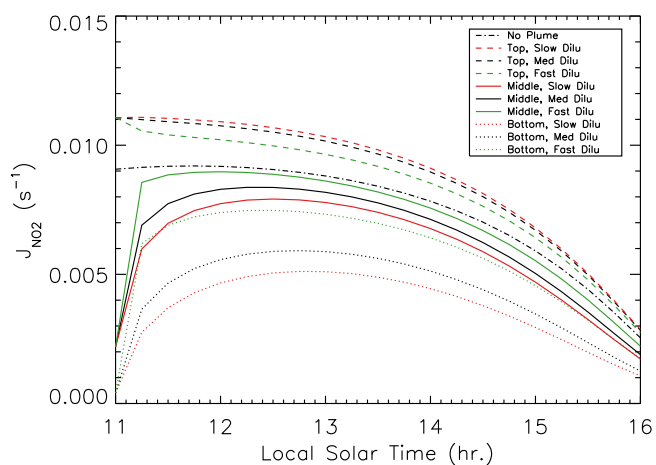

Figure 4. NO_2 photolysis rates (s^{-1}) vs. local time. Red, black, and green are for the slow, best fit (medium), and fast plume dilution rates. Dashed lines are for photolysis rates above the plume, solid lines are for the middle of the plume, and dotted lines are for the bottom of the plume, as described in the text. The black dot-dashed line is the clear-sky (no plume) photolysis rate.

photolysis rates vs. time. The nine values for the NO_2 photolysis rate (J_{NO_2}) are compared with the clear-sky (no aerosol) case in Fig. 4. In the middle of the plume (1.6 km), J_{NO_2} was reduced from an initial clear-sky value of $9 \times 10^{-3} \, s^{-1}$ to an initial value of $2 \times 10^{-3} \, s^{-1}$. However, by 15 min after emission J_{NO_2} in the middle of the plume increased to 6–$8.5 \times 10^{-3} \, s^{-1}$ depending on the dilution rate, showing that the plume reduced photolysis rates by 5–33 % after the initial, rapid dilution of the plume. J_{NO_2} was slightly enhanced above the plume (initially $1.1 \times 10^{-2} \, s^{-1}$) over the clear-sky value, and the photolysis rates were lowest in the bottom of the plume. As expected, the impact of the plume was larger for lower dilution rates, but the difference between the different dilution rates was largest for the bottom of the plume. Note that, while our assumption of a constant SSA is questionable as aerosol absorption is likely to change with both smoke age and with wavelength, our use of three dilution rates and three altitudes in the plume results in a wide range of photolysis rates used in this study, which can also account for uncertainties in the aerosol optical properties and other parameters used to calculate the photolysis rates.

5.2 ASP results with no unidentified SVOC chemistry

We first ran ASP assuming the unidentified SVOCs emitted by the fire are unreactive. Deficiencies in these simulations provide information on what the average chemistry of the unidentified SVOCs needs to be in order to explain the observations.

Figure 5 shows the ASP v2.1 results and Akagi et al. (2012) observations for the enhancement ratios (EnR; mol mol^{-1}) of O_3 and PAN in the Williams fire smoke plume
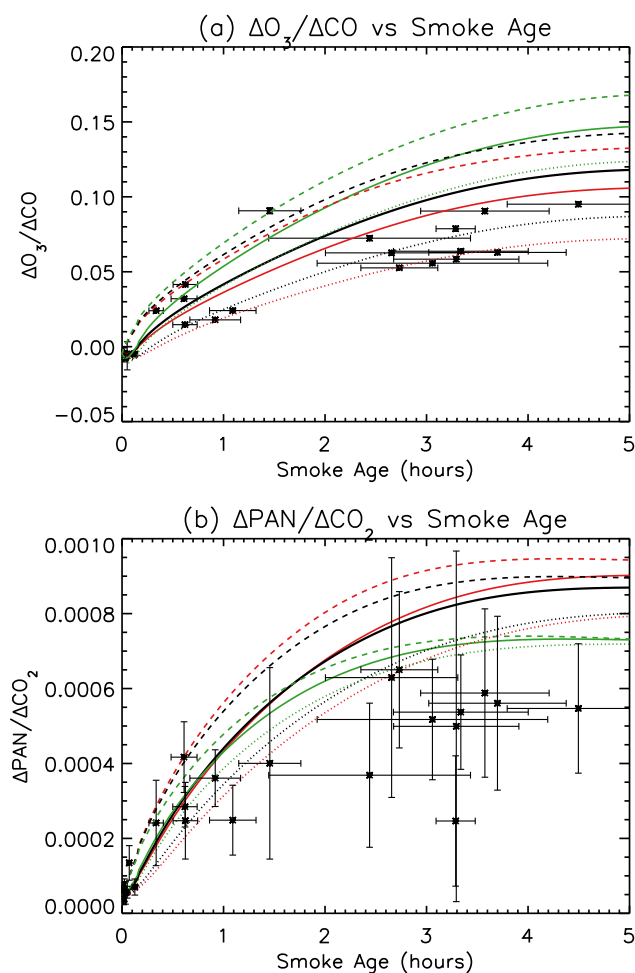
Figure 5. Enhancement ratios (mol mol^{-1}) of **(a)** O$_3$ to CO and **(b)** PAN to CO$_2$ vs. estimated smoke age when the chemistry of the unidentified SVOCs is not included in the model. Asterisks are the measured mixing ratios, with the horizontal error bars showing the uncertainty in the estimated age and the vertical error bars showing the uncertainty in the measurement. Red, black, and green are ASP results for the slow, best fit (medium), and fast plume dilution rates. Dashed lines are for above-plume photolysis rates, while solid lines are for the middle of the plume, and dotted lines are for the bottom of the plume (see the legend in Fig. 4).

vs. time after emissions. The EnR is defined as the ratio of the enhancement of a species X within the smoke plume ($\Delta X = C_x - C_x^a$; Akagi et al., 2011) to the enhancement of a less reactive species, such as CO$_2$ or CO. The choice of whether to use CO or CO$_2$ in the denominator of the EnR was made on a species-by-species basis to match the choices made in Table 2 of Akagi et al. (2012), which were in turn chosen to minimize the impact of measurement and sampling errors on the EnRs. We can see that the range of dilution rates and photolysis rates simulated for this case capture the general rate of the secondary formation of O$_3$ and PAN, but ASP v2.1 appears to be overestimating the rate of for-

mation of these compounds. This is in contrast to Alvarado and Prinn (2009), who found that ASP v1.0 underestimated the much faster O$_3$ formation in the Timbavati savannah fire smoke plume by about 50 %.

