## SYNTHETIC COBALTOMENITE, CoSeO<sub>3</sub>.2H<sub>2</sub>O, AND RELATED Co-COMPOUNDS AS MODEL PHASES FOR THE APPLICATION OF THE SUPERPOSITION MODEL (SM) OF CRYSTAL FIELDS IN GEOSCIENCES

by

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The Superposition Model (SM) of crystal fields was originally developed to separate the geometrical and physical information inherent in crystal field parameters [1, 2]. It is based on the assumption that the crystal field can be expressed as the sum of axially symmetric contributions of all i nearest neighbour ligands of the transition metal ion. The crystal field parameters  $B_{kq}$ can then be obtained from  $(R)^{t_k}$ 

$$B_{kq} = \sum \overline{B}_{k}(R_{0}) \left(\frac{R_{0}}{R}\right)^{t_{k}} K_{kq}(\Theta_{i}, \Phi_{i})$$

where  $\overline{B}_k$  are the intrinsic parameters (related to a reference metal-ligand distance  $R_0$ ),  $t_k$  are the power law exponent parameters, both for each rank k of the crystal field,  $R_i$  are the individual metal-ligand distances, and  $K_{kq}(\Theta_i, \Phi_i)$  are the coordination factors calculated from the angular polar coordinates of the ligands.

Since its introduction by BRADBURY & NEWMAN [3], the SM has been applied with considerable success for the analysis of lanthanide crystal fields. Contrary, SM investigations on d-block transition elements are scarce up to date. Hence, the necessary intrinsic and power law exponent parameters for application to systems containing d-elements, e.g. in geosciences, are practically missing or were obtained from natural, dilute phases, where the local structure around the particular transition ion is not exactly known, e.g. Cr<sup>3+</sup> in alumosilicates [4, 5]. In order to provide reliable SM parameters of 3d<sup>n</sup> elements for future applications in geosciences, we started to investigate inorganic model compounds, i.e. pure synthetic endmember phases, where the (local) structure is precisely known from concurrent or recent structure investigations.

The first complete SM parameter set for  $Co^{2+}$  ( $\bar{B}_4 = 4740$ ,  $\bar{B}_2 = 7000 \text{ cm}^{-1}$ ;  $t_4 = 3.1$ ,  $t_2 = 5.5$ ) was extracted from polarised electronic absorption spectra of  $Li_2Co_3(SeO_3)_4$ , which is characterised by a strong bond length and angle distortion of its  $CoO_6$  polyhedra with low symmetries  $C_1$  and  $C_i$  [6]. Subsequently, we performed SM analyses of temperature dependent polarised spectra of brucite-type  $Co(OH)_2$  [7], where the  $Co^{2+}$  ions occupy a high-symmetry site ( $D_{3d}$ ) within a compressed hexagonal close packing of oxygen atoms. In general, the magnitude of the intrinsic  $B_k$  parameters was confirmed ( $\bar{B}_4 = 5320$ ,  $\bar{B}_2 = 3900 \text{ cm}^{-1}$ ; the  $t_k$  were not refined due to symmetry restrictions). On closer inspection, however, it appeared that the specific position of closely related ligands within the spectrochemical series (i.e.  $Q_3Se^{2-}$  vs. OH<sup>-</sup>) affects the intrinsic  $\bar{B}_k$ .

Therefore, especially for mixed-ligand coordinations, it might be worthwhile to apply some empirical correction factor for the ligand type in the SM equation – resembling the *f*-factor formerly introduced by JØRGENSEN [8] – instead of refining different SM parameter sets for each ligand type.

Consequently, we modified our program suite [9] (performing the SM calculations with varying parameter sets, preparing the input files for a modified HCFLDN2 crystal field program package by YEUNG [10], and interpreting its output files) in such a way that

$$B_{kq} = \sum_{i} \overline{B}_{k}(R_{0}) \left(\frac{R_{0}}{R_{i}}\right)^{\prime k} K_{kq}(\Theta_{i}, \Phi_{i}) \cdot f$$

where  $f_i$  represents a fixed or adjustable correction factor for different ligand types, according to their position in the spectrochemical series of ligands. By default, fi is set to unity. As a Co<sup>2+</sup> model phase with a distorted mixed-ligand coordination relevant for geosciences, synthetic cobaltomenite, CoSeO<sub>3</sub>·2H<sub>2</sub>O, was chosen. There, the Co<sup>2+</sup> cations are octahedrally coordinated to four oxygen atoms belonging to SeO<sub>3</sub> groups (Co–O = 2.07 - 2.16 Å) as well as two water molecules (Co–O = 2.04 and 2.19 Å) in *cis*-position [11]. However, the calculations reveal only a marginal deviation of  $f_{H2O}$  from unity. In this case, a reason for this unexpected result might be an additional perturbation and hence complication of the crystal field due to the linkage of the CoO<sub>6</sub> octahedra to Co<sub>2</sub>O<sub>10</sub> dimers. Preliminary SM,  $f_{H2O}$ , and interelectronic repulsion parameters for CoSeO<sub>3</sub>·2H<sub>2</sub>O are:  $\mathbf{B}_4 = 5080 \text{ cm}^{-1}$ ,  $\mathbf{t}_4 = 5.4$ ,  $\mathbf{B}_2 = 8000 \text{ cm}^{-1}$ ,  $\mathbf{t}_2 = 0$ ,  $f_{H2O} = 1.01$ , Racah B = 800 cm<sup>-1</sup>, and Racah C = 3300 cm<sup>-1</sup>. With the exception of the second rank exponent  $\mathbf{t}_2$ , which might also be biased by the polyhedral linkage, the obtained SM and Racah parameters comply quite well with the expectations for Co<sup>2+</sup> in oxygen crystal fields.

## References

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