The ASP v2.1 value for the best estimate dilution and photolysis case (i.e., the best-fit dilution combined with the middle of the plume photolysis rates, plotted as a solid black line in Fig. 5) at 4.5 h for ΔO$_3$/ΔCO is 0.116 mol mol^{-1} which is within the uncertainty associated with the average value measured for the Williams fire (0.095 ± 0.022). This overestimate is similar to the positive bias ($\sim$ 25 %) of the SAPRC-07 mechanism vs. the EPA smog chamber results for low ROG / NO$_x$ (< 20 ppb C ppb^{-1} N^{-1}) ratios when aromatics are not part of the surrogate. As aromatics are a minor constituent in biomass burning smoke, and the ROG / NO$_x$ ratio for savannah/scrubland fires like the Williams fire (without including unidentified species) is $\sim$ 10 ppb C ppb^{-1} N^{-1}, we would expect the mechanism in ASP v2.1 to show a similar positive bias. ASP v2.1 predicts an "average" value of ΔPAN / ΔCO$_2$ of 8.4×10^{-4} at 4.5 h downwind, 65 % larger than the observed value of $(5.10 \pm 1.21) \times 10^{-4}$.

Figure 6 shows the ASP v2.1 results and observations for ΔNO$_x$/ΔCO$_2$ and ΔC$_2$H$_4$/ΔCO vs. time after emission. Figure 6a shows that the ΔNO$_x$/ΔCO$_2$ values are correctly simulated by ASP v2.1, with the best estimate dilution and photolysis case EnR of 3.4×10^{-4} matching the observed value of $4.6 \pm 2.3 \times 10^{-4}$ 4 to 4.5 h downwind. However, the observations show a faster rate of decay in the first 2 h after emission than is seen in the model results. Figure 6b shows that the decay of C$_2$H$_4$ is also well matched by the model results, suggesting that the modeled OH is similar to the actual OH concentrations. This can also be seen by comparing the modeled OH concentration for the best estimate dilution and photolysis case (5.3×10^6 molecules cm^{-3}) to that derived by Akagi et al. (2012) using the observed decay of C$_2$H$_4$ ($5.27 \pm 0.97 \times 10^6$ molecules cm^{-3}). This is again in contrast with Alvarado and Prinn (2009), who found that ASP underestimated the observed OH radical concentrations for the Timbavati smoke plume (1.7×10^7 molecules cm^{-3}; Hobbs et al., 2003).

We can explore this contrast further by looking at the rate of loss of HONO in the smoke plume, shown in Fig. 7. Note that unlike the previous figures, Fig. 7 only shows the first hour after emission as the observations showed no detectable HONO further downwind. As noted in Sect. 1, to explain the underestimate of O$_3$ and OH in the Timbavati fire, Alvarado and Prinn (2009) posited that a heterogeneous reaction of NO$_2$ to make HONO and HNO$_3$ was taking place in that plume. However, the O$_3$ and OH results for the Williams fire show no evidence of this chemistry, and the HONO decay seen in Fig. 7 also shows little evidence for a secondary source of HONO except for a few points within the first 12 min after emission that have more HONO than is predicted by the model. While explaining the discrepancy between the Williams and Timbavati results is beyond
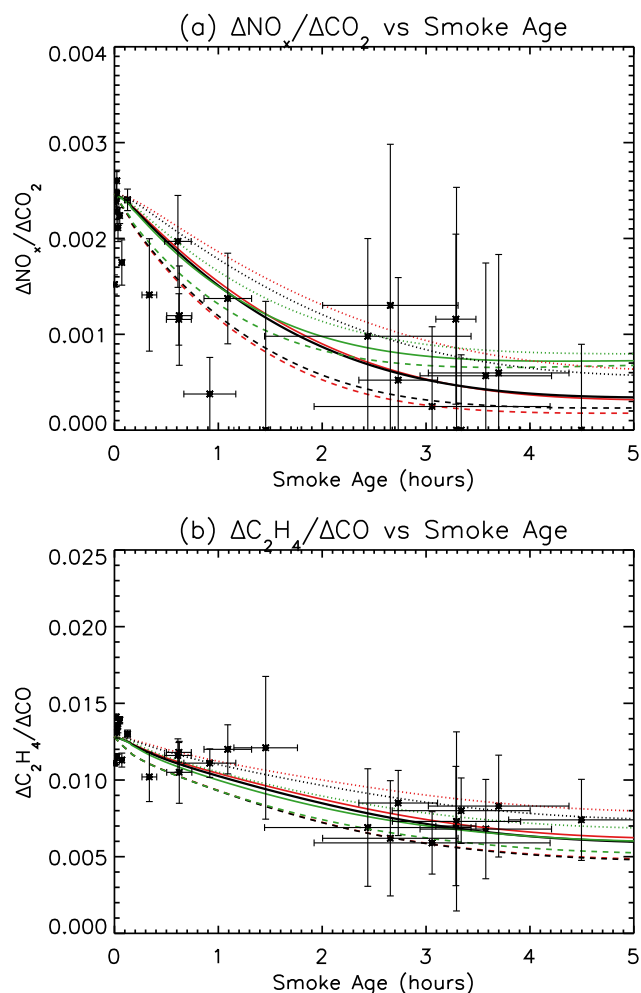
Figure 6. (a) NO_x enhancement ratio (EnR; $mol\,mol^{-1}$) to CO_2 vs. estimated smoke age when the chemistry of the unidentified SVOCs is not included in the model. **(b)** EnR of C_2H_4 to CO vs. estimated smoke age. Asterisks are the measured mixing ratios, with the horizontal error bars showing the uncertainty in the estimated age and the vertical error bars showing the uncertainty in the measurement. Red, black, and green are ASP results for the slow, best-fit (medium), and fast plume dilution rates. Dashed lines are for above-plume photolysis rates, while solid lines are for the middle of the plume, and dotted lines are for the bottom of the plume (see the legend in Fig. 4).

Figure 7. HONO mixing ratio (ppbv) vs. estimated smoke age for the first hour after emission (note difference in x axis scale from Figs. 4–6) when the chemistry of the unidentified SVOCs and a downwind HONO source is not included in the model. Asterisks are the measured mixing ratios, with the horizontal error bars showing the uncertainty in the estimated age and the vertical error bars showing the uncertainty in the measurement. Red, black, and green are ASP results for the slow, best fit (medium), and fast plume dilution rates. Dashed lines are for above-plume photolysis rates, while solid lines are for the middle of the plume, and dotted lines are for the bottom of the plume (see the legend in Fig. 4).

the scope of this paper, we note that the Timbavati fire took place closer to the Equator ($24°$ S vs. $35°$ N), earlier in the year (7 September vs. 17 November) than the Williams fire, and that the RH was higher as well (45.0 % vs. 15.7 %). All of these differences would tend to increase photolysis rates and the formation of OH. In addition, the higher actinic flux and RH in Timbavati may have increased the speed of reactions for forming HONO from NO_2 that are not included in standard chemical mechanisms, either via aqueous chemistry (Jacob et al., 2000), sunlight-activated humic acid sur-
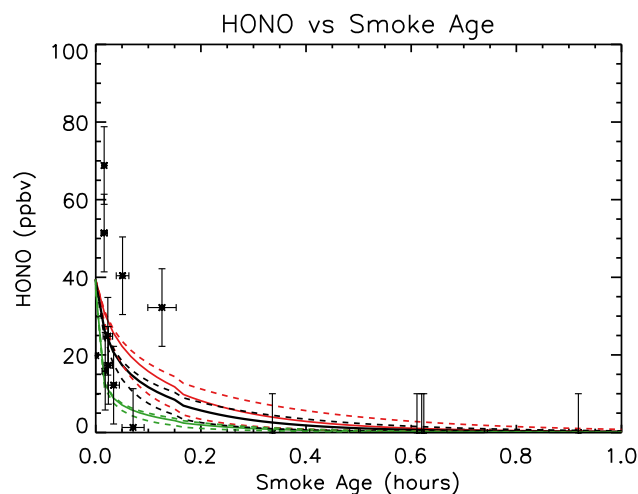
faces (Stemmler et al., 2006, 2007), or photo-excited NO_2 reacting with H_2O (Ensberg et al., 2010). Though we find no evidence for secondary HONO production in the Williams fire data, this does not preclude that some HONO was made, but remained below the AFTIR detection limit of 10 ppbv as the plume diluted.

Figure 8 shows the ASP results for two aldehydes, HCHO and glycoaldehyde ($HCOCH_2OH$), and two organic acids, HCOOH and CH_3COOH, in the Williams fire plume in terms of EnRs to CO. We can see that ASP generally underestimates the formation of these species. Part of this underestimate may be due to errors in the chemical mechanism for known precursor compounds, as was seen for HCHO in the smog chamber results, but neglecting the chemistry of the SVOCs and their ability to form these smaller organic compounds is also likely responsible for this underestimate.

Figure 9 shows the modeled OA enhancement ratios ($\Delta OA\,/\,\Delta CO_2$, $g\,g^{-1}$) at 4.5 h downwind using the parameters listed in Table 2 in addition to the observed average OA enhancement ratio ($2.83 \pm 1.08 \times 10^{-3}$) and the modeled OA enhancement ratio for the case where the chemistry of the unidentified SVOCs is not included (2.27×10^{-3}). When SVOC chemistry was not included, some of the original OA evaporated into the gas phase as the plume diluted, and as there was no chemistry to make these SVOC species less volatile, they stayed in the gas phase leading to a net de-
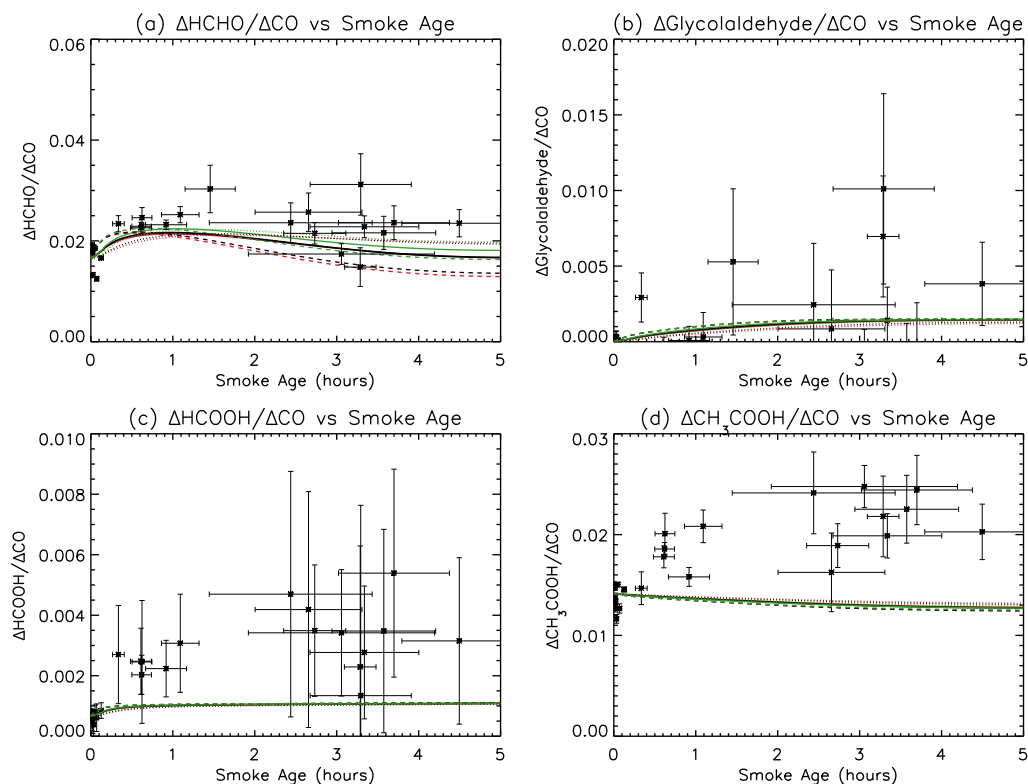
Figure 8. Enhancement ratio (EnR; mol mol^{-1}) of **(a)** HCHO, **(b)** glycoaldehyde (HCOCH$_2$OH), **(c)** formic acid (HCOOH), and **(d)** acetic acid (CH$_3$COOH) to CO vs. estimated smoke age when the chemistry of the unidentified SVOCs is not included in the model. Asterisks are the measured mixing ratios, with the horizontal error bars showing the uncertainty in the estimated age and the vertical error bars showing the uncertainty in the measurement. Red, black, and green are ASP results for the slow, best fit (medium), and fast plume dilution rates. Dashed lines are for above-plume photolysis rates, while solid lines are for the in-plume rates, as described in the text.

crease in $\Delta OA / \Delta CO_2$ with time. Changing the gas-phase concentrations of the unidentified SVOC by $\pm 50\%$ has a small impact ($\sim 3\%$) on these results, but the match between the model and observation could be improved by using a less volatile POA distribution than that given by Grieshop et al. (2009a). However, the modeled decrease without SVOC chemistry is larger (but still within the error bars) than the decrease that was reported by Akagi et al. (2012). In addition, the assumption that the SVOCs do not react is unrealistic – as large multi-functional organic compounds, they should have a relatively fast reaction rate with OH (see Sect. 5.3 below). Thus, in Sect. 5.3 through 5.5 we test different, more realistic implementations for the chemistry of these SVOCs.

5.3 OH reaction rate and fragmentation probability of the unidentified SVOCs

Here we evaluate the ability of the parameters from the original VBS paper of Robinson et al. (2007), a study of SOA formation in wood smoke by Grieshop et al. (2009a, b), and the implementation of the VBS scheme into WRF-Chem by Ahmadov et al. (2012) to simulate the observed evolution of OA in the Williams fire plume. Table 2 shows
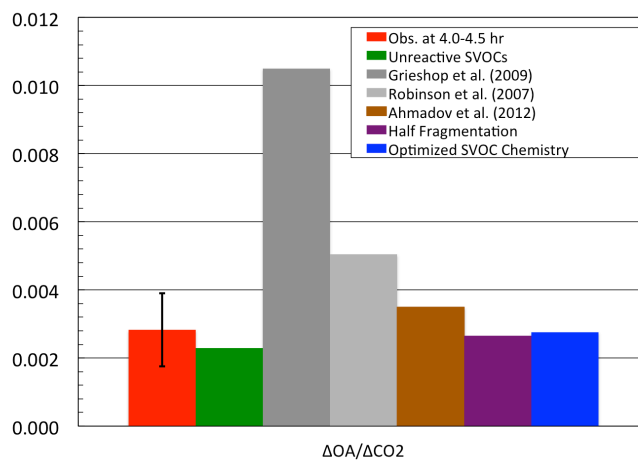

Figure 9. Enhancement ratio (EnR; g g^{-1}) of organic aerosol (OA) to CO$_2$ after 4 to 4.5 h of smoke aging. The error bars on the observed values are based on the 36% uncertainty in the AMS observations of OA. All model results assume the best-estimate dilution rate and the photolysis rates corresponding to the middle of the plume (solid black line in Fig. 4).

the values for the parameters in Reactions (R3) and (R4) that define these various SVOC mechanisms. Figure 9 shows that the SVOC mechanisms of Robinson et al. (2007) and Grieshop et al. (2009a, b) overestimated the OA downwind by a factor of 1.8 and 3.7, respectively. This is primarily due to their relatively large values for k_{OH}. For the Grieshop et al. (2009a, b) case, the overestimation is also partially due to the large increase in mass (μ) and decrease in volatility (n) for each OH reaction. The OA formed using these mechanisms can be reduced by a further 25 % if we assume the unidentified SVOCs are mainly the more volatile SVOC$_9$ ($C^* = 10^6\,\mu g\,m^{-3}$) instead of SVOC$_8$ ($C^* = 10^5\,\mu g\,m^{-3}$), but are fairly insensitive to errors in the POA volatility distribution. The scheme of Ahmadov et al. (2012), with $k_{OH} = 10^{-11}\,cm^3\,molecule^{-1}\,s^{-1}$, was consistent with the uncertainty in the observations, but slightly higher than the observed value (3.48×10^{-3} vs. the observed value of $2.83 \pm 1.08 \times 10^{-3}$). One approach to further reduce the modeled OA would be to reduce k_{OH} even further. However, it seems unlikely that the average OH reaction rate of the unidentified SVOC species would be less than $10^{-11}\,cm^3\,molecule^{-1}\,s^{-1}$, as this is close to the reaction rate for large alkanes (k_{OH} (298 K) $= 1.1 \times 10^{-11}\,cm^3\,molecule^{-1}\,s^{-1}$; Goliff et al., 2013) and the presence of other functional groups (double bonds, aldehydes) would be expected to result in even higher reaction rates. For example, α-pinene has a k_{OH} (298 K) of $5.0 \times 10^{-11}\,cm^3\,molecule^{-1}\,s^{-1}$ (Goliff et al., 2013), and other monoterpenes can have even faster reaction rates with OH. Thus, we think a more likely explanation for the remaining overestimate is that a substantial fraction of the SVOC and OH reactions resulted in the fragmentation of the primary SVOC into more volatile compounds, as in the 2D-VBS schemes of Jimenez et al. (2009) and Donahue et al. (2011).

Figure 9 shows that a k_{OH} of $10^{-11}\,cm^3\,molecule^{-1}\,s^{-1}$ and a fragmentation probability of 50 % (the "half fragmentation" case; see Table 2) provided a reasonably good match with the observed $\Delta OA / \Delta CO_2$ 4.5 h downwind in the smoke plume (2.63×10^{-3} vs. the observed value of $2.83 \pm 1.02 \times 10^{-3}$). Here we assumed that the SVOC fragmented into a small VOC and another, more volatile, SVOC, as in Reactions (R3) and (R4). While this is a relatively large fragmentation probability, we note that it seems reasonable given the likely complex and multifunctional nature of the unidentified SVOCs in a biomass burning smoke plume.

This fragmentation of the SVOCs after reaction with OH could also help to explain the underestimate of aldehydes and organic acids seen in Sect. 5.2 when SVOC chemistry was neglected. For example, Fig. 10 shows the ASP modeled EnR of CH$_3$COOH when we assumed that the VOC fragment produced in Reaction (R4) is CH$_3$COOH. This provided a remarkably good match with the observed CH$_3$COOH formation, providing additional evidence to support the fragmentation hypothesis. While we are not claiming to have proven this is the source of the missing CH$_3$COOH, we note
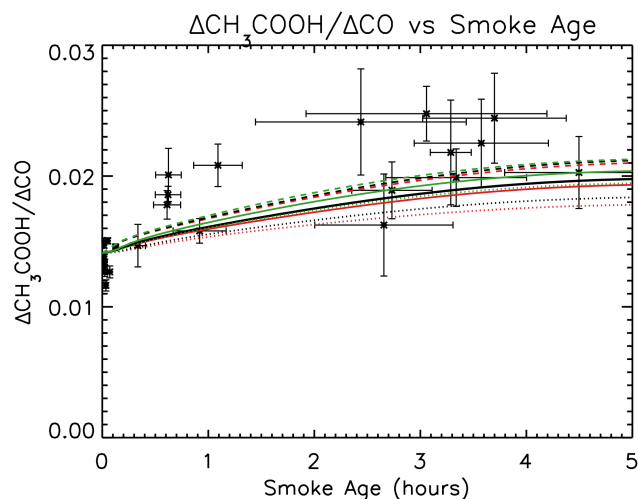

Figure 10. As in Fig. 8d, but for the "half fragmentation" SVOC mechanism (see Table 1) where the VOC fragment produced by fragmentation of the parent SVOC is assumed to become acetic acid (CH$_3$COOH).

that the fragmentation hypothesis is thus consistent with the initial underestimate of the secondary formation of aldehydes and organic acids in ASP v2.1. In addition, there is some evidence from biomass burning plume observations that the formation of CH$_3$COOH and OA are inversely correlated with each other. In the Yucatan plume studied by Yokelson et al. (2009), a large amount of SOA was formed, but CH$_3$COOH did not increase downwind, while in the Williams fire, CH$_3$COOH increased, but OA did not. Thus, the limited amount of relevant airborne data in BB plumes is so far consistent with the idea that the branching between functionalization and fragmentation in BB plumes is variable and future work should identify what environmental and combustion factors control the outcome.

An additional potential explanation for the SOA overestimate observed when the schemes of the Robinson et al. (2007), Grieshop et al. (2009a, b), and Ahmadov et al. (2012) were used is that the OA was becoming more viscous and "glassy" with time (i.e., the particles had a lower bulk diffusivity), thereby reducing the amount of quasi-liquid OA for SVOC compounds to dissolve into (e.g., Kidd et al., 2014; Zaveri et al., 2014). There has been some recent evidence for this process occurring in smoke plumes from biomass burning in the western USA (A. Sedlacek, personal communication, March 2014). ASP v2.1 is not able to examine this possibility in detail, but we do note that while the formation of "glassy" OA would reduce SOA formation, it likely would not increase the formation of aldehydes or organic acids as in the fragmentation hypothesis.
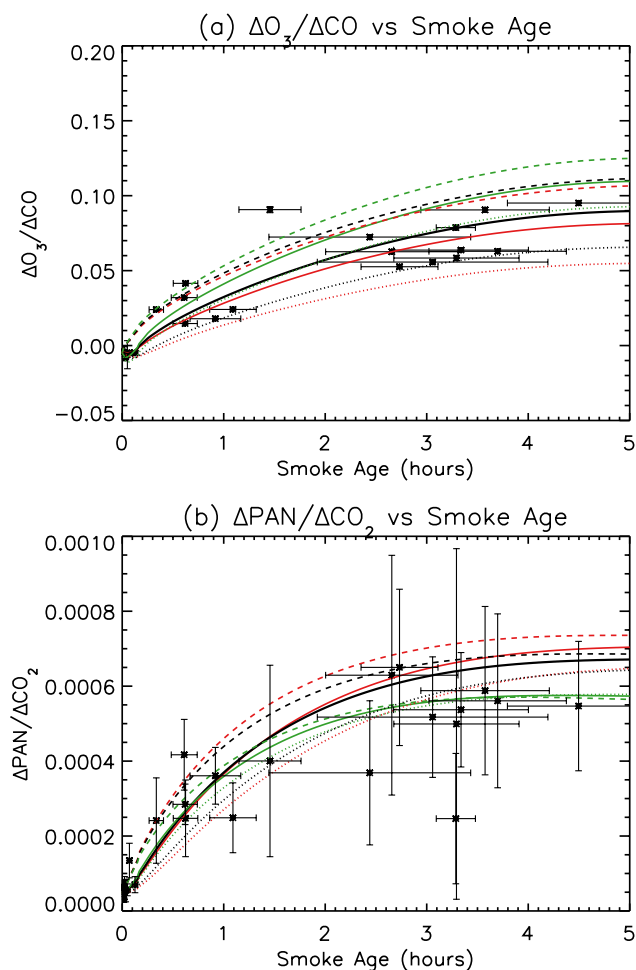
Figure 11. As in Fig. 5, but for the "half fragmentation" SVOC mechanism rather than no fragmentation (see Table 2).

Figure 12. As in Fig. 6, but for the "half fragmentation" SVOC mechanism (see Table 2).

5.4 HO$_x$ and NO$_x$ chemistry of the unidentified SVOCs

Section 5.3 showed that an SVOC mechanism following Reaction (R2) with a k_{OH} of 10^{-11} cm^3 molecule^{-1} s^{-1} and a fragmentation probability α of up to 0.5 (the "half fragmentation" scheme in Table 2) could explain the observed evolution of OA in the Williams fire. However, neglecting the regeneration of HO$_x$ and reaction of the peroxy radical with NO, as in Reaction R2, can lead to substantial underestimates of OH in the concentrated smoke plumes. This is because including Reaction (R2) in ASP leads to a loss of OH with no corresponding regeneration of HO$_2$.

For example, Fig. 11 shows that using the "half fragmentation" scheme reduced the ASP v2.1 estimates of the enhancement ratios of O$_3$ and PAN downwind by 24 and 23%, respectively (for the best estimate dilution and photolysis case, the black line in Fig. 5), while Fig. 12 shows that it increased the ASP v2.1 estimates of C$_2$H$_4$ and NO$_x$ downwind by 33 and 151%, respectively. The "half fragmentation" O$_3$ and PAN estimates are more consistent with the observations –
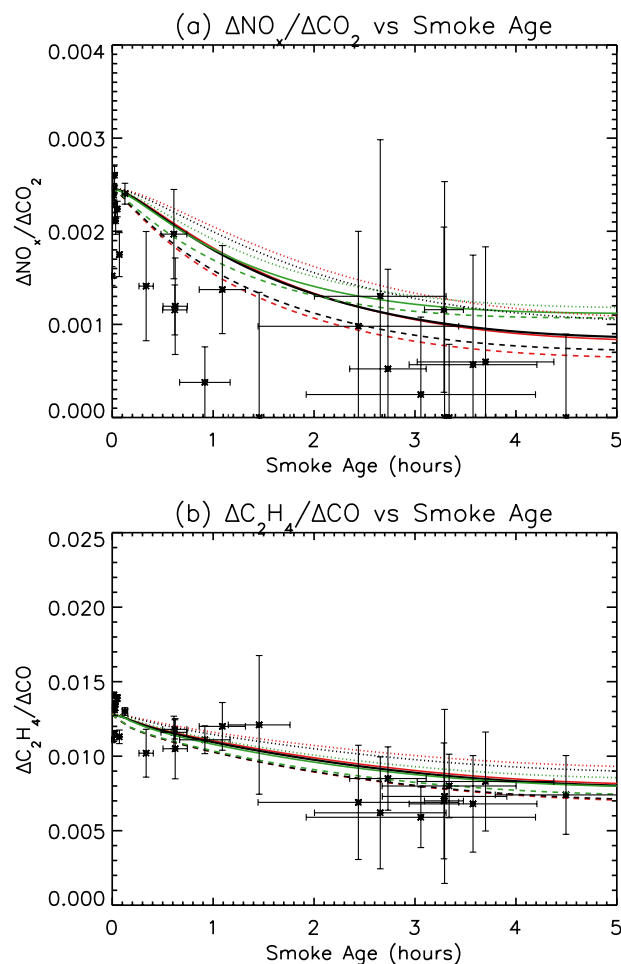
the overestimate of PAN seen with the unidentified SVOC chemistry neglected in Sect. 5.2 has disappeared – but the large overestimate of NO$_x$ (i.e., underestimate of NO$_x$ loss in the plume presumably due to missing organic nitrogen formation) is a serious problem. While this underestimate of NO$_x$ loss reduces the amount of O$_3$ and PAN formed within 5 h after emission, it would lead to large overestimates of the impact of biomass burning plumes on O$_3$ and PAN formation further downwind.

In addition, the chemistry of Reaction (R2) is unrealistic, in that it implies a total loss of OH and no effect of the SVOC oxidation on NO$_x$. One approach for addressing the first concern is to artificially regenerate the OH by simply adding it as an additional product to Reaction (R2). While this makes sense as a "first do no harm" modeling approach to keep the gas-phase results the same regardless of the SVOC scheme, it is equally unrealistic, as it assumes that the SVOCs are oxidized without having any impact on NO$_x$ or HO$_x$.

We prefer to approach this problem by recognizing that SVOCs are going to have impacts on the HO$_x$ and NO$_x$ rad-

ical budgets just like any other organic species, and that this chemistry can be approximated via Reactions (R3) and (R4). Including this more realistic, yet still simplified, chemistry allows ASP to simultaneously simulate the observed changes in OA and O_3 while still making reasonable, chemically plausible assumptions about the chemistry of the unidentified SVOCs emitted by the fire.

Our approach thus used the observations of OA, O_3, PAN, NO_x, and C_2H_4 in the Williams fire as constraints on β and δ, the amount of NO_2 and HO_2 produced in Reaction (R4), respectively, while assuming that $\chi = 1$ throughout. For the Williams fire, we know from the above results that we want the optimized SVOC chemistry to (a) increase O_3, PAN, and OA formation as little as possible; (b) increase the loss of NO_x, either through organic nitrate formation or increased OH concentrations; and (c) increase the OH concentration, thereby increasing C_2H_4 loss. We found that using the parameters for large alkanes from RACM2 ($\delta = 0.63$ and $\beta = 0.74$) generally produced too much O_3 and PAN and too little OH, but did a reasonable job for NO_x loss. However, attempts to increase OH by increasing δ led to too much O_3 formation except for unrealistically low values of β (~ 0.1). Thus, we set $\delta = 0.6$ and reduced β to 0.5, implying that 1.1 O_3 is formed per molecule of SVOC reacted. These parameters (arrived at by trial and error) appear to give the best balance of reducing modeled NO_x and C_2H_4 mixing ratios while minimizing the increase in O_3, PAN, and OA. The following section discusses the ASP v2.1 model results for the Williams fire smoke plume using these parameters in detail. Note that while slightly different, more precise parameters might provide a slightly better match with observation, our goal here is not to derive exact best-fit parameters, but rather to roughly classify the average chemistry of the SVOCs in the Williams fire smoke plume, both for modeling this fire and for future comparisons with other smoke plumes.

5.5 Results with optimized SVOC chemistry

Figure 9 shows the $\Delta OA / \Delta CO_2$ 4.5 h downwind in the smoke plume using the optimized SVOC chemistry discussed in Sect. 5.4. For the best estimate dilution and photolysis model case (i.e., the solid black line in Fig. 5), $\Delta OA / \Delta CO_2$ is 2.75×10^{-3} ($g\,g^{-1}$), very close to the observed value of $2.83 \pm 1.02 \times 10^{-3}$. As in Section 5.2, changing the gas-phase concentrations of the unidentified SVOC by $\pm 50\%$ has a small impact ($\sim 3\%$) on these results, as does assuming that the unidentified SVOCs are mainly the more volatile $SVOC_9$ ($C^* = 10^6\,\mu g\,m^{-3}$) instead of $SVOC_8$ ($C^* = 10^5\,\mu g\,m^{-3}$). However, this result is still sensitive to the POA volatility distribution – for example, moving all the mass in $SVOC_3$ ($C^* = 1\,\mu g\,m^{-3}$) to $SVOC_7$ ($C^* = 10^4\,\mu g\,m^{-3}$) decreases the modeled $\Delta OA / \Delta CO_2$ downwind by 12 % for this case.

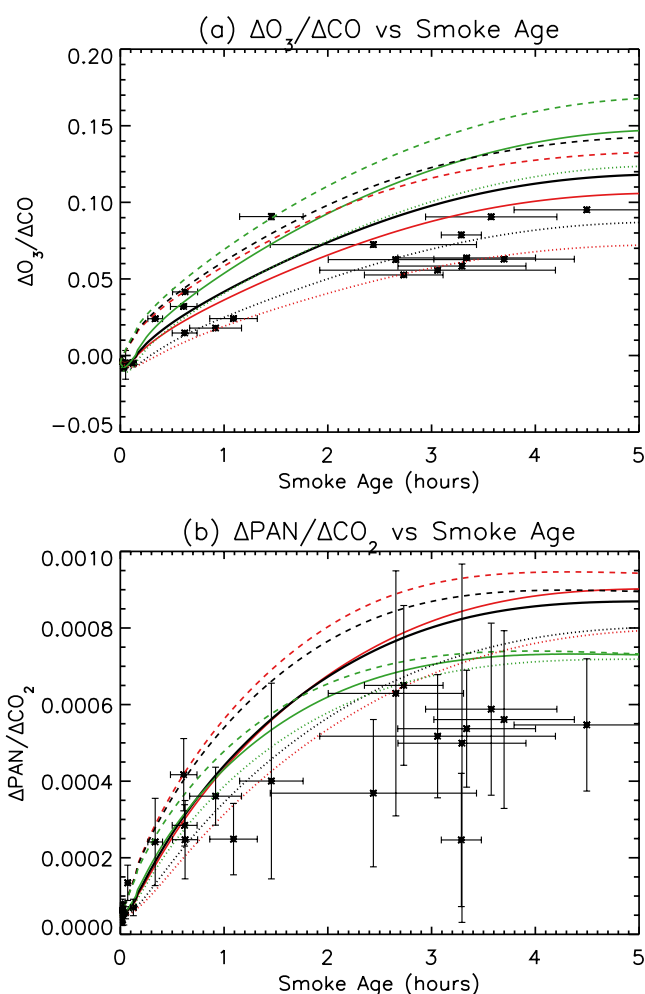Figure 13 shows the enhancement ratios of O_3 and PAN for the optimized SVOC chemistry, and Fig. 14 shows the

Figure 13. As in Fig. 5, but for the optimized SVOC chemistry (see Table 2).

results for NO_x and C_2H_4. The O_3 results are very similar to those from Sect. 5.2 (where SVOC chemistry was not included in the model), while the PAN results are slightly lower (and closer to the observed values) than in that case. For the best estimate dilution and photolysis model case $\Delta O_3/\Delta CO$ is 0.119 at 4.5 h downwind, about 25 % larger than the observed value of 0.095 ± 0.022, while the $\Delta PAN / \Delta CO_2$ is now 7.56×10^{-4} at 4.5 h downwind, about 48 % larger than the observed value of $(5.10 \pm 1.21) \times 10^{-4}$. However, the O_3 and PAN values are reasonably close given the uncertainties in the concentrations and in the estimated smoke ages for the observations.

The NO_x results were much improved from the "half-frag" case in Sect. 5.4, with the best estimate dilution and photolysis case $\Delta NO_x/\Delta CO_2$ of 1.6×10^{-4} being below the mean observed value of $4.6 \pm 2.3 \times 10^{-4}$, but consistent with the error bars of the individual samples as shown in Fig. 14. We could attempt to get a closer match by increasing β, but at the cost of increases in the modeled O_3,
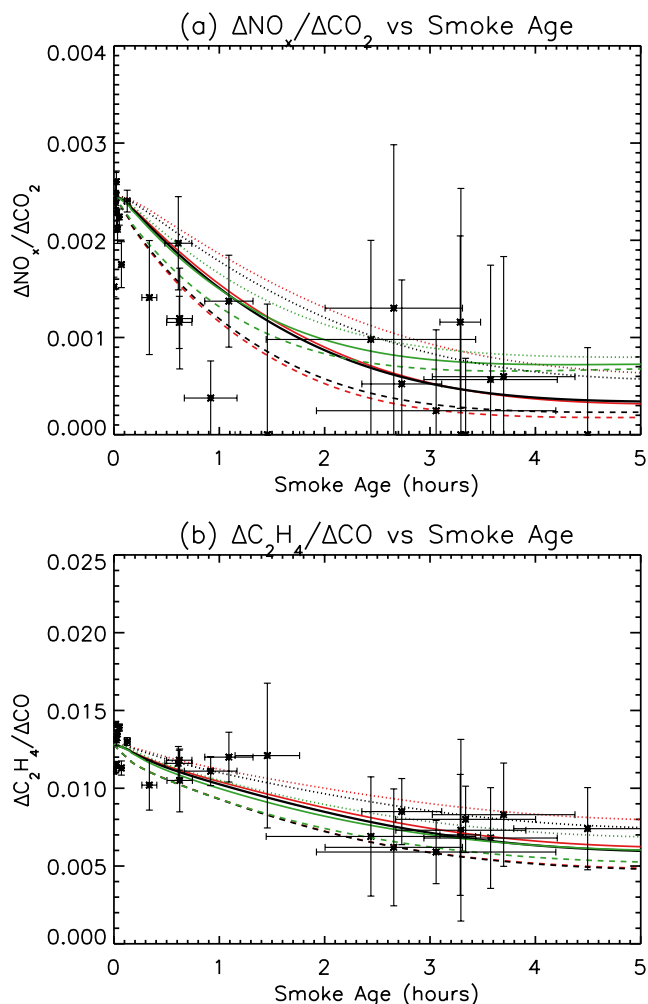
Figure 14. As in Fig. 6, but for the optimized SVOC chemistry (see Table 2).

SVOCs can successfully simulate the observations within the uncertainties of the measurements, the estimated smoke ages of the samples, the plume dilution rate, and the vertical location of the samples in the plume. For the Williams fire, these assumptions were (1) a reaction rate constant with OH of $\sim 10^{-11} \, \mathrm{cm^3 \, s^{-1}}$; (2) a significant fraction (up to $\sim 50\%$) of the $RO_2 + NO$ reaction resulted in fragmentation, rather than functionalization; (3) ~ 1.1 molecules of O_3 were formed for every molecule of SVOC that reacts; and (4) 60 % of the OH that reacted with the SVOC was regenerated as HO_2 by the $RO_2 + NO$ reaction, which implied (5) that 50 % of the NO that reacted with the SVOC peroxy radicals was lost, likely due to organic nitrate formation. However, this chemistry still overestimates PAN formation downwind by about 50 %, suggesting the need for further refinements to the chemistry and estimated emission rates of PAN precursors like acetaldehyde. Furthermore, these specific, quantitative results only apply to the Williams fire analyzed in this paper. Further analysis of other smoke plume observations is needed to determine how these parameters vary between individual smoke plumes.

The method used in this study can provide a way of classifying different smoke plume observations in terms of the average chemistry of their unidentified SVOCs. Similar studies of other young biomass burning plumes would allow us to see how the chemistry of the unidentified SVOCs varies with combustion efficiency, fuel type, and other combustion and environmental parameters, providing an additional constraint on the reactivities of the unidentified SVOCs. These constraints could then provide more insight into the formation of O_3, PAN, and OA in young biomass burning smoke plumes and serve as the basis of parameterizing this process for regional or global-scale models. Future field experiments, focused on quasi-Lagrangian sampling of biomass burning smoke plumes, should ideally also provide data beyond that available for the Williams fire that will increase our understanding of the chemistry of these plumes. These field experiments should include (a) observations of changes in particle size distribution to test model simulations of condensational growth, coagulation, and new particle formation; (b) observations of a larger suite of NO_y species, such as $HNO_{3(g)}$, peroxy nitrates, and alkyl nitrates, for use in studying and constraining the transformations of reactive nitrogen; (c) direct measurements of photolysis rates within the smoke plumes; (d) measurements of organic aerosol volatility, viscosity, and mixing state with black carbon and inorganic aerosols; and (e) more detailed measurements of the currently unidentified organic species present in the smoke plumes, including acetaldehyde, an important PAN precursor.

PAN, and OA formation. The decay of C_2H_4 is also better modeled than in the half fragmentation case, suggesting the model OH is also improved. The modeled OH concentration for the best estimate dilution and photolysis case is now 5.3×10^6 molecules $\mathrm{cm^{-3}}$, matching the observed value of $5.27 \pm 0.97 \times 10^6$ molecules $\mathrm{cm^{-3}}$.

6 Conclusions

We have used version 2.1 of the ASP model, which includes extensive updates to the gas-phase chemistry and SOA formation modules, to simulate the near-source chemistry within the smoke plume from the Williams fire, as sampled by Akagi et al. (2012). We find that the assumptions made about the chemistry of the unidentified SVOCs emitted by the fire have a large impact on the simulated secondary formation of O_3, PAN, and OA within the plume. We showed that reasonable assumptions about the chemistry of the unidentified

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